

2nd line:

1. "Operator Formalism" and $\psi(s)$ as a maximal eigenvalue
2. Evaluating $\psi(s)$ using a "cloning algorithm"

"OPERATOR FORMALISM" & $\psi(s)$ AS A MAXIMAL EIGENVALUE

1. a. Generic Observables. System with discrete configurations $\{e\}$
(\hookrightarrow even a finite number)

x Dynamics of the system: given by transition rates $W(e \rightarrow e')$

ie btw t & $t+dt$: (small dt)

$$\begin{cases} e \rightsquigarrow e' & \text{with probability } dt W(e \rightarrow e') \\ e \rightsquigarrow e & \text{[no event] } \text{---} \quad 1 - \underbrace{\sum_{e'} dt W(e \rightarrow e')} \end{cases}$$

$$P(e, t+dt) = \sum_{e'} dt W(e' \rightarrow e) P(e', t) + (1 - dt \sum_{e'} W(e \rightarrow e')) P(e, t)$$

$$dt \rightarrow 0 \hookrightarrow \partial_t P(e, t) = \sum_{e'} W(e' \rightarrow e) P(e', t) - \underbrace{\sum_{e'} W(e \rightarrow e')}_{\substack{\hookrightarrow = \text{the "escape rate"} \\ \text{[} r(e) \text{ is called } \text{]}}} P(e, t)$$

x Description of the dynamics (see Tibbo's lecture)

- . the system stays in conf e for a duration Δt with prob. $r(e) e^{-\Delta t r(e)}$ (\hookrightarrow not small)
- . the system then jumps from $e \rightsquigarrow e'$ with prob. $\frac{W(e \rightarrow e')}{r(e)}$

x Generic observable: $A = \sum_{k=0}^{\infty} \alpha(e_k \rightarrow e_{k+1})$

for a history $e_0 \dots e_k$ of the system;

i.e. at each "jump" (or event) $A \rightsquigarrow A + \alpha(e \rightarrow e')$

1.6 - The simplest case : $A = K$ (activities)

ie, upon each event $K \mapsto K+1$

x evolution of the probability density of being in \mathcal{E} at time t ,
having observed a value K of the activity - $P(\mathcal{E}, K, t)$:

$$\partial_t P(\mathcal{E}, K, t) = \sum_{\mathcal{E}'} P(\mathcal{E}', K-1, t) W(\mathcal{E}' \rightarrow \mathcal{E}) - r(\mathcal{E}) P(\mathcal{E}, K, t) \quad (*)$$

↑
before the "jump"

the problem is difficult to solve
because this term is non-diagonal
along direction K

However not that if one obtains $P(\mathcal{E}, K, t)$, one has $P(K, t) = \sum_{\mathcal{E}} P(\mathcal{E}, K, t)$

x Going to the s-ensemble : As Tobias said: this def. of the generating function is akin to Fourier Transforms.
 $\hat{P}(\mathcal{E}, s, t) = \sum_{K \geq 0} e^{-sK} P(\mathcal{E}, K, t)$

Remarks - this is a "Laplace transform" (\leftrightarrow generating functions)

- by definition $\langle e^{-sK} \rangle = \sum_{K \geq 0} \sum_{\mathcal{E}} e^{-sK} P(\mathcal{E}, K, t) = \sum_{\mathcal{E}} \hat{P}(\mathcal{E}, s, t)$
not normalized!

- More generally, $\langle \Theta(\mathcal{E}) e^{-sK} \rangle = \sum_{\mathcal{E}} \Theta(\mathcal{E}) \hat{P}(\mathcal{E}, s, t)$

From (*) above : $\partial_t \hat{P}(\mathcal{E}, s, t) = \sum_{\mathcal{E}'} e^{-s} W(\mathcal{E}' \rightarrow \mathcal{E}) \hat{P}(\mathcal{E}', s, t) - r(\mathcal{E}) \hat{P}(\mathcal{E}, s, t)$

vector of components $\hat{P}(\mathcal{E}, s, t)$ ↑ matrix
this is the only difference with the $s=0$ dynamics

Linear evolution :

$$\partial_t |\hat{P}(s, t)\rangle = W(s) |\hat{P}(s, t)\rangle$$

$$(W(s))_{\mathcal{E}\mathcal{E}'} = e^{-s} W(\mathcal{E}' \rightarrow \mathcal{E}) - r(\mathcal{E}) \delta_{\mathcal{E}\mathcal{E}'}$$

* Eigenvector & eigenvalue problem :

Let's call $\psi(s)$ the largest eigenvalue¹ of eigenvector $|\hat{P}_0(s)\rangle$

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In the large-time limit :

$$[\partial_t |\hat{P}\rangle = W|\hat{P}\rangle \text{ is solved by } |\hat{P}\rangle = e^{tW} |\hat{P}_{t=0}\rangle]$$

$$|\hat{P}(s, t)\rangle = e^{tW(s)} |\hat{P}(s, 0)\rangle \xrightarrow{t \rightarrow \infty} e^{t\psi(s)} |\hat{P}_0(s)\rangle$$

Hence: all components $\hat{P}(p, s, t)$ evolve exponentially in time $\sim e^{t\psi(s)}$

$$\Leftrightarrow \langle e^{-sK} \rangle \sim e^{t\psi(s)}, \quad \psi(s) = \max_p W(p, s)$$

Maximal eigenvalue

$= \sum_p \hat{P}(p, s, t)$, see p II.2

The problem of determining the fluctuations of K thus amounts to finding the extremal eigenvalue of a matrix.

One can thus borrow methods from quantum mechanics & mathematics where this sort of problems occurs rather regularly.

* Remark: everything we wrote also works for the l.d.f $\psi(s)$

of quantities $A = \sum_{k=0}^{K-1} \alpha(p_k \rightarrow p_{k+1})$:

One simply has to replace $e^{-s} W(p \rightarrow p')$ by $e^{-s\alpha(p \rightarrow p')} W(p \rightarrow p')$

2. EVALUATING $\Psi(s)$ THROUGH A "CLONING" ALGORITHM

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Consider the equation of evolution of $\hat{P}(e, s, t)$ $\left\{ \begin{array}{l} n(e) = \sum_{e'} W(e \rightarrow e') \\ \downarrow \\ e' \end{array} \right.$

$$\partial_t \hat{P}(e, s, t) = \sum_{e'} e^{-s} \underbrace{W(e' \rightarrow e)}_{= W_s(e' \rightarrow e)} \hat{P}(e', s, t) - n(e) \hat{P}(e, s, t)$$

It does not correspond to a Markov evolution of a probability $\hat{P}(e, s, t)$

because: $n(e) = \sum_{e'} W(e \rightarrow e')$ and not $\sum_{e'} W_s(e \rightarrow e')$

(And indeed: $\sum_e \hat{P}(e, s, t) \stackrel{t \rightarrow \infty}{\sim} e^{t\psi(s)}$: "probability is not conserved")

2-a. Rewriting in terms of a population dynamics:

$$= \sum_{e'} W_s(e \rightarrow e')$$

$$\partial_t \hat{P}(e, s, t) = \underbrace{\sum_{e'} W_s(e' \rightarrow e) \hat{P}(e', s, t)}_{\text{probability-preserving evolution with } s\text{-modified rates } W_s(e' \rightarrow e)} - n_s(e) \hat{P}(e, s, t) + \underbrace{(n_s(e) - n(e)) \hat{P}(e, s, t)}_{\text{reproduction at a rate } n_s(e) - n(e)}$$

(probability-preserving) evolution with s -modified rates $W_s(e' \rightarrow e)$

reproduction at a rate $n_s(e) - n(e)$

"mutations"

"selection"

Physical interpretation: [don't forget we study a large deviation \checkmark characterized by s]

- s -dependent transition rates slightly biasing the trajectories ("mutations" allowing to probe the available space of configurations)
- selection rules actually favor the configurations, which render the large deviation (determined by s) typical.

↳ In case of phase coexistence, this allows to select a phase and to study its properties

2. b. Implementation in continuous time :

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Start from a large number N_0 of copies of the system

- evolve each copy of the system with rates $W_s(\mathcal{E} \rightarrow \mathcal{E}')$

↳ time Δt before one exits from a configuration \mathcal{E} distributed exponentially with rate $r_s(\mathcal{E})$

• after this, change $\mathcal{E} \rightarrow \mathcal{E}'$ with probability $\frac{W_s(\mathcal{E} \rightarrow \mathcal{E}')}{r_s(\mathcal{E})}$

- "clone" with rate $r_s(\mathcal{E}) - r(\mathcal{E})$:

on each interval Δt during which \mathcal{E} does not change, replicate the copy by a factor $e^{\Delta t \cdot (r_s(\mathcal{E}) - r(\mathcal{E}))}$.

↳ with those rules, the number $N(\mathcal{E}, t)$ of copies of the system in configuration \mathcal{E} evolves with the same eq. as $\hat{P}(\mathcal{E}, s, t)$

Hence $N(\mathcal{E}, t) \sim e^{\underbrace{t \Psi(\mathcal{E})}_{\substack{\text{measure the exponential evolution} \\ \text{its rate yields } \Psi(\mathcal{E})}}}$

Besides, average among the set of copies yields $\langle \mathcal{O} \rangle_s$.

Remark: in practice, one uses tricks to keep total population constant !

[also linked with Wright-Fisher process - see Tobias Galla's lecture.]

2-c Implementation in discrete time:

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$$\hat{P}(e, s, t+dt) = \sum_{e'} dt W_s(e' \rightarrow e) \hat{P}(e', s, t) + \underbrace{[1 - dt r_s(e)]}_{\substack{\text{make no change} \\ \text{with probability} \\ 1 - dt r_s(e)}} \hat{P}(e, s, t) + dt \underbrace{[r_s(e) - n(e)]}_{\substack{\text{in any case,} \\ \text{duplicate the} \\ \text{copy with} \\ \text{rate } dt(r_s(e) - n(e))}} \hat{P}$$

⚠ CAVEAT: depending on s , the choice of dt might require a very small dt !
→ better to use continuous-time -