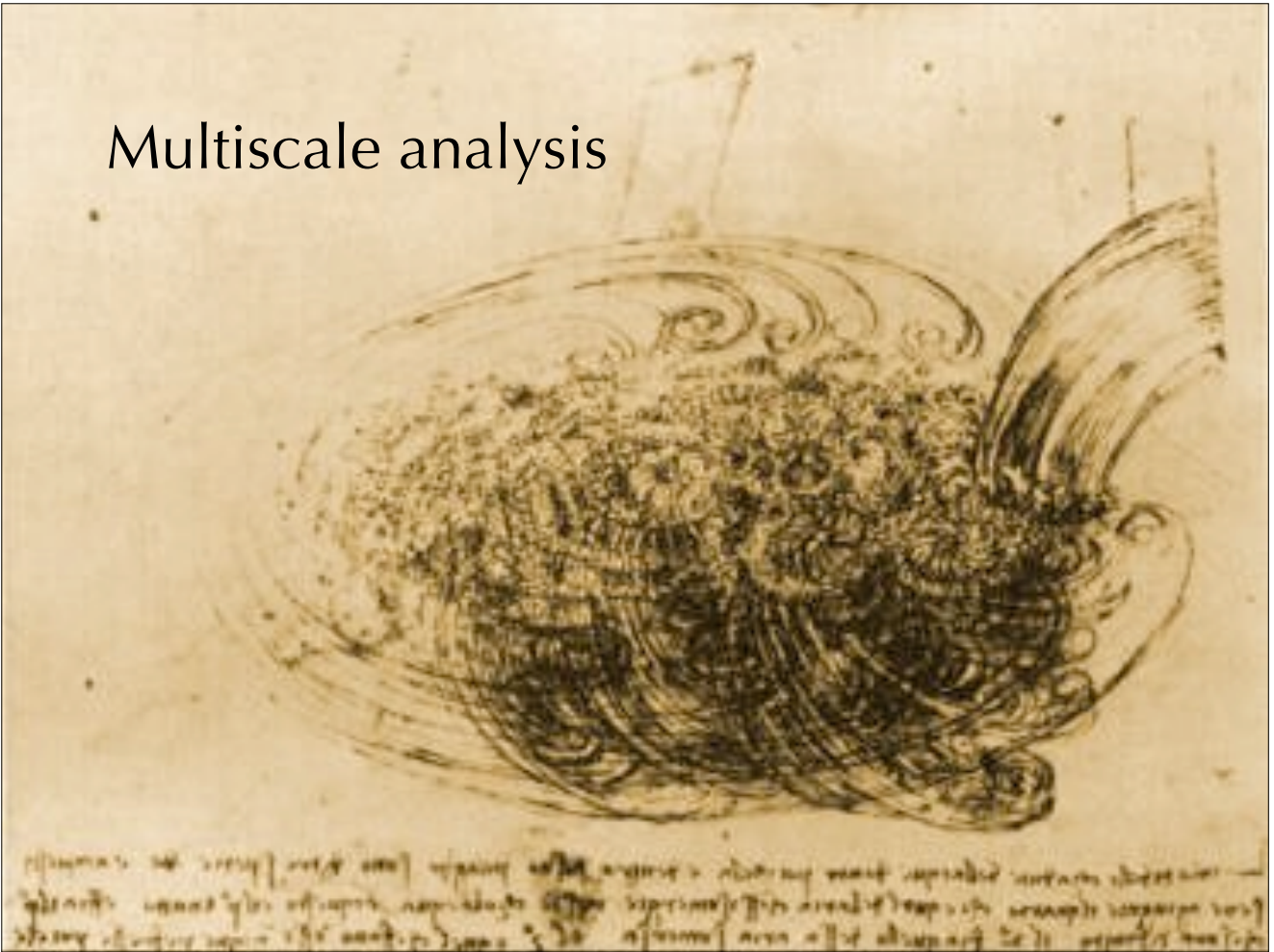


# Multiscale analysis



2

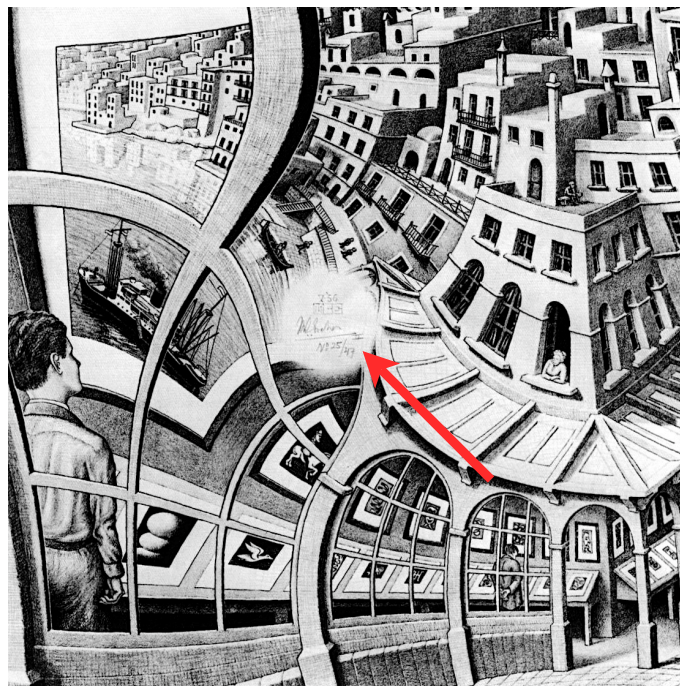


- Most of our data analysis techniques involve strong assumptions on linearity, stationarity, etc.
- Yet the world we live in is
  - non-linear
  - non-stationary
  - non-Gaussian
  - spatio-temporal
  - ...

- The proper understanding of these “*non-properties*” requires new (more advanced) tools.
- Many mistakes have been made by ignoring these non-properties
- These “non-properties” can however give deeper insight into the physics

## An old mystery

- M. C. Escher didn't know how to properly finish his drawing



# An old mystery

- The solution was found last year : self-similarity !

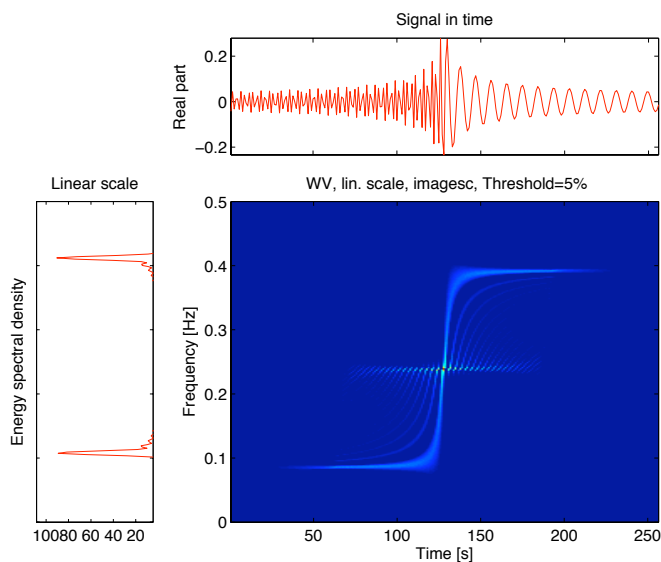


Warwick, 2/2008

# Wigner-Ville

- Unfortunately, *interferences* arise when signals are mixed

$$W_{x+y}(t, \nu) = W_x(t, \nu) + W_y(t, \nu) + 2\Re\{W_{x+y}(t, \nu)\}$$

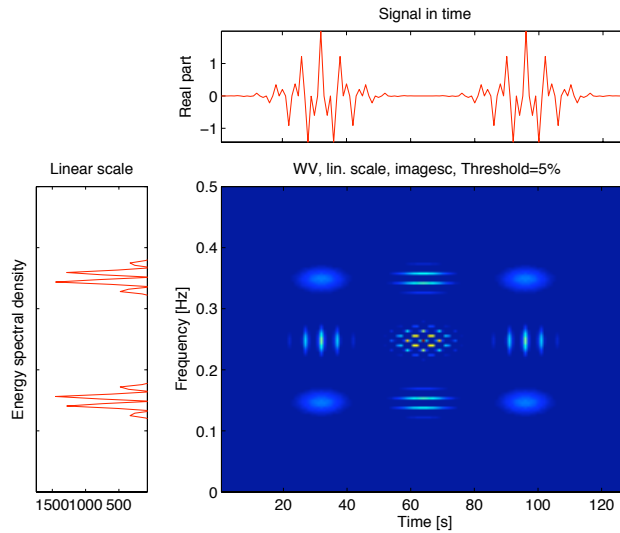


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# Wigner-Ville

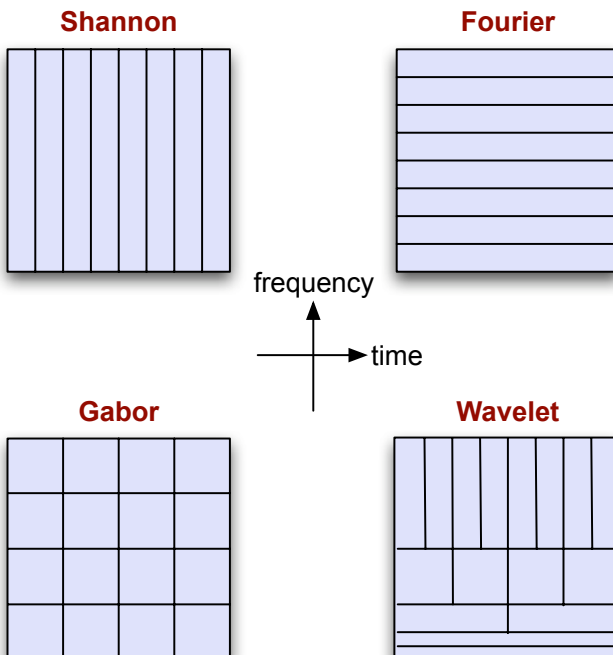
- Various methods have been developed for attenuating interferences in the Wigner-Ville transform

but these cross-terms remain a severe problem



# Tiling the plane

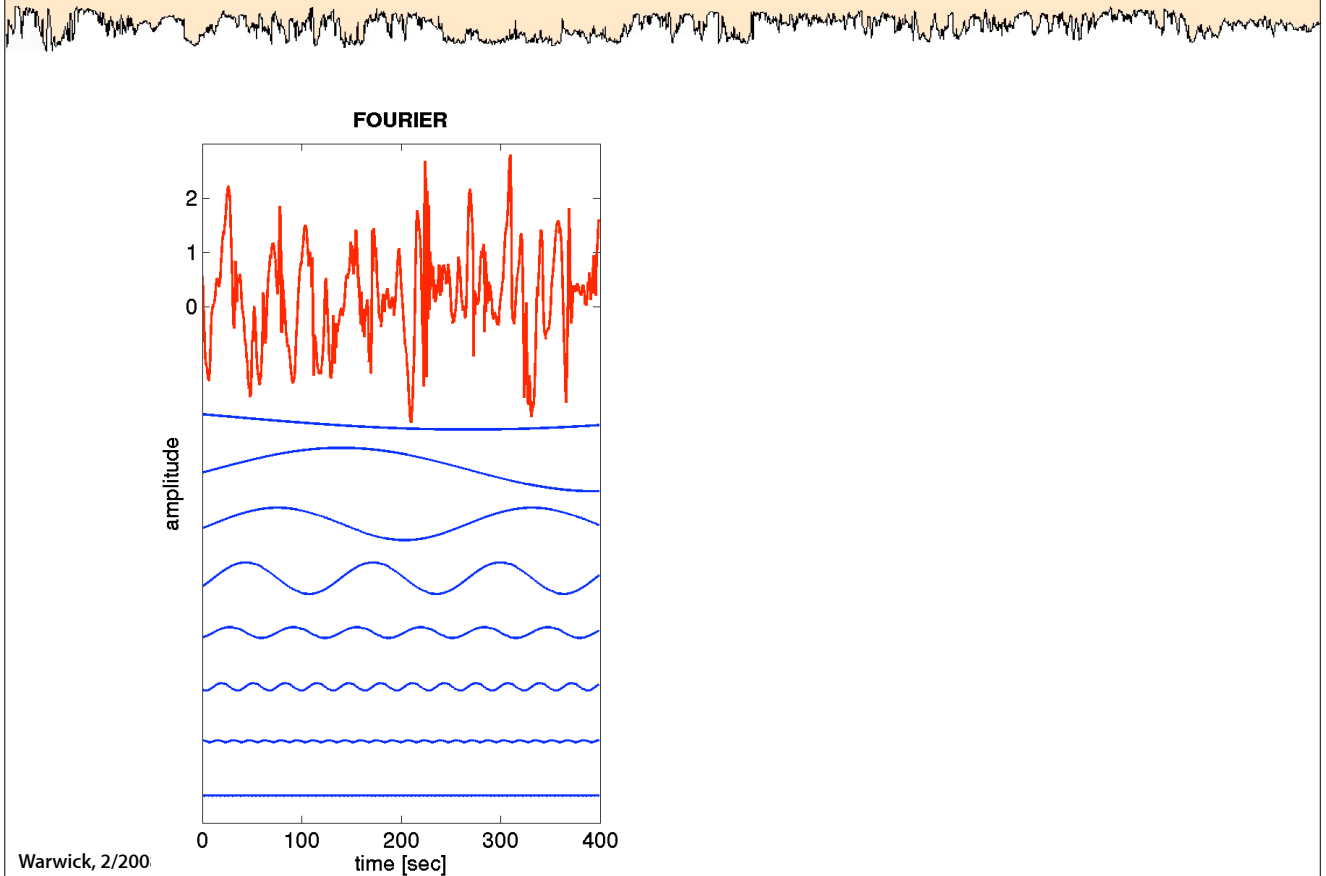
- The problem is : *how to properly partition the time-frequency plane ?*



$$\Delta t \cdot \Delta \nu \geq \frac{1}{4\pi}$$



## Fourier vs multiscale analysis



## Wavelet transform : definition

**Fourier transform**

$$X(\omega) = \langle x(t), e^{j\omega t} \rangle = \int_{-\infty}^{+\infty} x(t) e^{-j\omega t} dt$$

**Continuous wavelet transform**

$$X(\tau, a) = \langle x(t), \psi_{\tau, a}(t) \rangle = \int_{-\infty}^{+\infty} x(t) \psi_{\tau, a}^*(t) dt$$

**mother wavelet**

$$\psi_{\tau, a}(t) = \frac{1}{\sqrt{a}} \psi\left(\frac{t - \tau}{a}\right)$$

**scale**

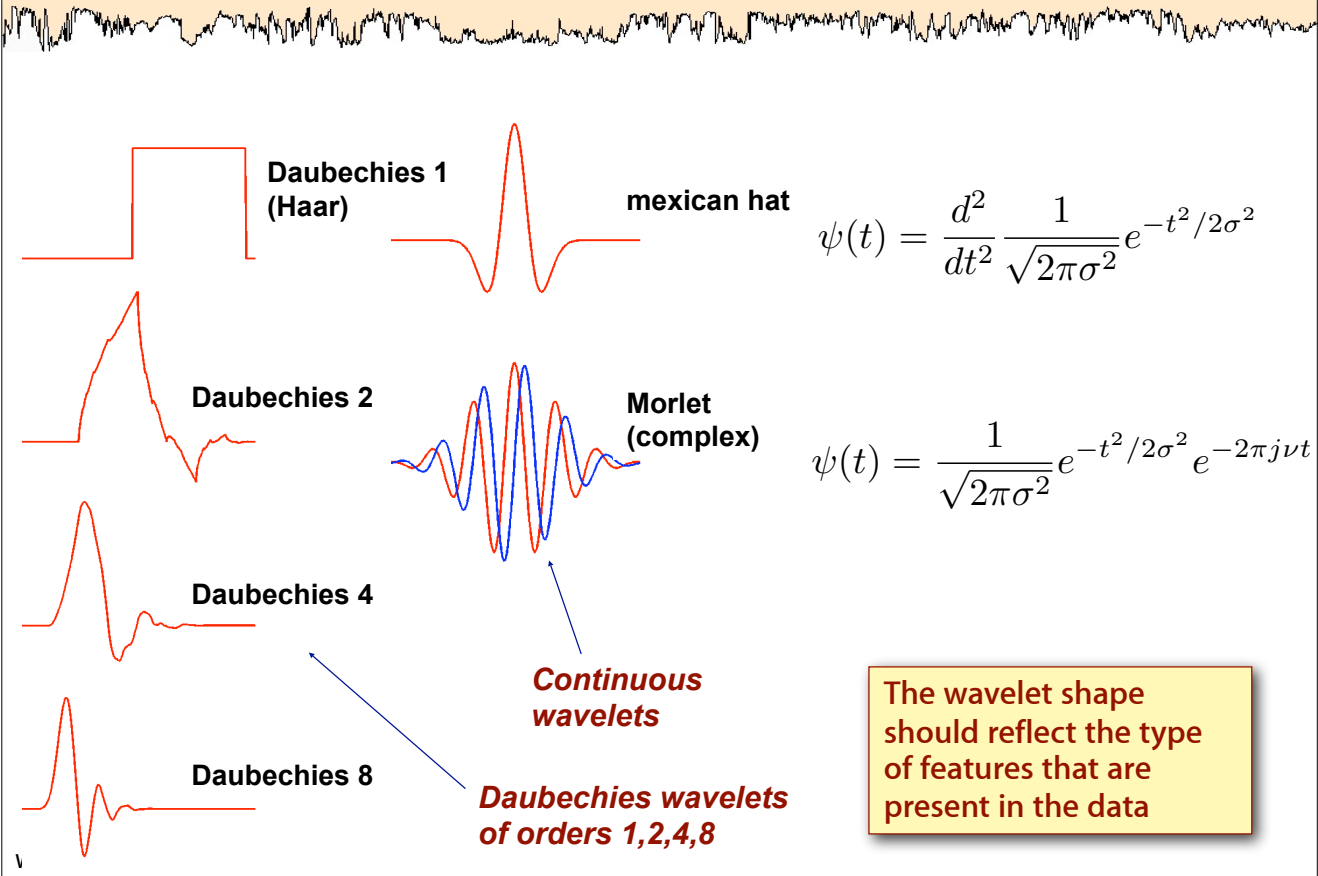
**Discrete wavelet transform**

$$X_{j, k} = \langle x(t), \psi_{j, k}(t) \rangle = \int_{-\infty}^{+\infty} x(t) \psi_{j, k}^*(t) dt$$

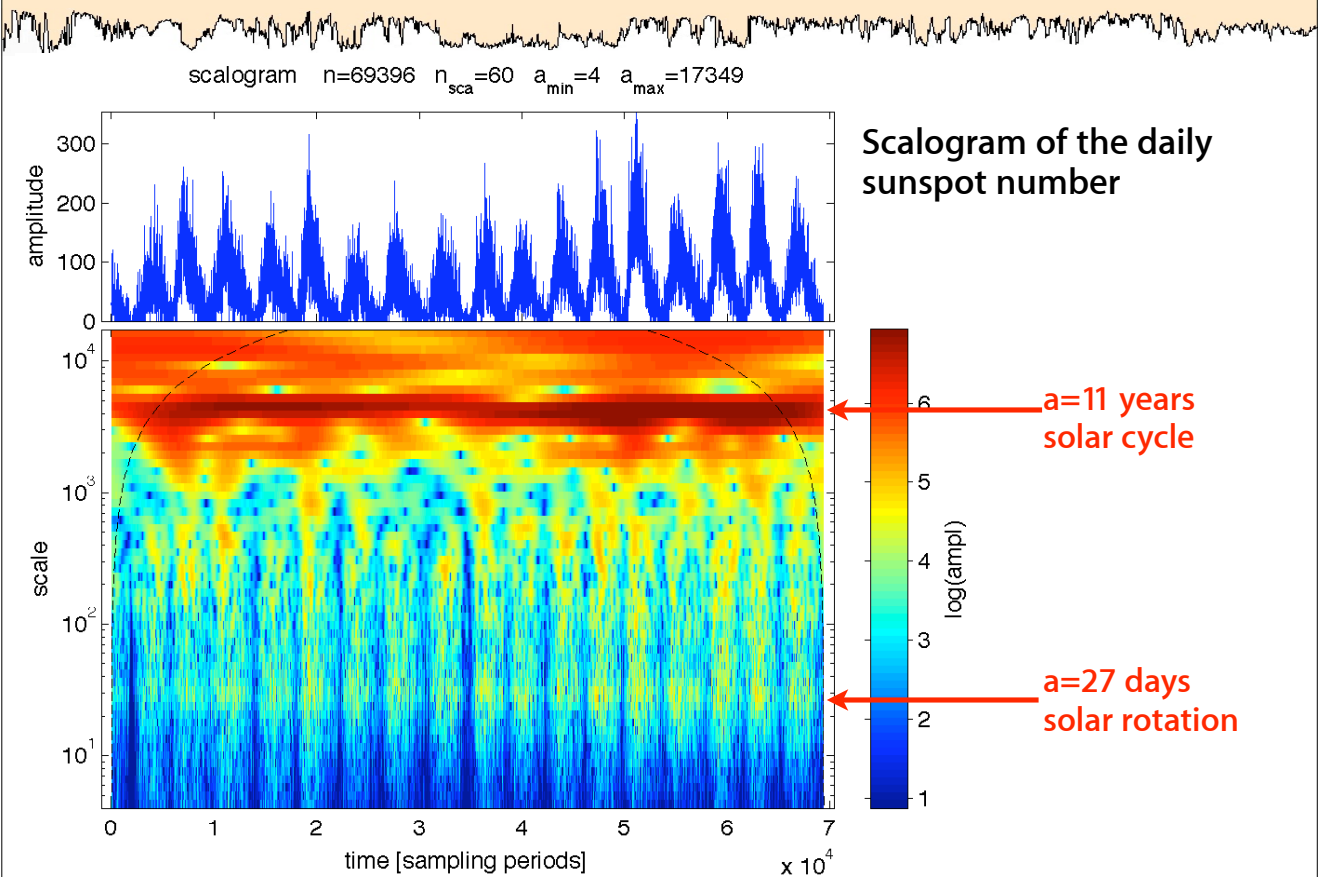
**mother wavelet**

$$\psi_{\tau, a, k}(t) = \frac{1}{2^{j/2}} \psi(2^{-j}t - k)$$

# Mother wavelets : examples



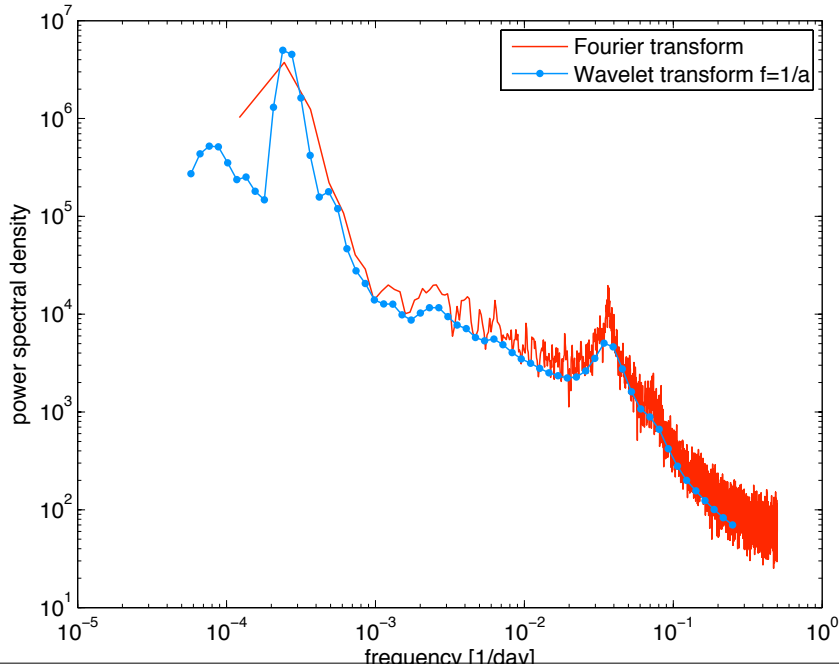
# Scalogram



## Power spectral density



- By integrating  $|\langle x(t), \psi_{a,\tau}(t) \rangle|^2$  over time, we recover a “power spectral density”

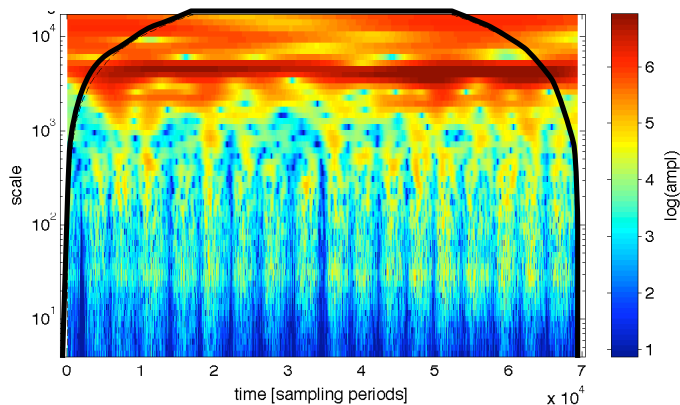


Morlet wavelet transform of the sunspot number, with  $frequency = 1/scale$

## What scales ?



- As in Fourier analysis, the scales are bounded
  - The smallest scale should be  $> 4$  sampling periods (2 for Fourier)
  - The largest scale should typically be  $< 0.25 \times$  sequence length. This defines the **cone of influence**

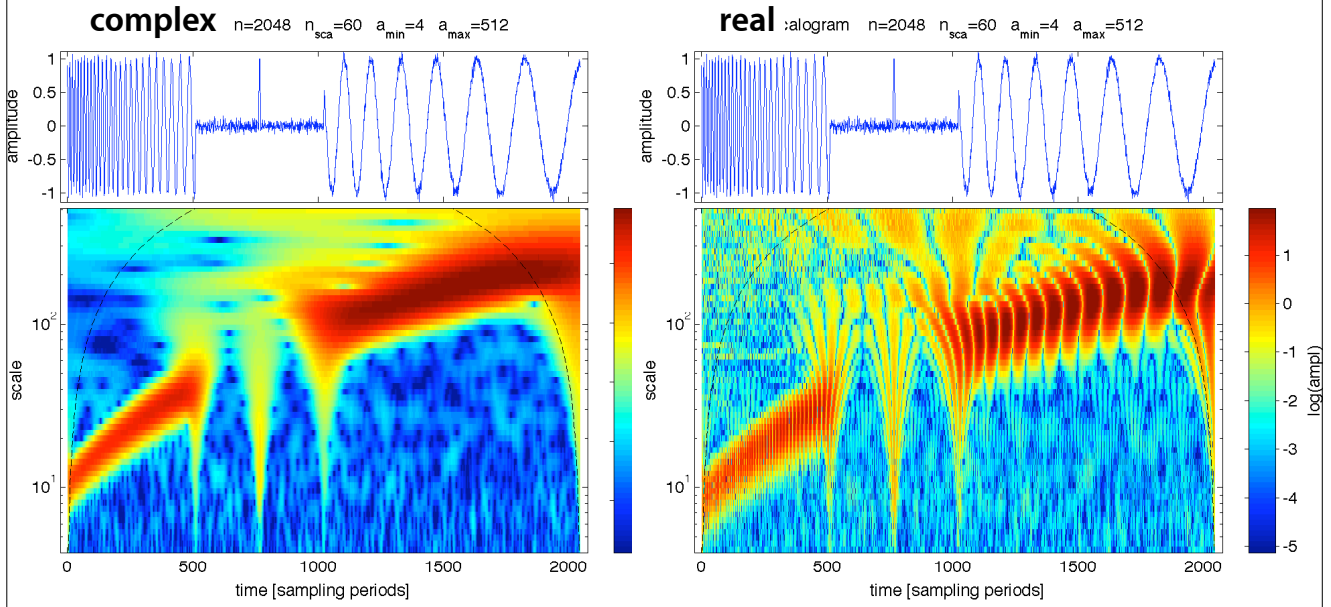


- Since the wavelets are self-similar, the scales should increase logarithmically ( $\neq$  linearly as in Fourier)
- For the discrete wavelet transform, a dyadic grid is imposed

# Complex or real ?



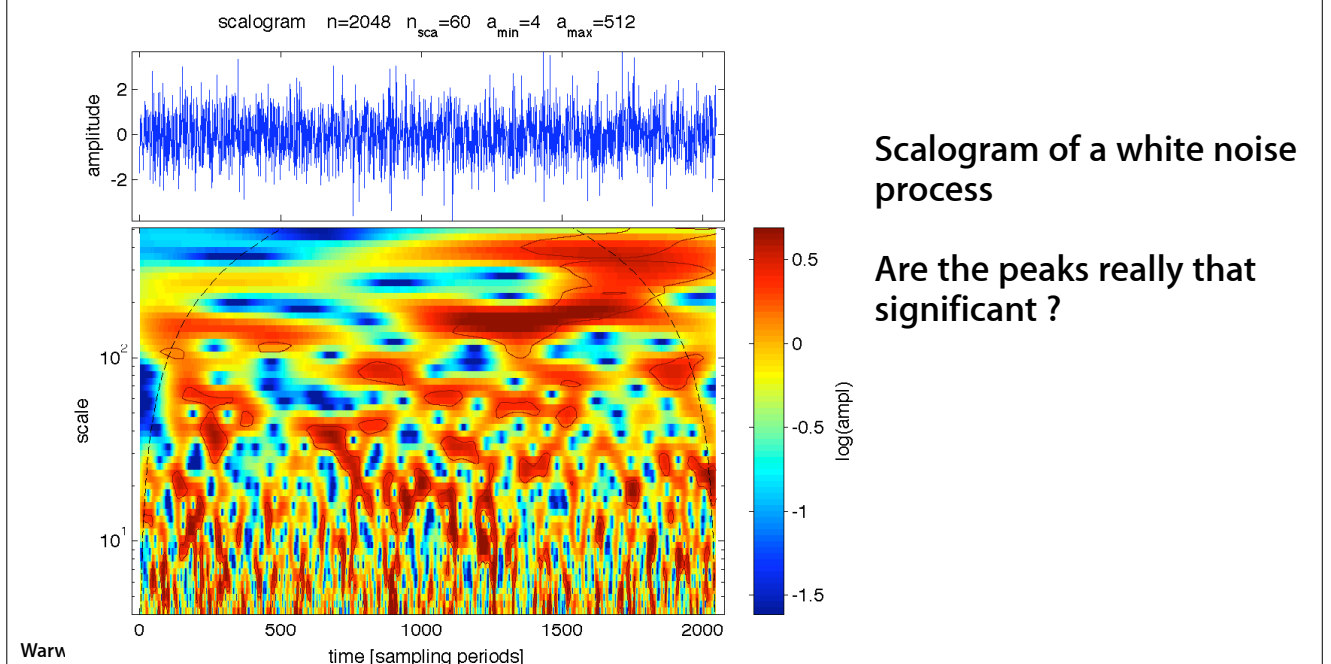
- **Complex wavelets** provide information about phase and amplitude. They are more appropriate for oscillatory behaviour.
- **Real wavelets** are used instead for probing peaks and discontinuities.



# How significant are the wavelet coeffs ?



- Plotting confidence intervals is essential for determining the significance of peaks in the wavelet transform



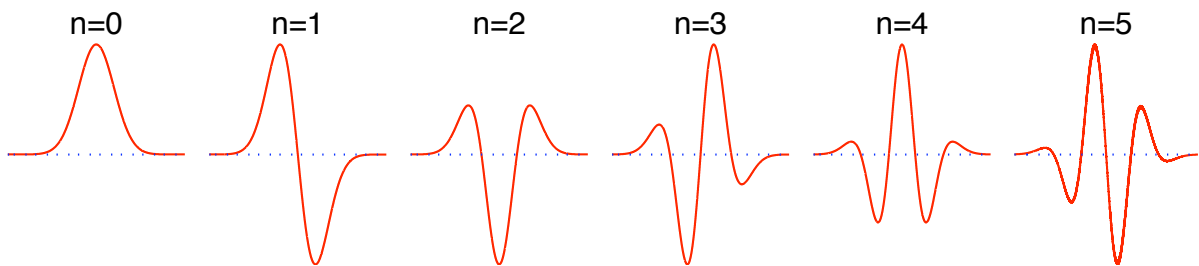


## Vanishing moments

A wavelet has  $n$  vanishing moments if

$$\int_{-\infty}^{+\infty} t^k \psi(t) dt = 0 \quad \text{for } 0 \leq k < n$$

such a wavelet is orthogonal to polynomials of degree  $n-1$



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## Hölder (Lipschitz) regularity

A function  $x(t)$  has a local Hölder exponent  $h$  if

$$|x(t + \tau) - x(t)| \leq K|\tau|^{h(t)}, \quad \tau \rightarrow 0$$

Then the wavelet transform (with  $n$ 'th order wavelets) will satisfy

$$|X(\tau, a)| \leq \begin{cases} K |a|^{h(t)}, & a \rightarrow 0 \quad \text{if } h(t) \leq n \\ K |a|^n, & a \rightarrow 0 \quad \text{if } h(t) \geq n \end{cases}$$

***If the wavelet order  $n$  is large enough,  
then the wavelet transform reveals local  
regularity***

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## Local regularity



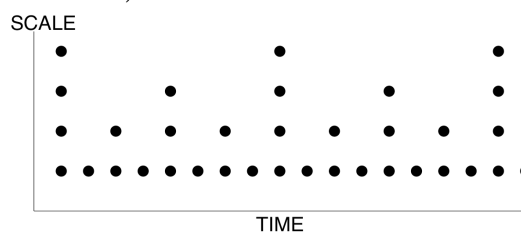
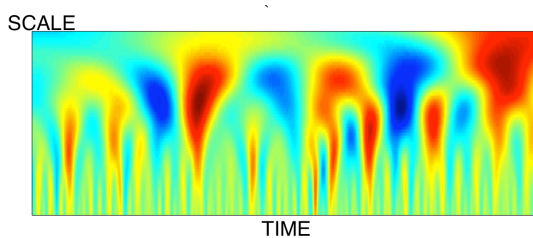
The consequences of the previous property are

- n'th order wavelets are insensitive to trends that can be described by polynomials of order  $< n$
- low order wavelets thus should be avoided, unless the signal has a low regularity

**Adapt you wavelet to the properties of your data set**

- **regularity**
- **symmetry**
- **time-frequency localisation**
- ...

## Discrete or continuous transform ?



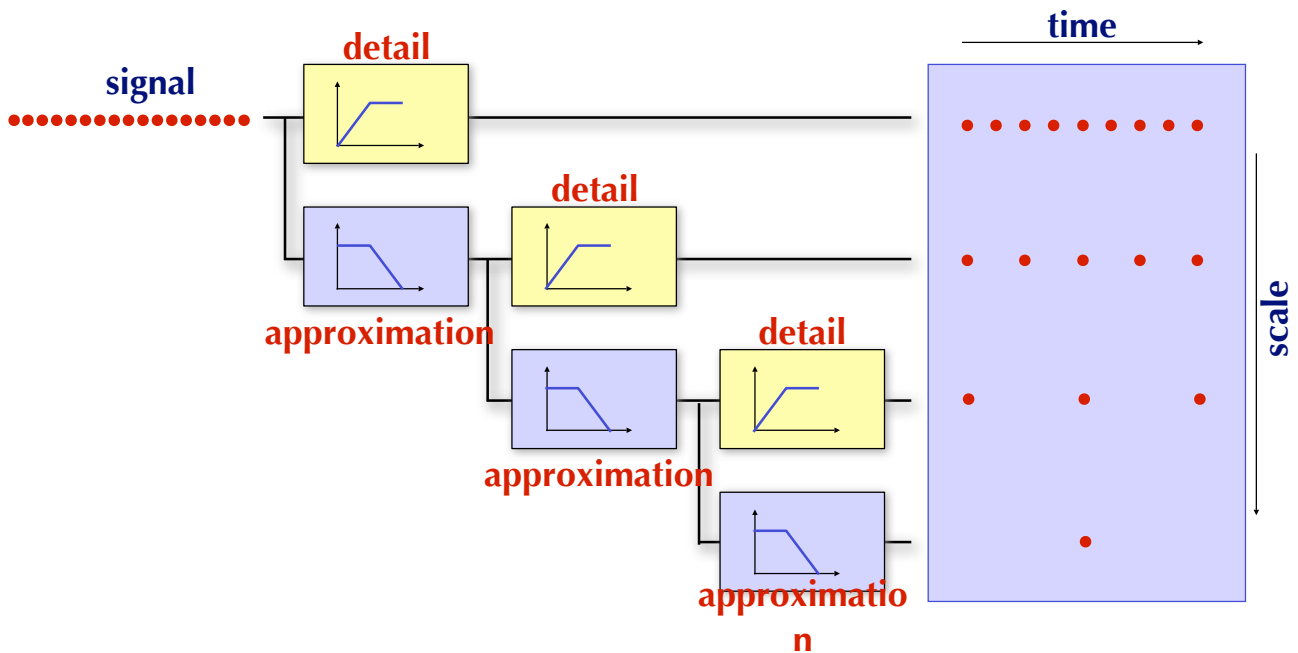
### Continuous transform

- highly redundant
- good for data analysis
- scales can be freely chosen
- computationally expensive

### Discrete transform

- non-redundant and uses orthogonal bases
- useful for multiresolution analysis (denoising) and compression
- scales are imposed
- very fast algorithms (faster than FFT)

## Multiresolution analysis

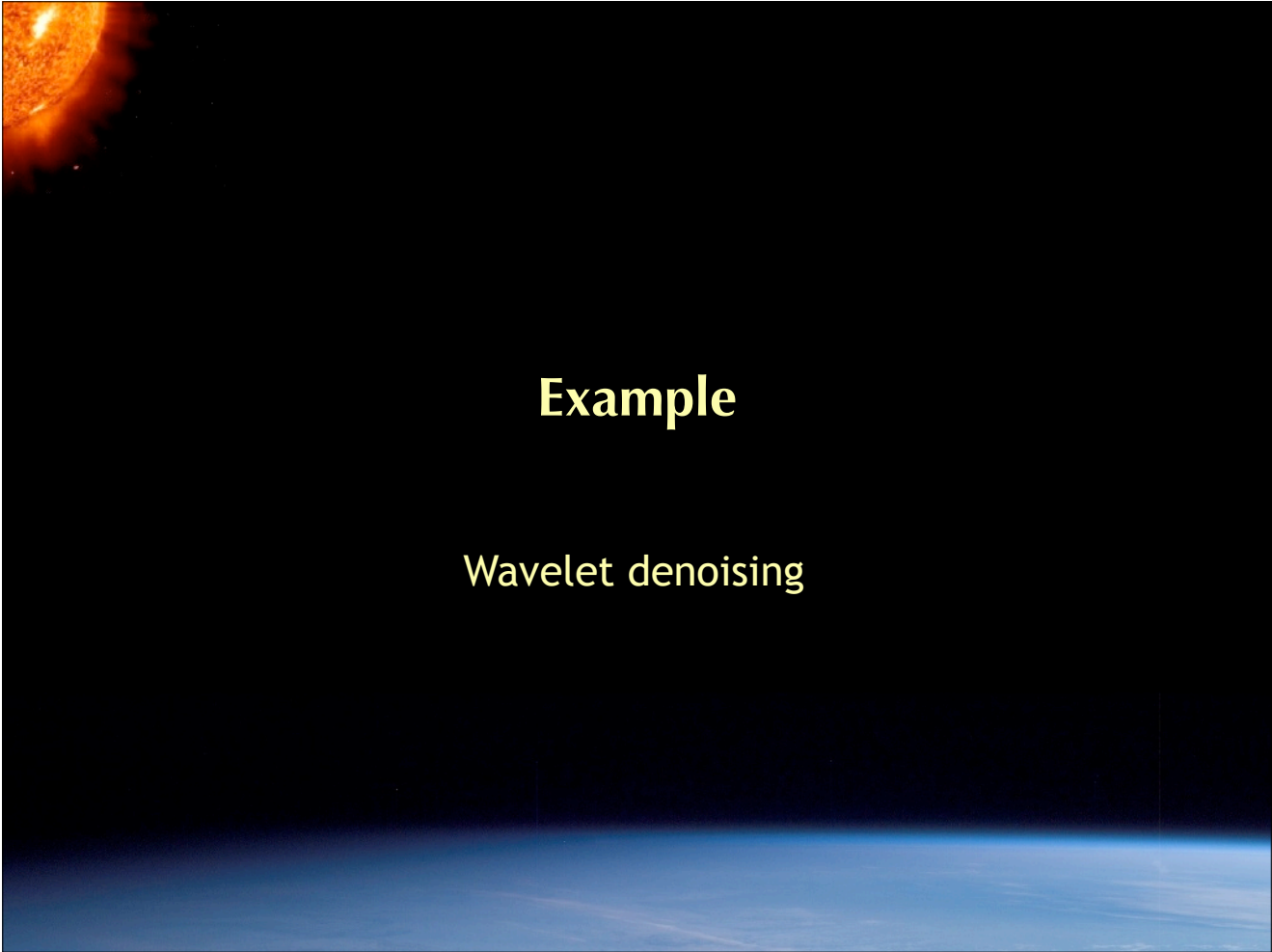


The discrete wavelet transform can be implemented by means of a fast pyramidal recursive filter bank

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## Why use wavelets then ?

- Because wavelets are local in time AND frequency, they're much more efficient than the Fourier transform for capturing transients.
- There have been attempts to use wavelet bases for simulating HD / MHD, with limited success so far.



# Example

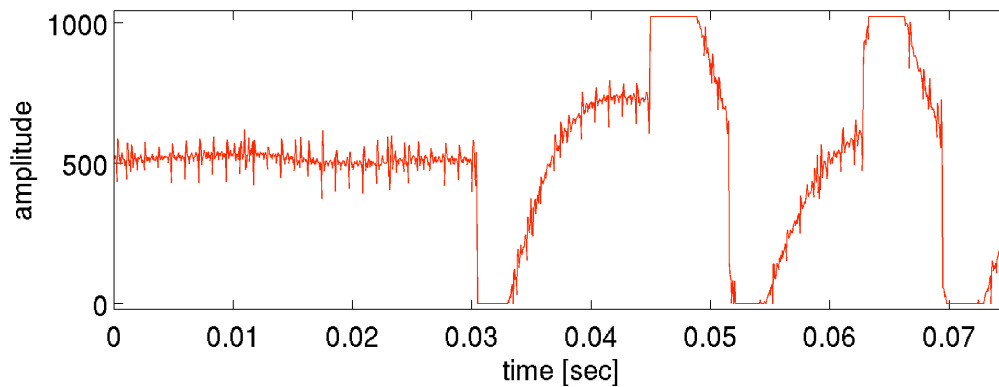
## Wavelet denoising

24

# Example



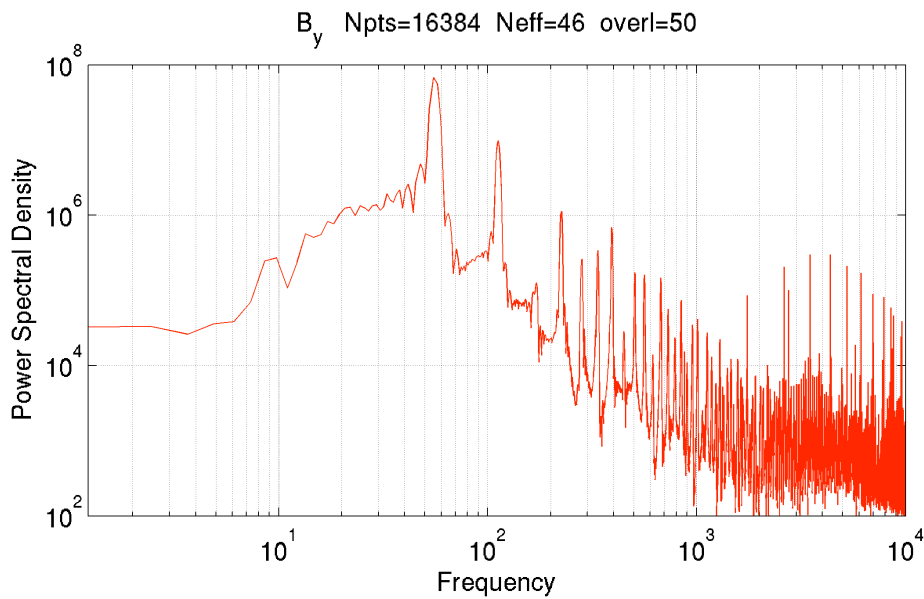
AC magnetic field measurements made by the a sounding rocket in the terrestrial cusp region



**The search coil data are heavily polluted by interference noise from nearby instruments. Is this hopeless ?**



# Spectral analysis



**Interference noise occurs at all scales, and is non-stationary**  
**→ there is no way we can filter it out by Fourier analysis**

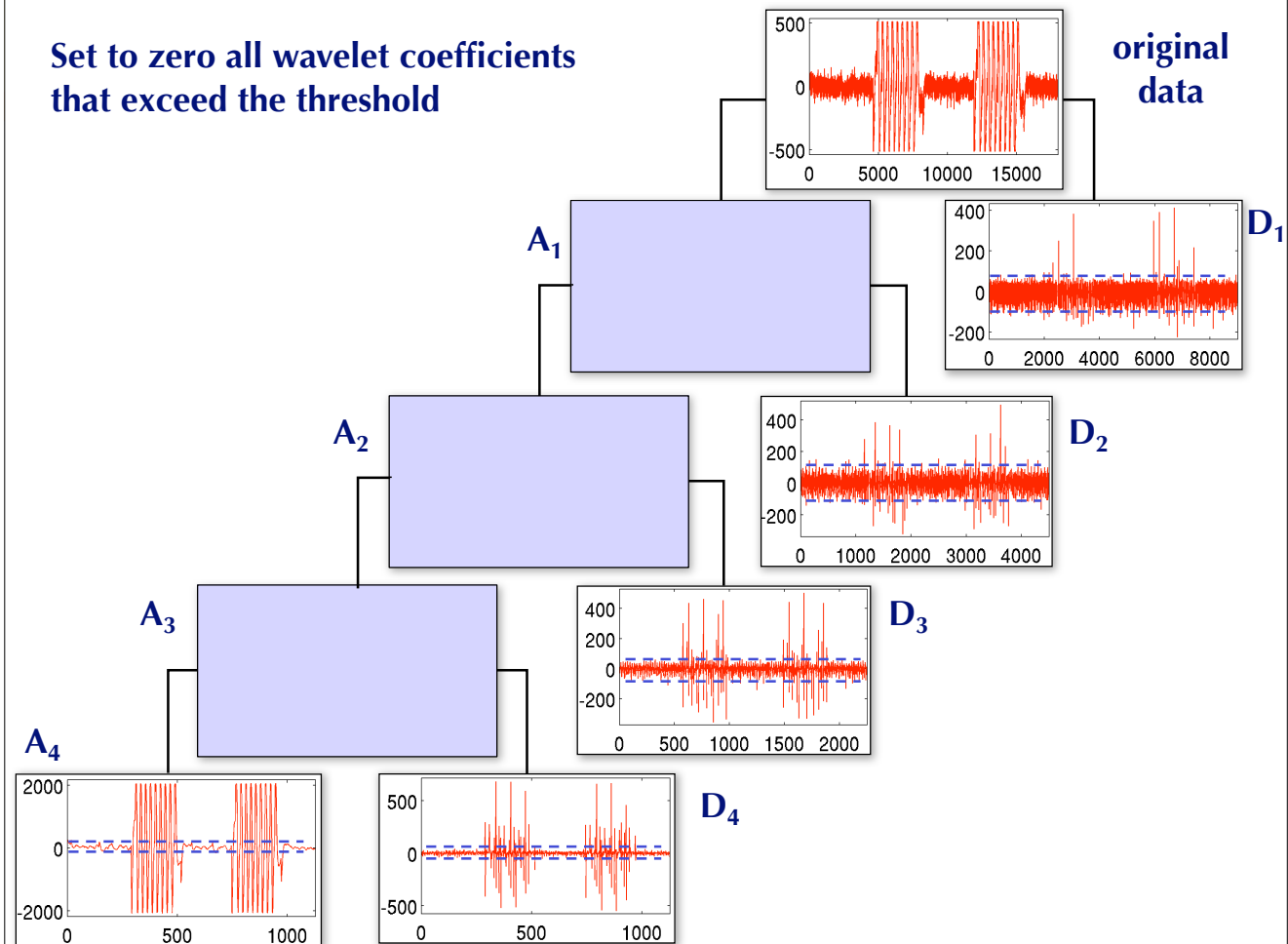
# Discrete wavelet analysis



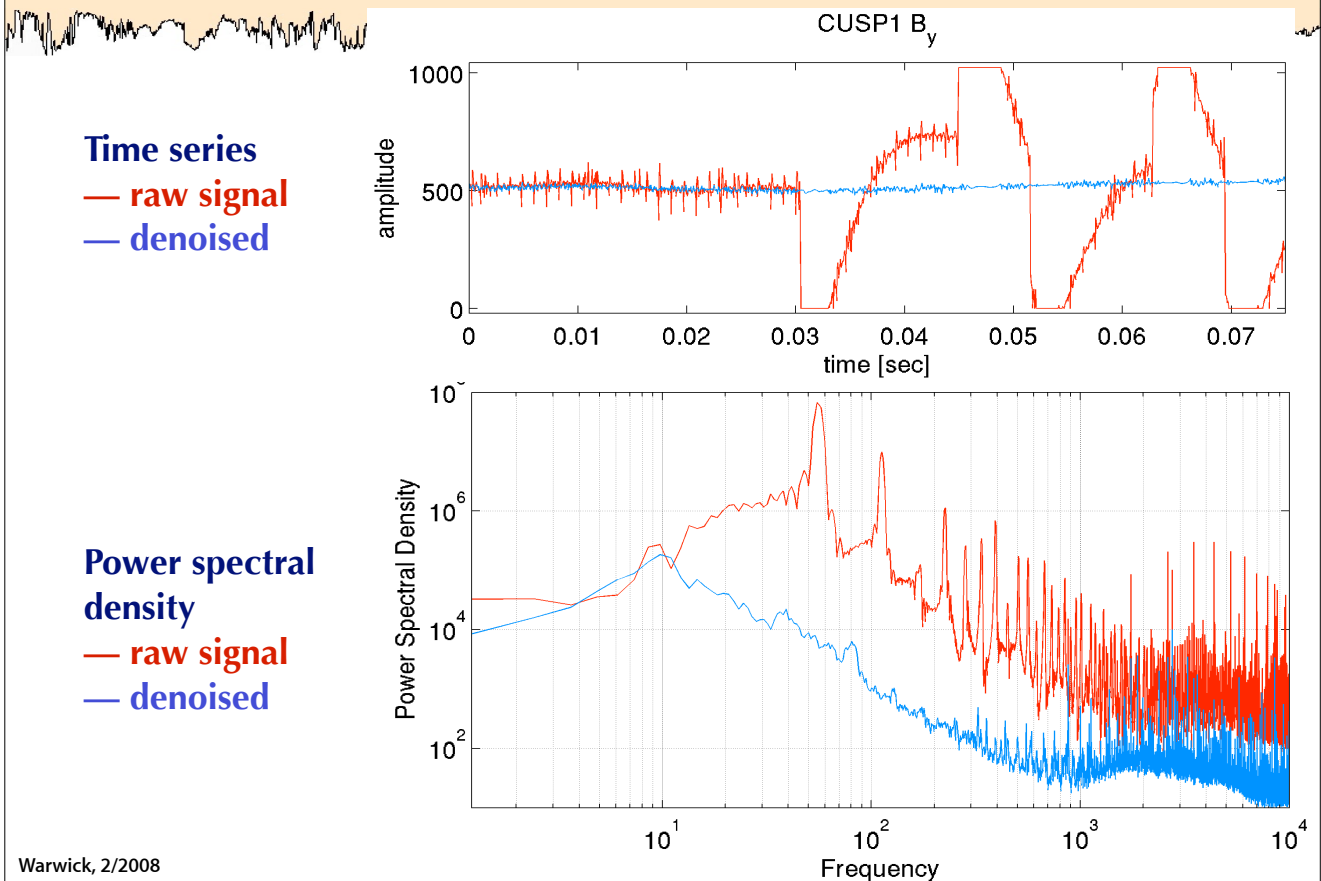
**Recipe** for denoising with the discrete wavelet transform

- Decompose the signal at different levels (**A**pproximations and **D**etails)
- At each level, compare the wavelet coefficients for noisy and quiet periods
- Set to 0 all coefficients corresponding to “noise”
- Reconstruct the data

Set to zero all wavelet coefficients that exceed the threshold



## Results



## Practical issues

- Which wavelet should I use ?
  - Not so important, but symmetry considerations may play a role.
  - Much more important is the choice of the wavelet order :
    - high order = smoother + larger support
    - low order = more discontinuous + smaller support
- How many levels ?
  - More levels = large scales are investigated more in detail
- Where should I set the threshold ?
  - Depends a lot on the type of « noise »
  - Automatic thresholding criteria have been proposed, but selection must be driven by knowledge of the noise

## Conclusion

- Discrete wavelet transforms are remarkably efficient for removing transient patterns
- Their lack of redundancy makes their interpretation difficult  
→ they're rarely used for analysis purposes

**« Wavelet denoising is like removing weed while saving daisies »  
(M. Wickerhauser)**



## Example

Spectral analysis of ion motion in a turbulent magnetic field

→ an example of non-stationarity

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## Simulation model



- Test particle simulations based on a 2-D model

(Hada and Kuramitsu, 1999)

$$\begin{pmatrix} B_y \\ B_z \end{pmatrix} = \sum_{k,l} B_k^l \begin{pmatrix} \cos \\ \sin \end{pmatrix} (kx - \omega_k^l t + \phi_k^l)$$

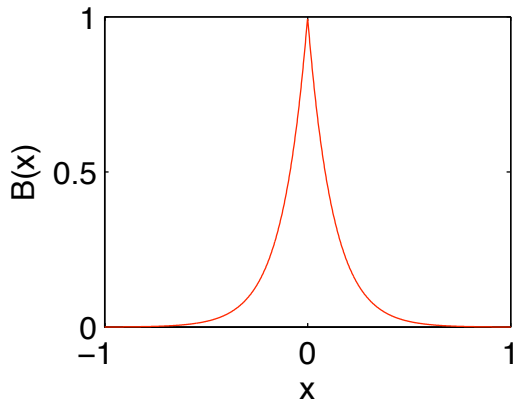
- Phase correlation of adjacent modes can be adjusted by modifying the phases
- Wavefield is periodic in L (L=512), wavefield spectrum goes as

$$B_k \propto k^{-2}$$

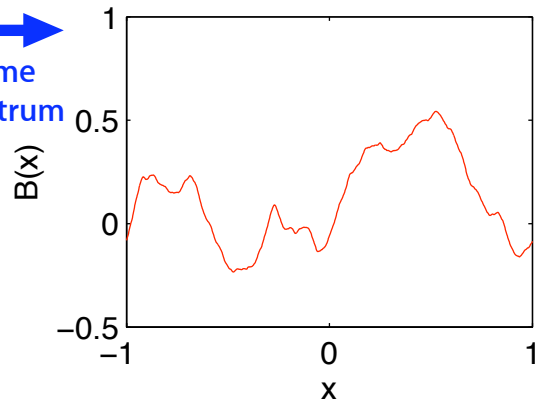
- amplitude and phase coherence are varied



# Simulation model



↕  
same  
spectrum

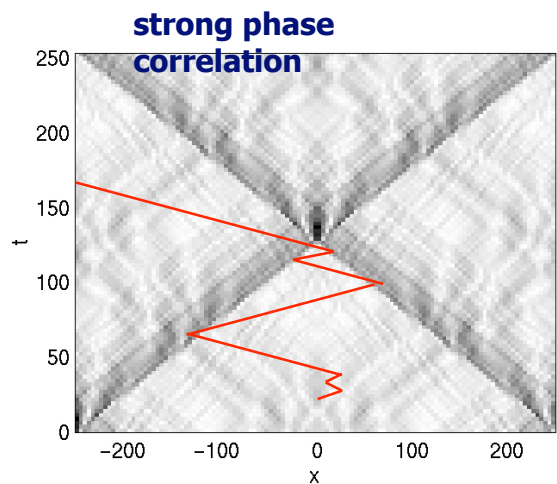
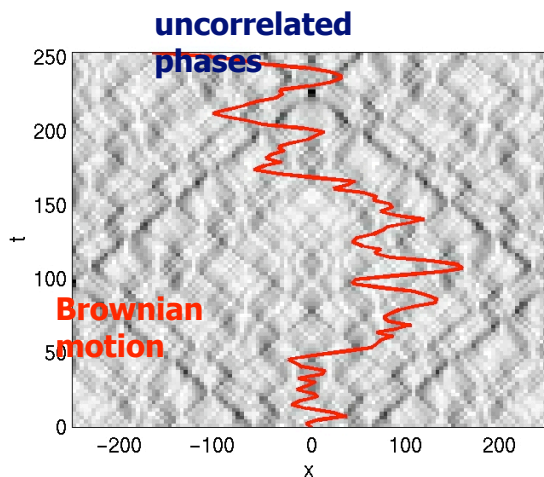


**phase coherence strongly affects ion diffusion through nonlinear trapping by large-amplitude waves**

# Ion trajectories

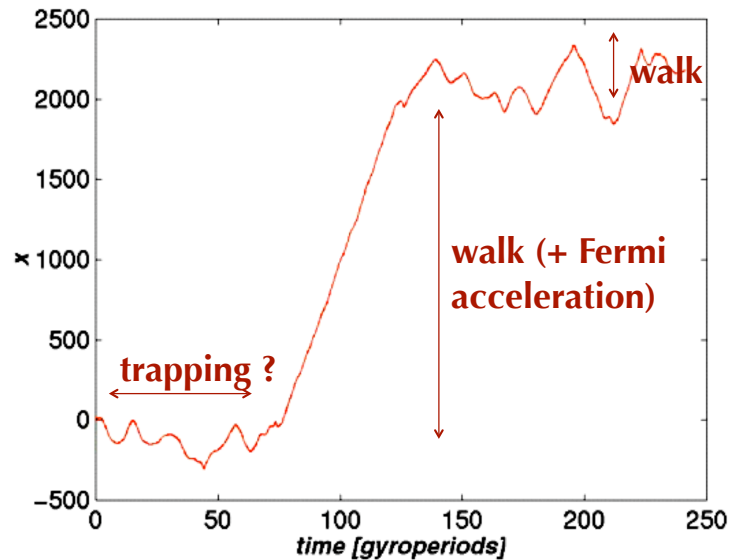


ion trajectories in the x-t plane (grey levels represent  $|B|$ )



## Ion motion

Particle motion is a mixture of **trappings** (gyration) and **walks** (reflection by coherent structures)



## Spectral analysis

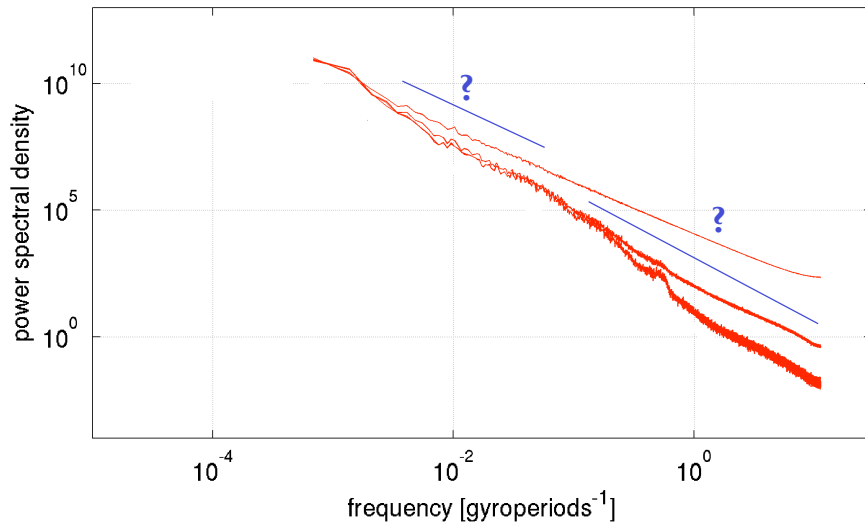
- We're dealing with a self-similar process

$$\Rightarrow u(x) = \lambda^\alpha u(\lambda x)$$

$$\Rightarrow u(k) = \mathcal{F}u(x) \sim k^{-\beta}$$

- Estimate the power spectral density of  $x(t)$  using windowed Fourier transform with various types of windows and detrendings

# Spectral analysis



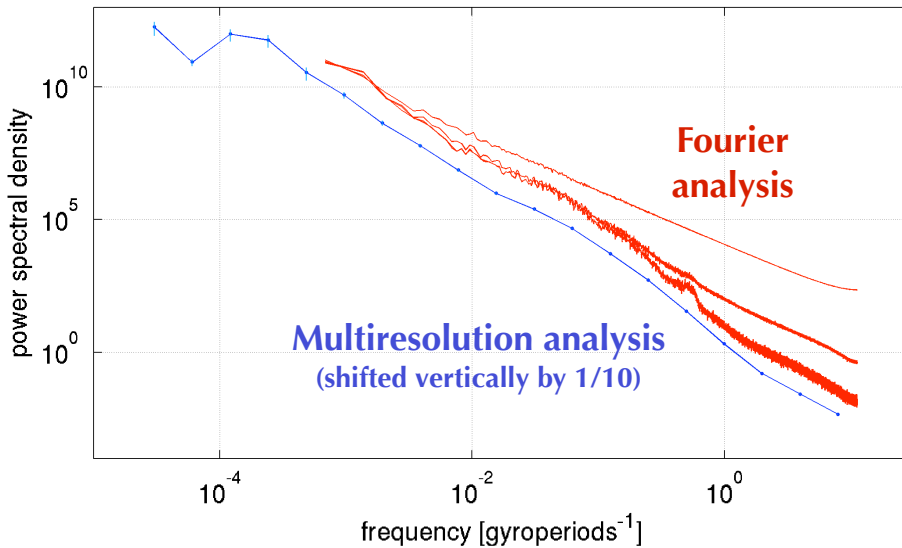
Self-similar ?

What is the correct spectral index ?

# Spectral analysis



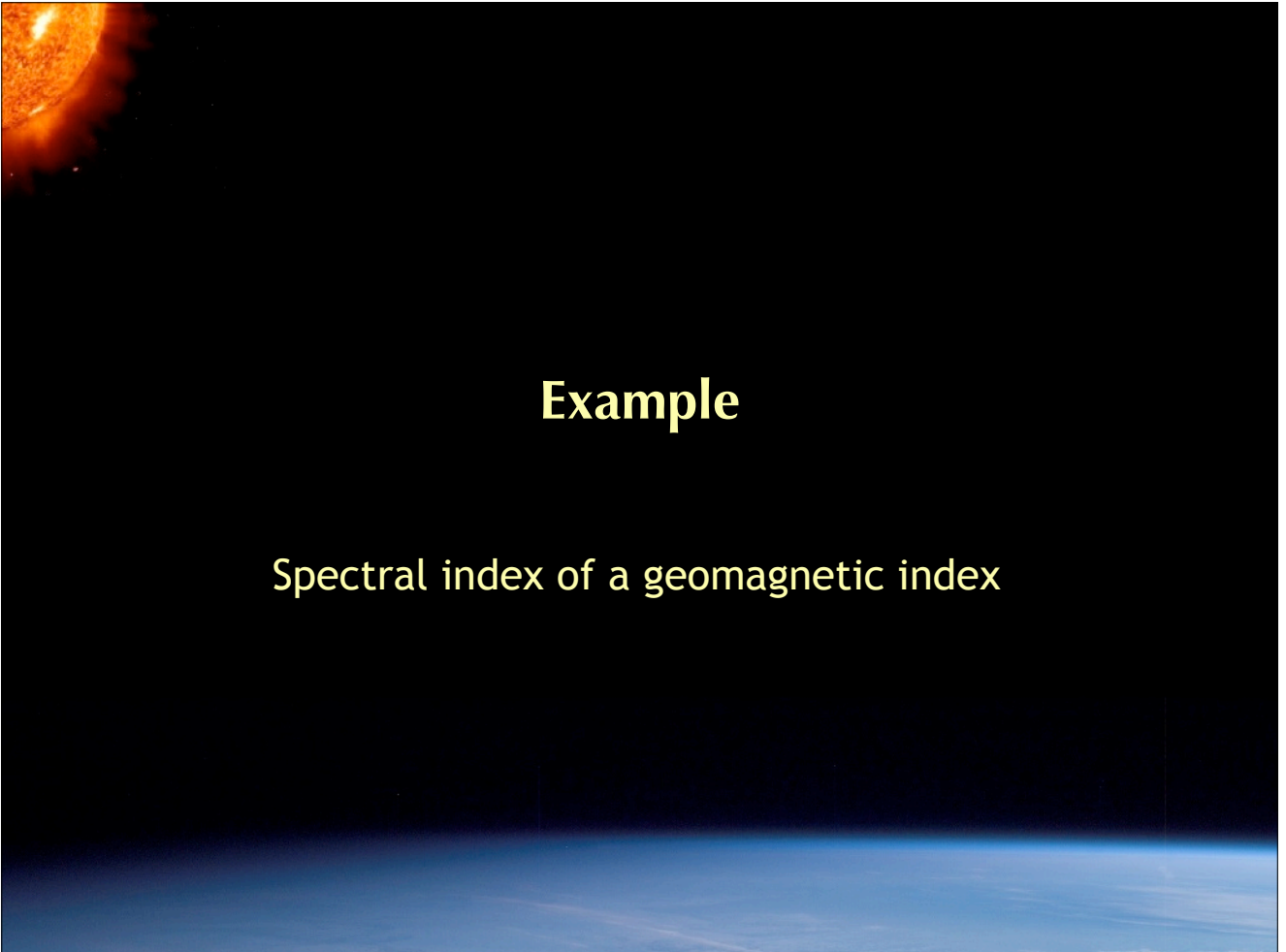
Multiresolution analysis provides a sound and unique estimate of the spectral index



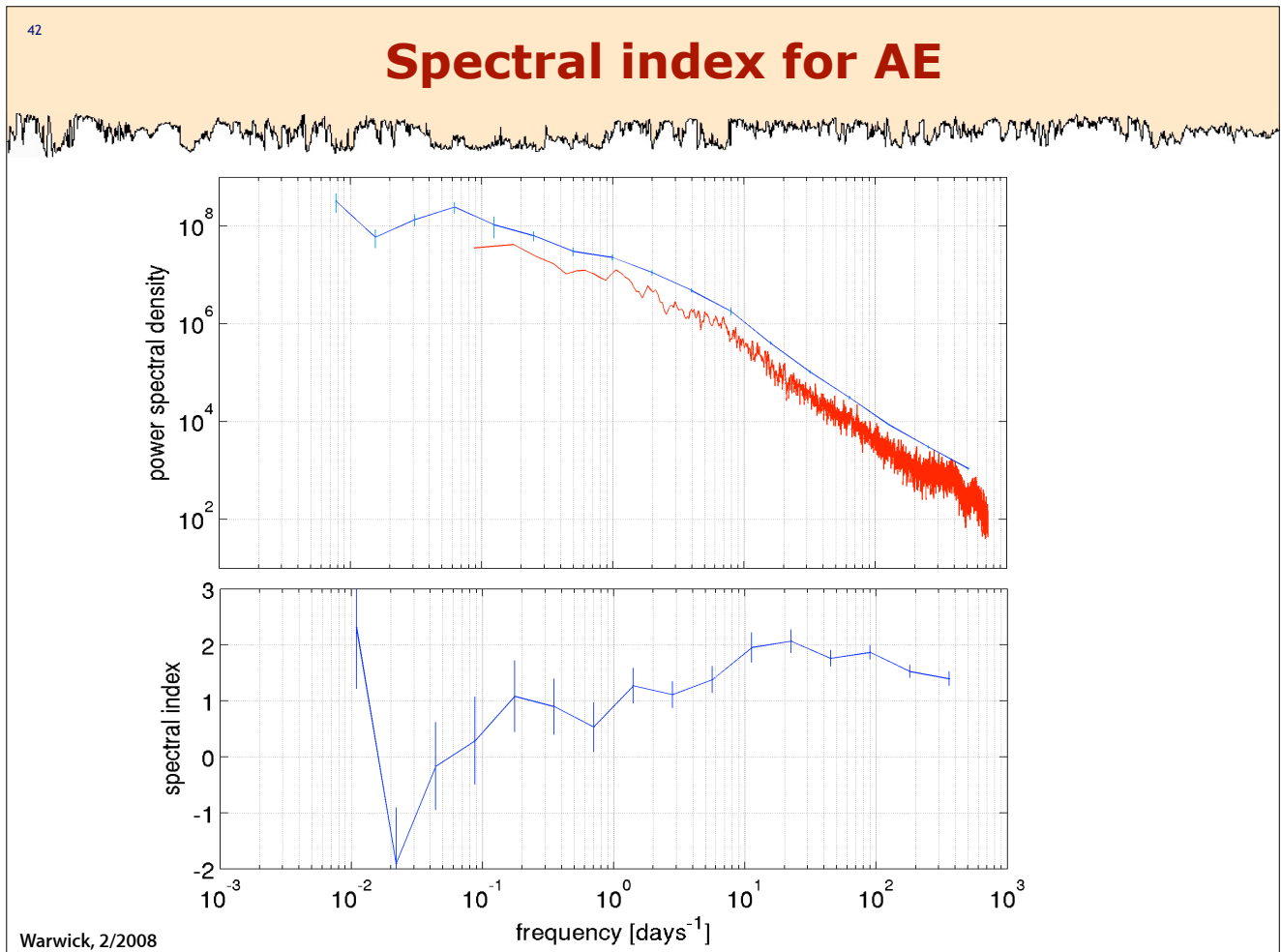
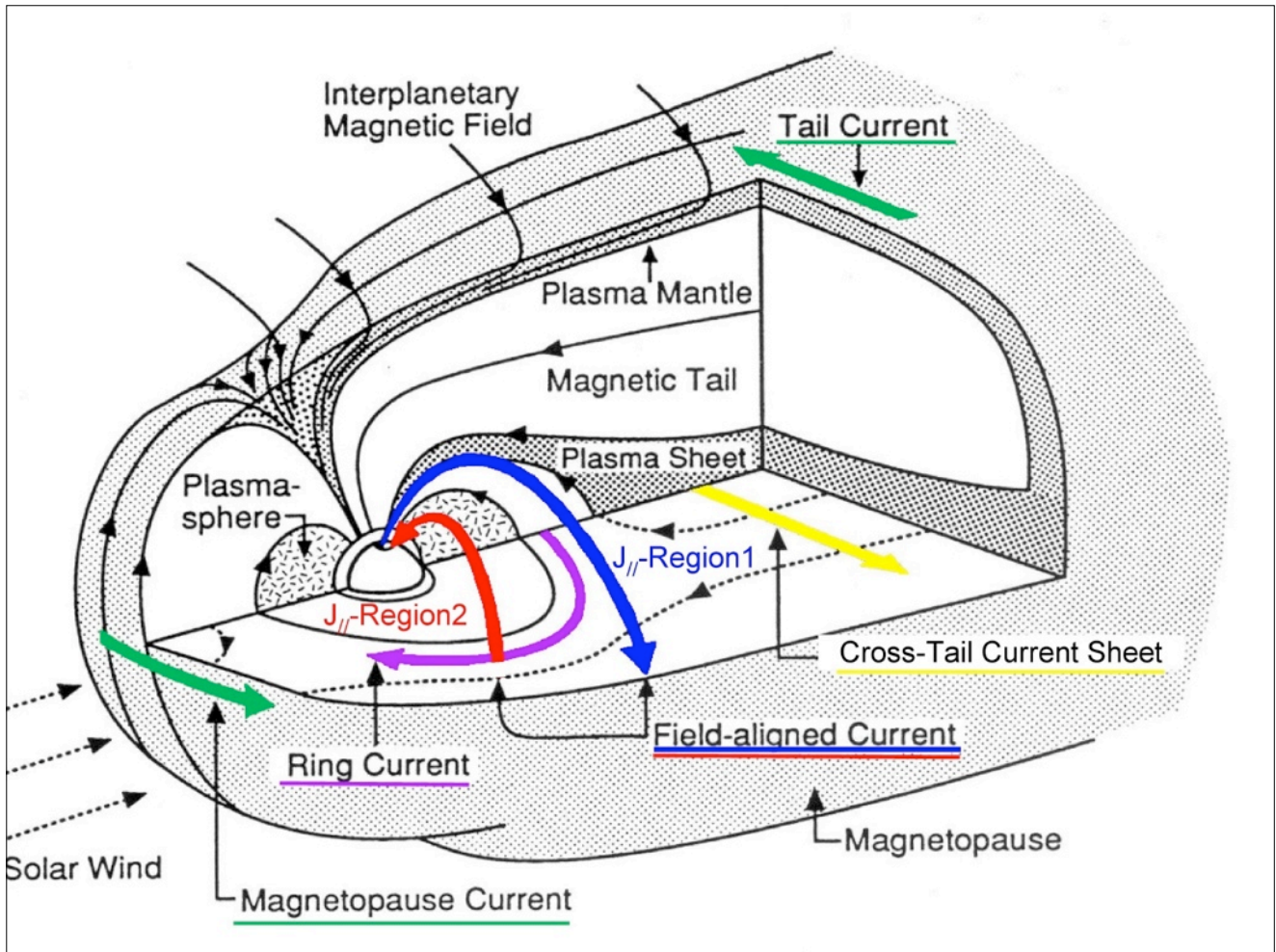
## Conclusion



- Fourier analysis can be **wrong** when the time series is not sufficiently stationary
- Even for stationary data, Fourier analysis gives a biased estimate of the spectral index
  - ➔ multiresolution analysis provides an **unbiased** and **more robust** estimate of the spectral index (Abry et al., 2000)
- But be careful to take wavelets of sufficiently high order







## Characterizing scale invariance

- Scale invariance is a key concept in the statistical analysis of turbulence
- The property of interest is the **interplay** between scales, rather than the role played by each individual scale
  - ➔ spectral indices reveal one aspect only of scale invariance
- To get a more complete picture, other estimators are needed
  - structure functions
  - waiting times
  - singularity spectra
  - ...
- For the particular case of our ion motion, it is interesting to look at the distance travelled by the ion

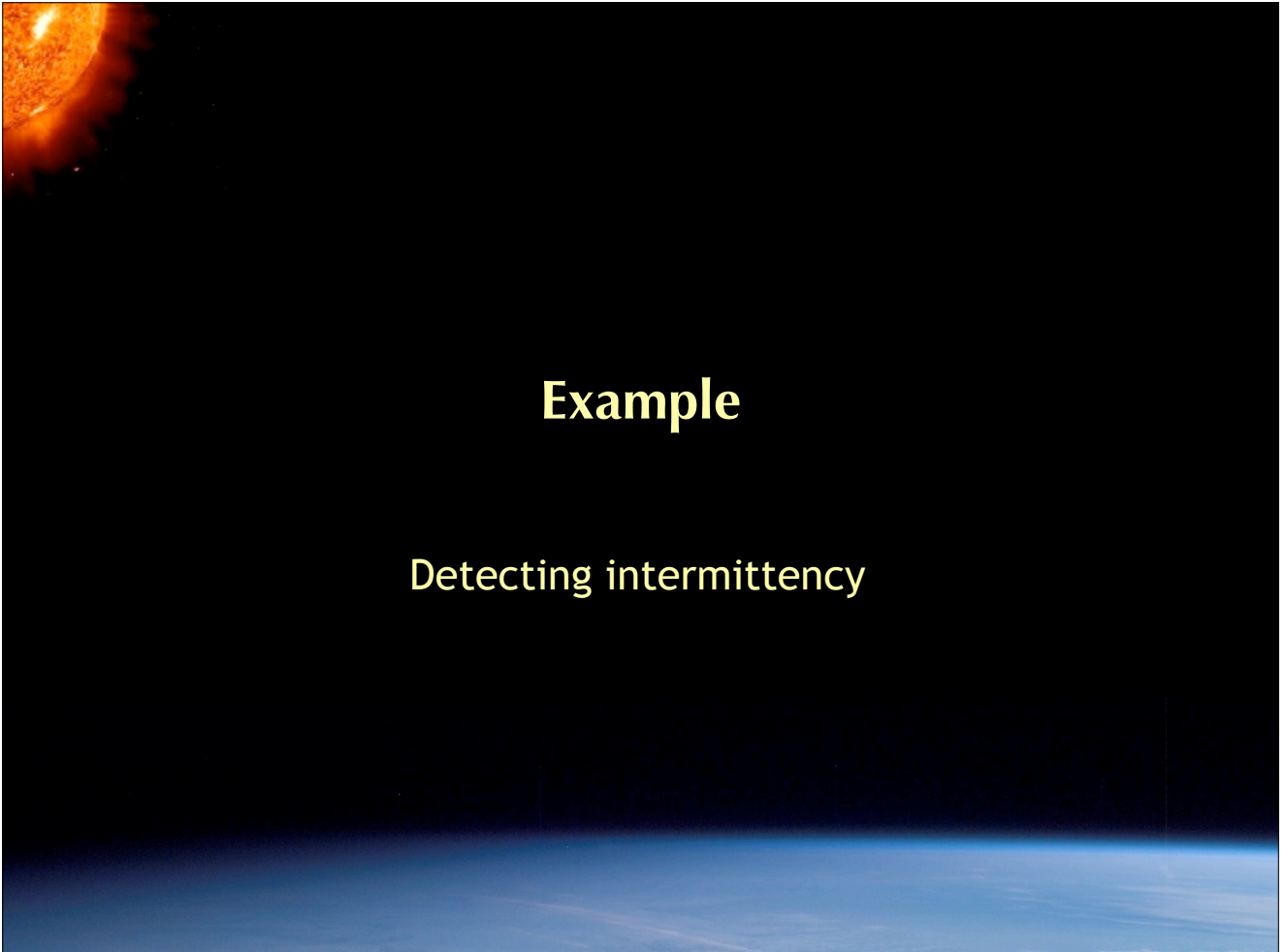
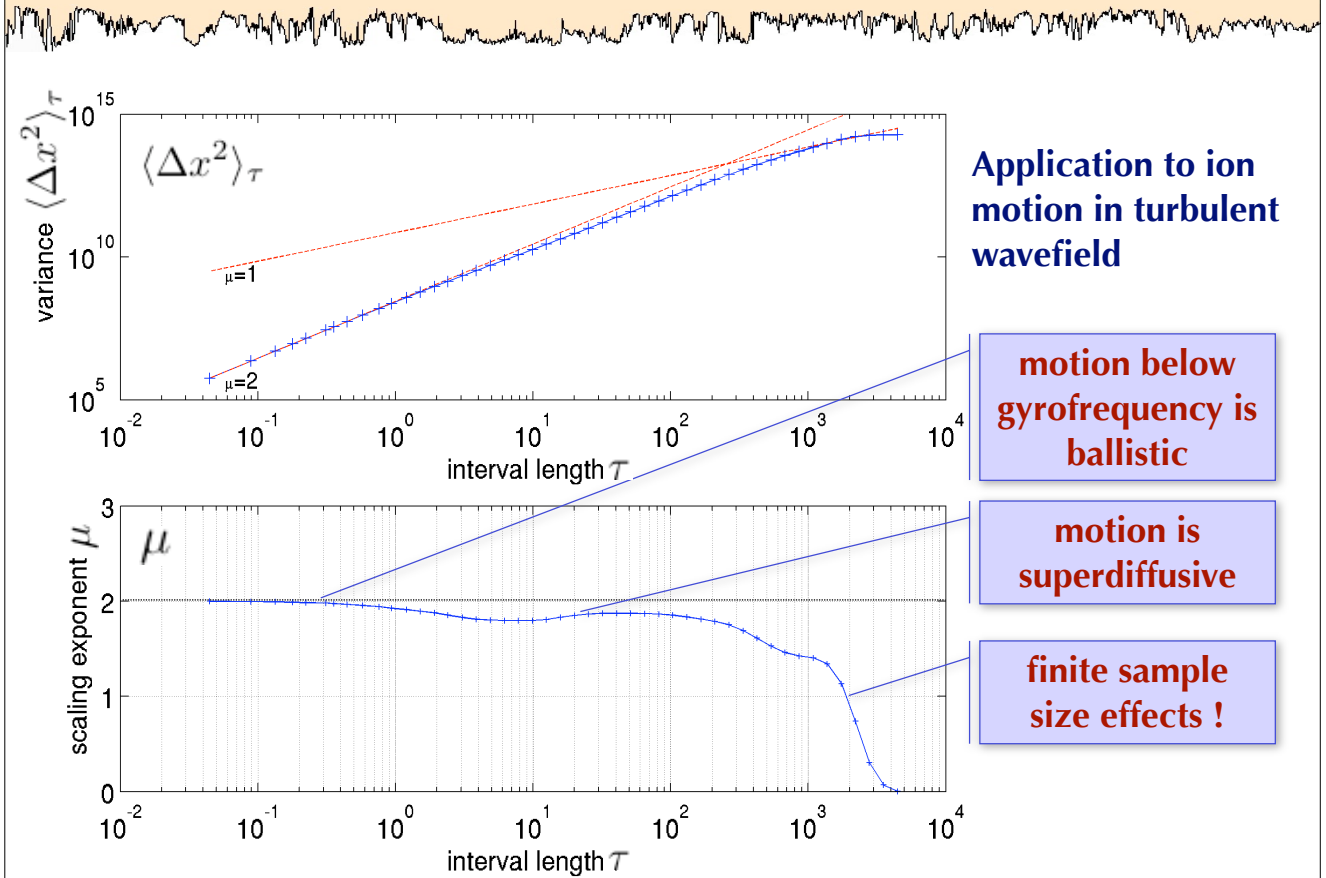
## Ion motion

Let  $\langle \Delta x^2 \rangle_\tau$  be the mean squared distance travelled by an ion after a time  $\tau$

- For diffusion (Brownian motion)  $\langle \Delta x^2 \rangle_\tau = 2 D \tau$
- For convection (ballistic motion)  $\langle \Delta x^2 \rangle_\tau = v^2 \tau^2$
- More generally, we have  $\langle \Delta x^2 \rangle_\tau = \alpha \tau^\mu$

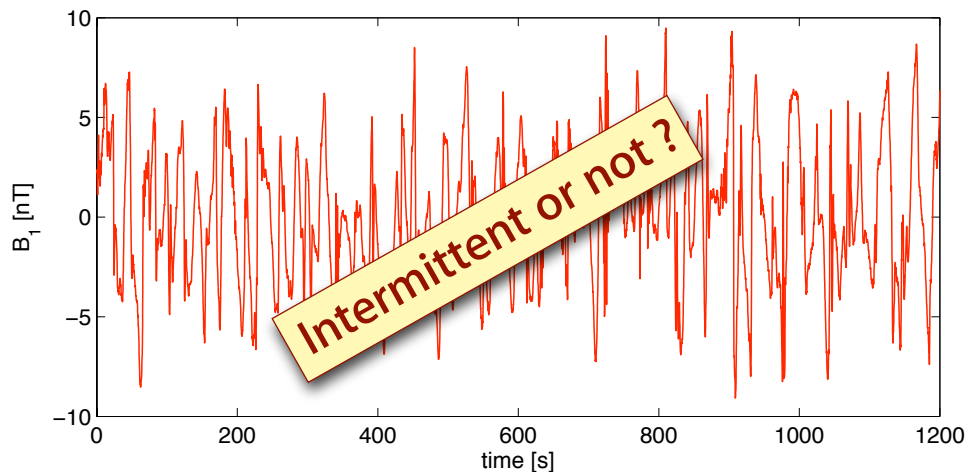
Regimes with  $1 < \mu < 2$  are particularly interesting, for the transport may then be neither diffusive (anomalous transport)

# Ion motion



## Detecting intermittency in flows

- Here **intermittency** = alternation of regimes with normal spectral content, and regimes with significant **excess of energy** in a given range of scales
- **Example** : magnetic field measurements upstream the Earth's bow shock



Warwick, 2/2008

## Intermittency

- To quantify the anomalous energy content, we define the **local intermittency measure** (LIM)

$$\gamma(t, \tau) = \frac{|X(t, \tau)|^2}{\int_{t_{min}}^{t_{max}} |X(t', \tau)|^2 dt'}$$

Any value of LIM that significantly exceeds 1 implies a local excess of energy

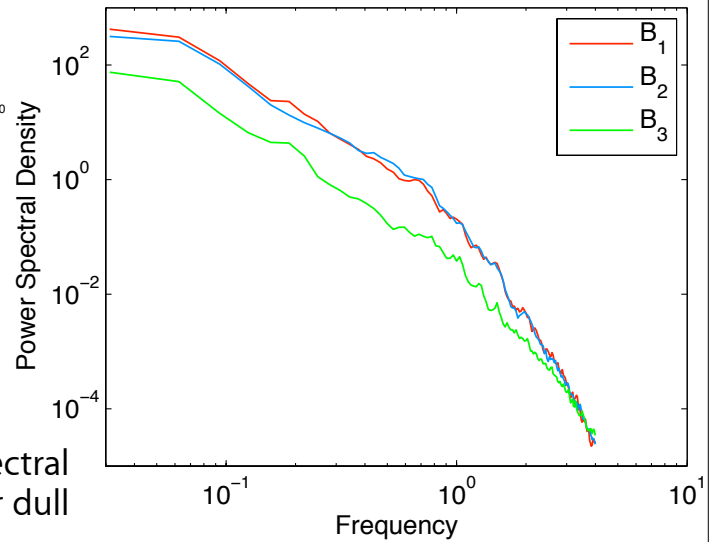
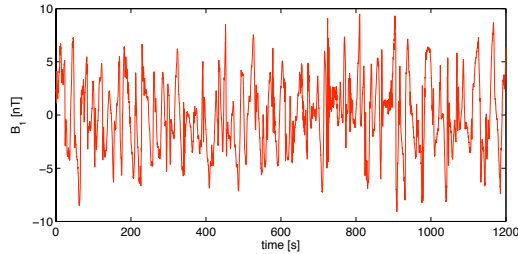
We then need to do a statistical test (chi-square)

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# Intermittency



- **Example from magnetic field at bow shock:** Fourier spectral analysis doesn't reveal anything unusual



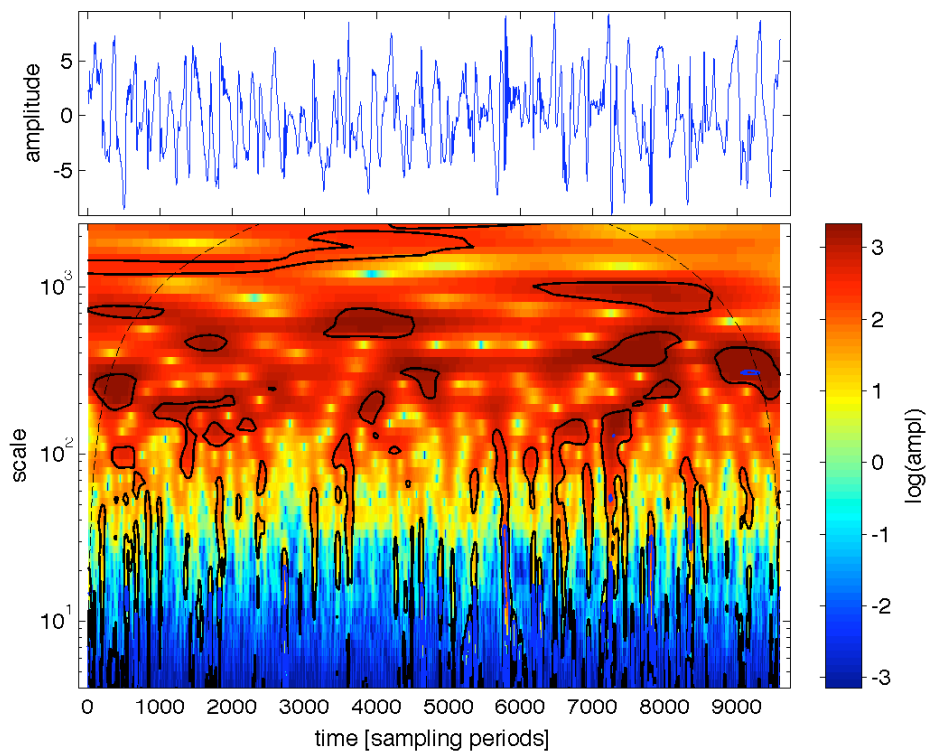
The power spectral density look rather dull

# Intermittency



scalogram  $n=9600$   $n_{sca}=60$   $a_{min}=4$   $a_{max}=2400$

A wavelet scalogram reveals the presence of structures at various scales, but is it really that intermittent ?



# Intermittency

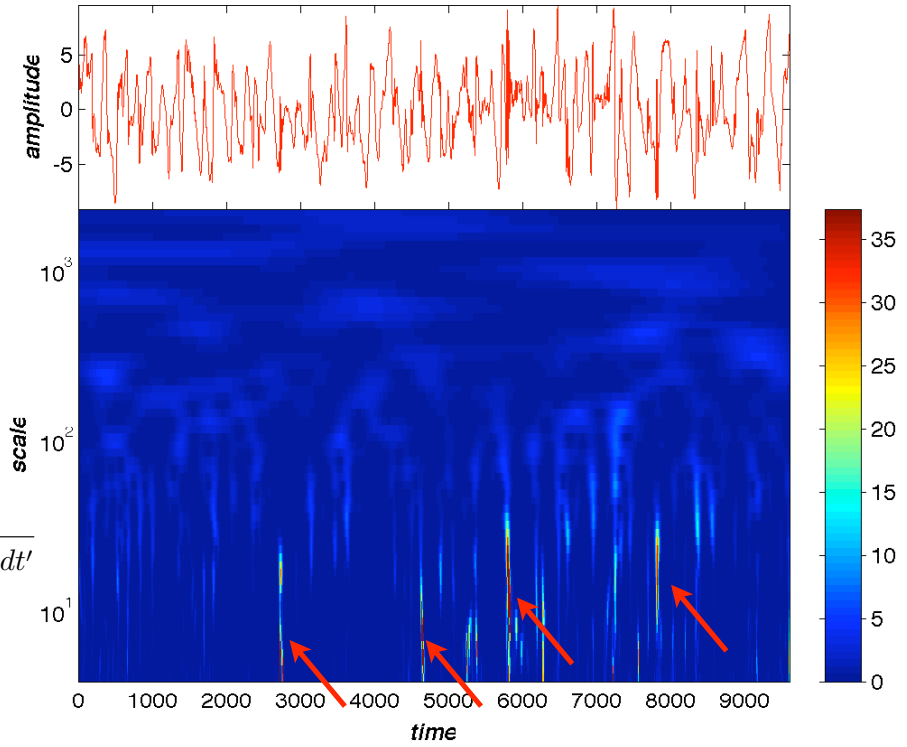


intermittency factor  $n=9600$   $n_{sca}=60$

The local intermittency measure reveals clear bursts

What are these ?

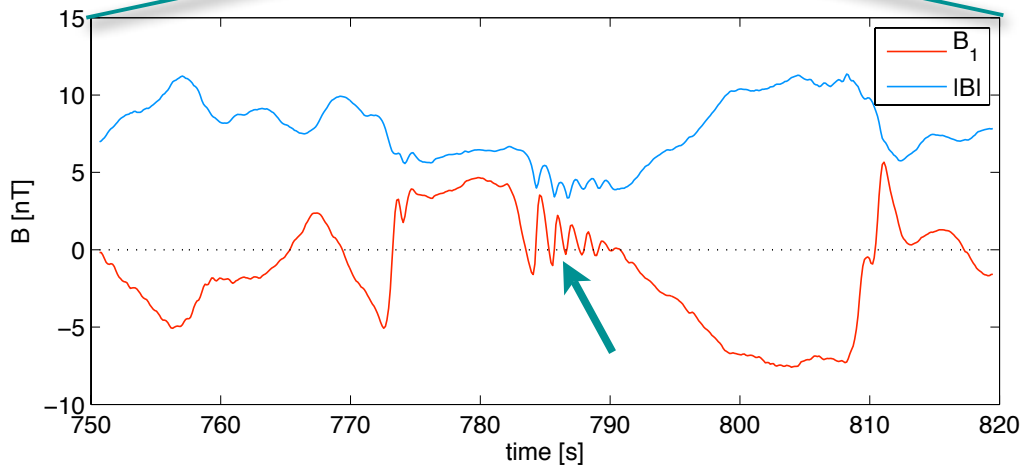
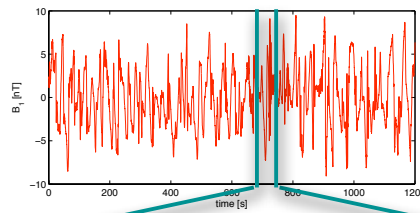
$$\gamma(t, \tau) = \frac{|X(t, \tau)|^2}{\int_{t_{min}}^{t_{max}} |X(t', \tau)|^2 dt'}$$



# Intermittency



- The scalogram reveals the existence of occasional bursts of energy



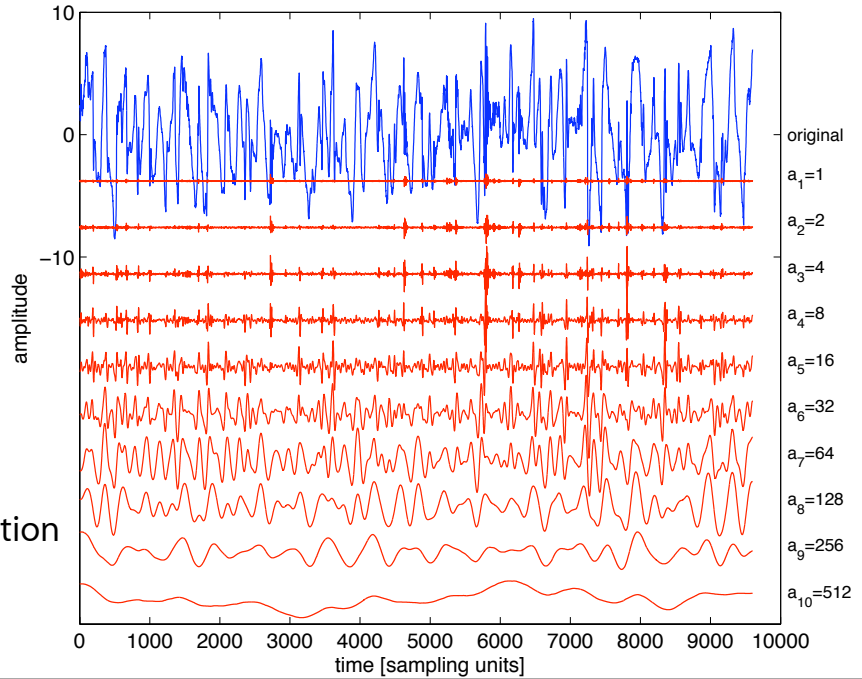


# Intermittency



- This bursty behaviour can be observed both in the continuous and in the discrete wavelet transform

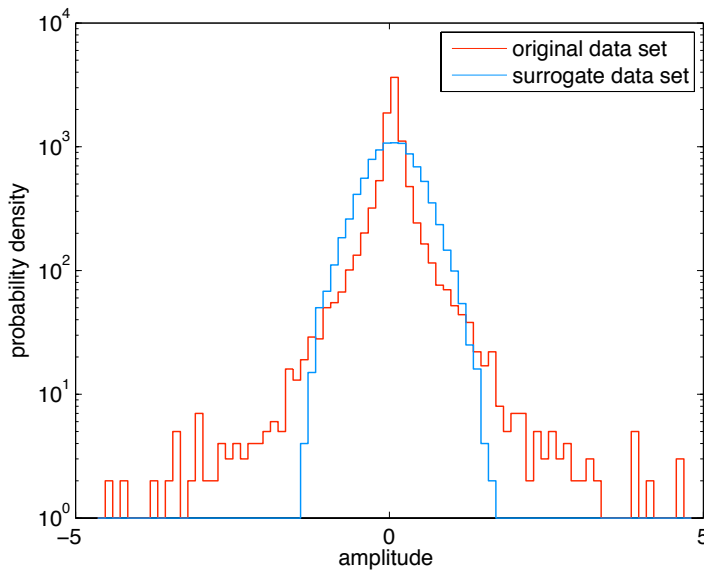
Multiscale decomposition of the magnetic field



# Intermittency



- Since the wavelet transform is linear, many interesting mathematical property remain valid



probability density function of wavelet coefficients (DOG wavelet) at scale=8

# Example

The same concept can be used to detect changes :

monthly mean of eastward component of geomagnetic field (de Michelis et al., 2003)

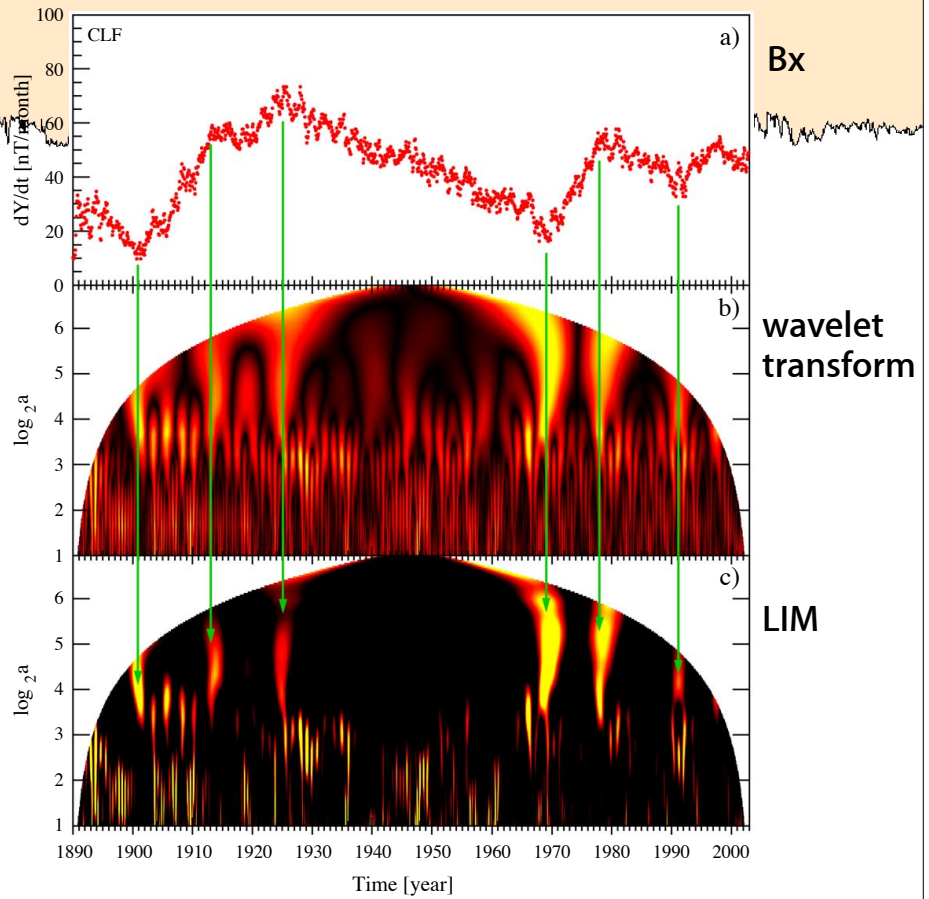
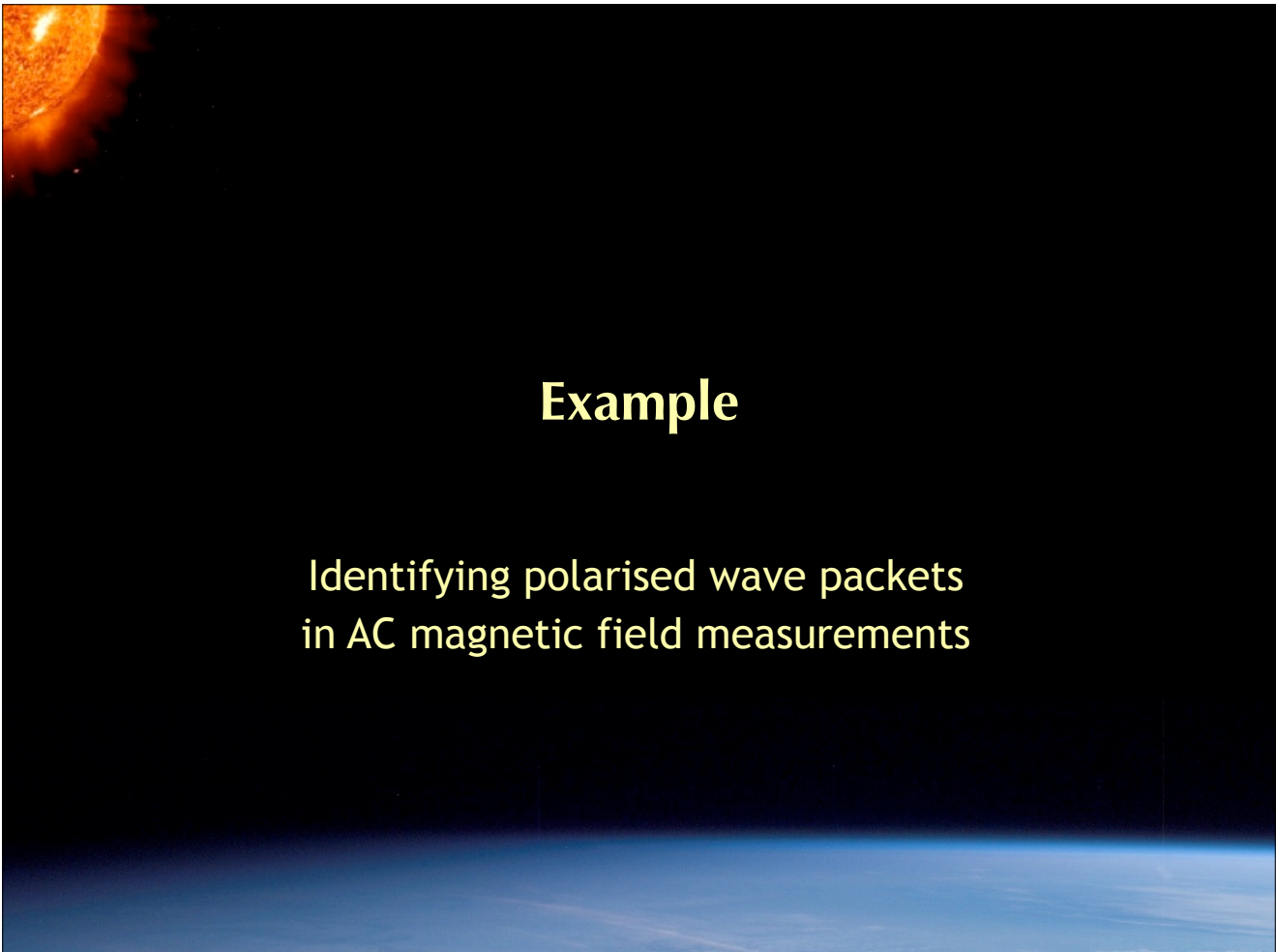


Plate I. Monthly mean values of the geomagnetic secular variation for the eastward component ( $Y$ ) measured at Chambon la Forêt (CLF) observatory from 1890 to 2002 (panel a). Wavelet transform absolute value obtained using Alexandrescu et al. [3] methodology (panel b). Results from LIM analysis using the condition  $LIM_{a,b}^2 > 3$  (panel c).

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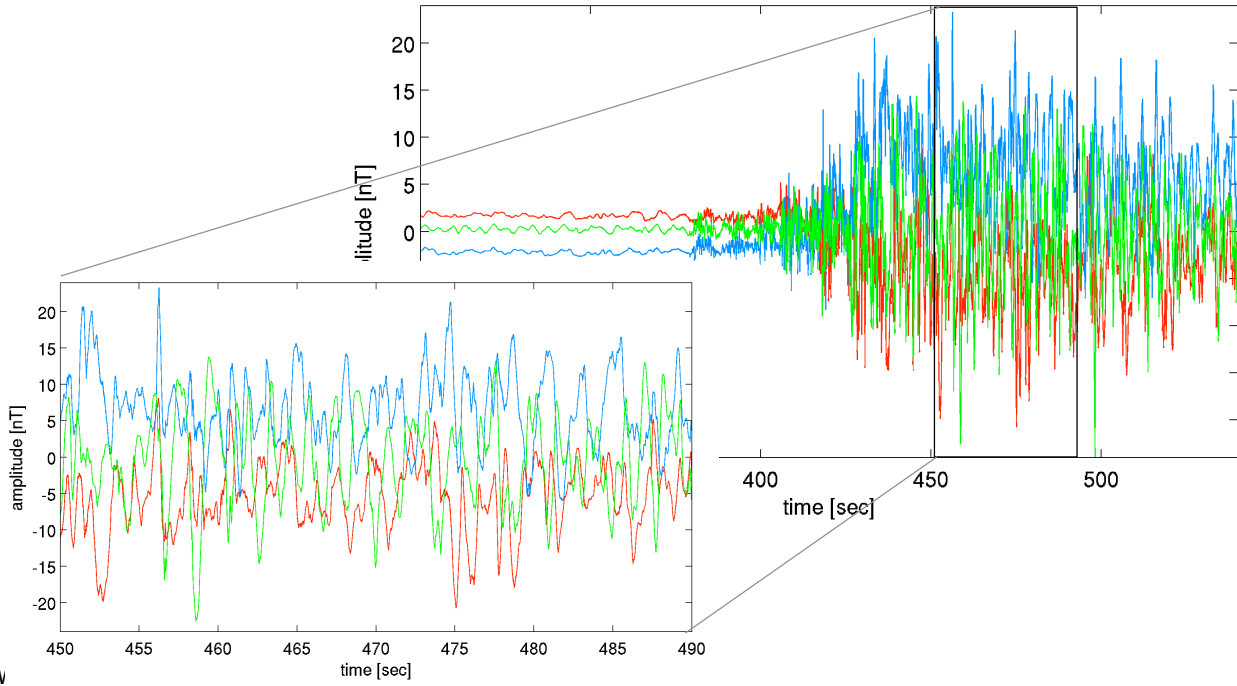
# Example

Identifying polarised wave packets in AC magnetic field measurements

# Magnetic field fluctuations

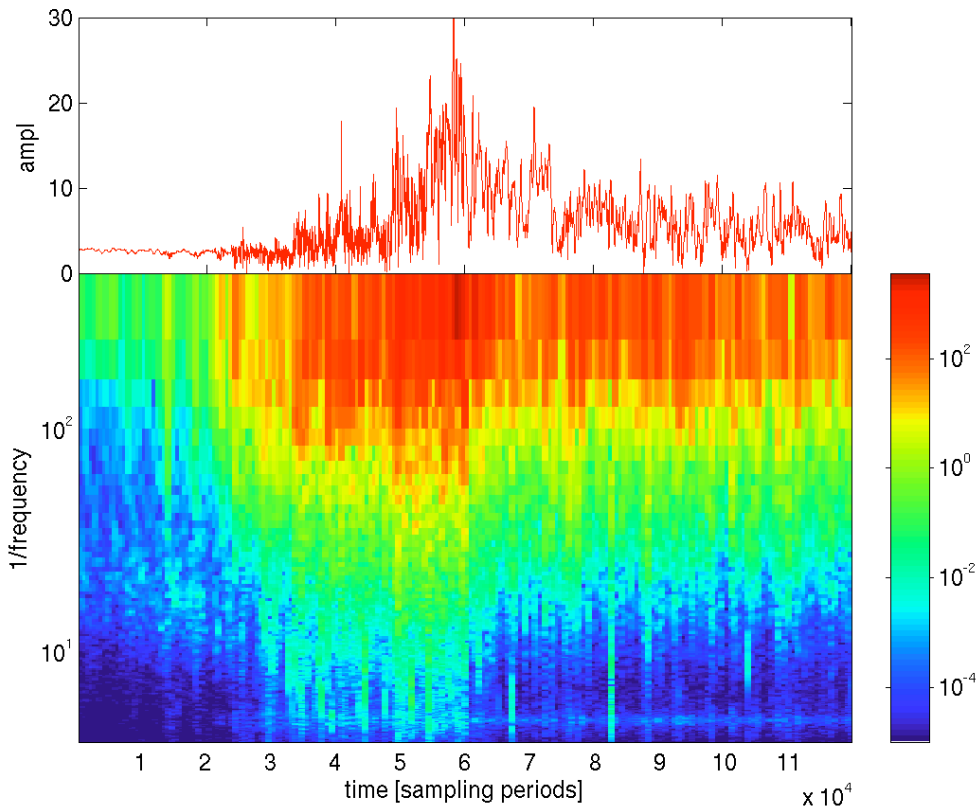


● AC magnetic field measurements at a quasiperpendicular bow shock (STAFF instrument)



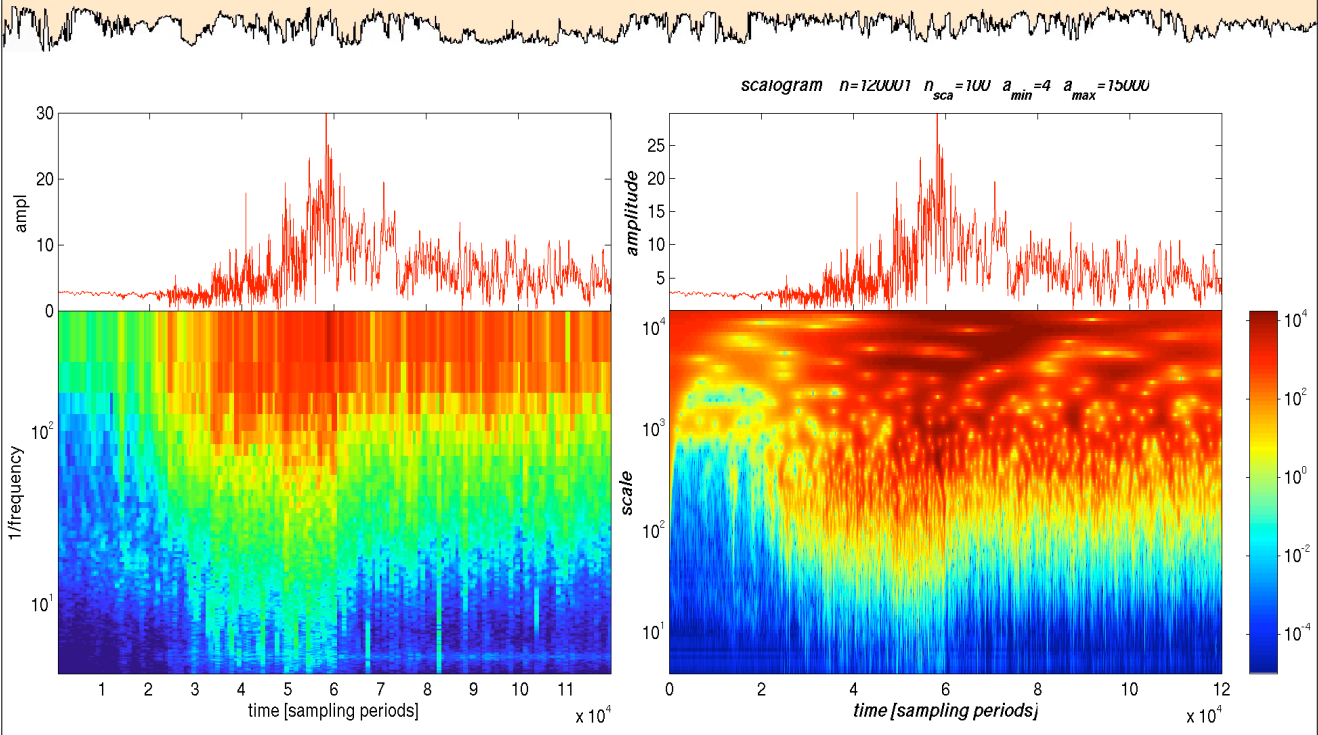
v

# STAFF spectrogram



Warwick

# Spectrogram vs scalogram



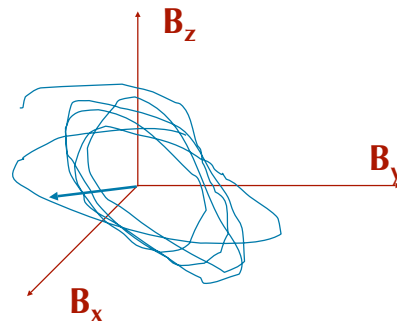
spectrogram (Fourier)

scalogram (Morlet wavelets)

there is nothing really special about the

# Polarisation

- Since spectral characteristics alone are not very appealing here, we look for the polarisation vs time and frequency



- Compute

$$S_{i,j} = \langle B_i(\omega)^* B_j(\omega) \rangle \quad \text{for } i, j = x, y, z$$

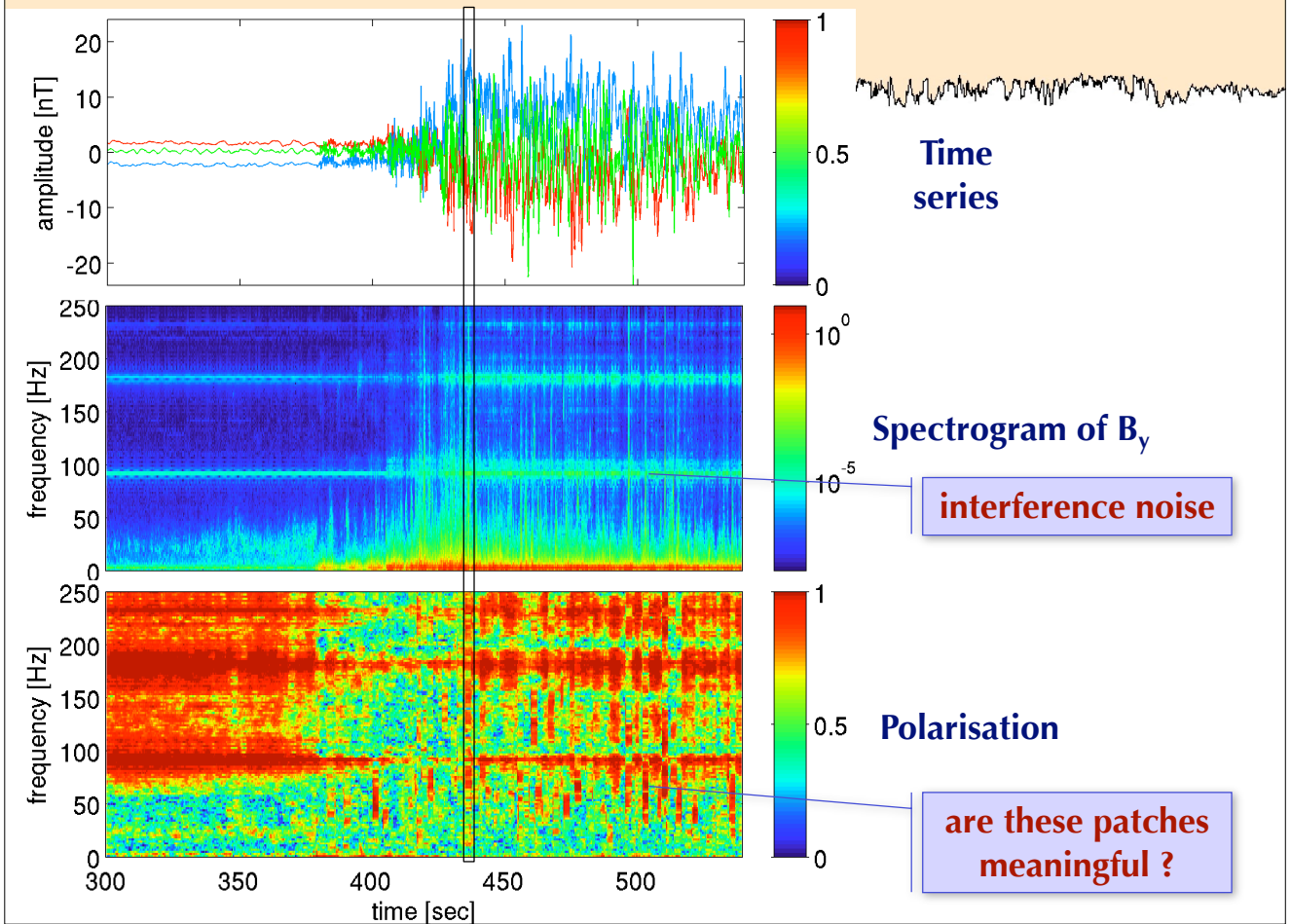
- Diagonalise

$$S(\omega) = \begin{pmatrix} S_{xx} & S_{xy} & S_{xz} \\ S_{yx} & S_{yy} & S_{yz} \\ S_{zx} & S_{zy} & S_{zz} \end{pmatrix}$$

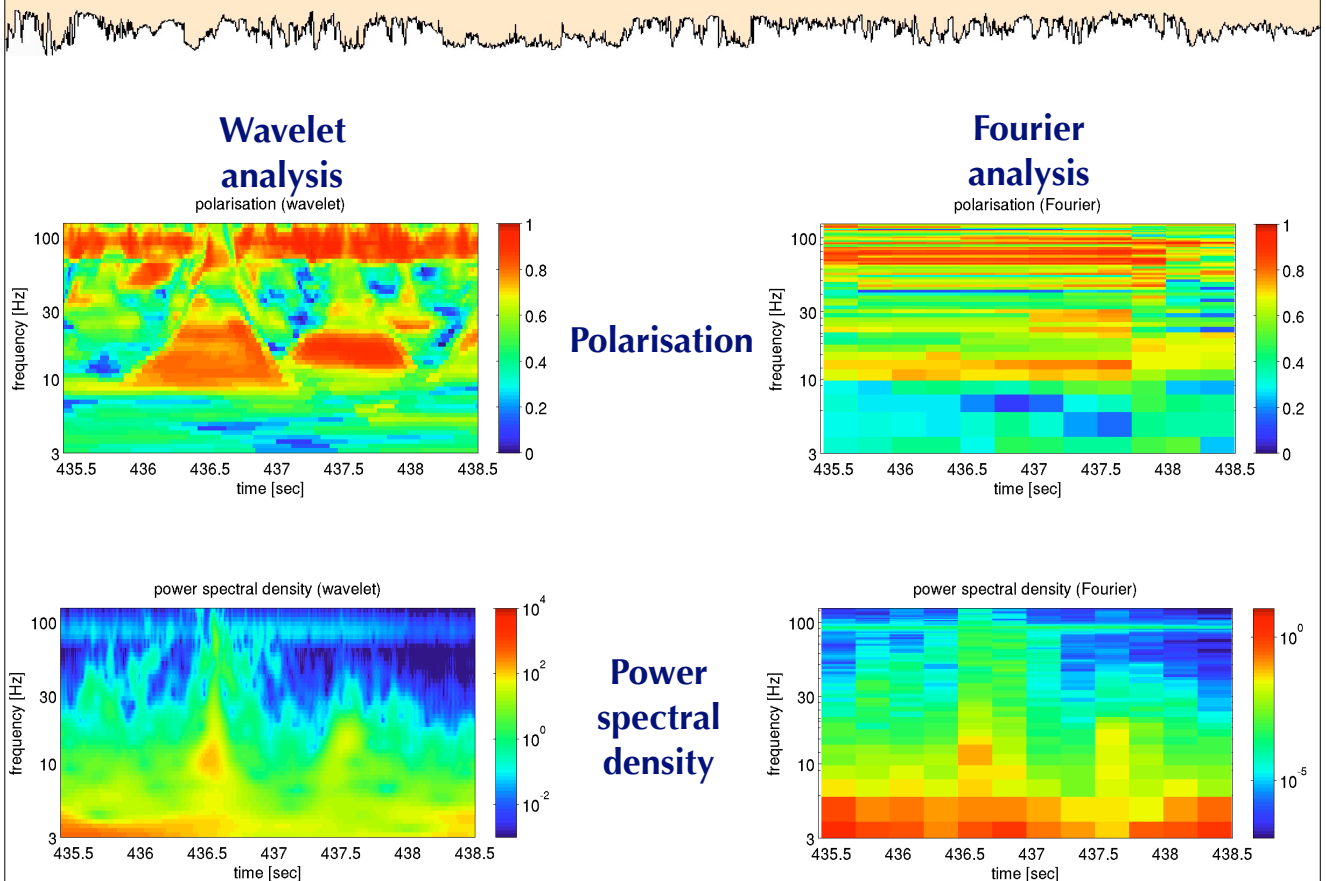
- The polarisation is

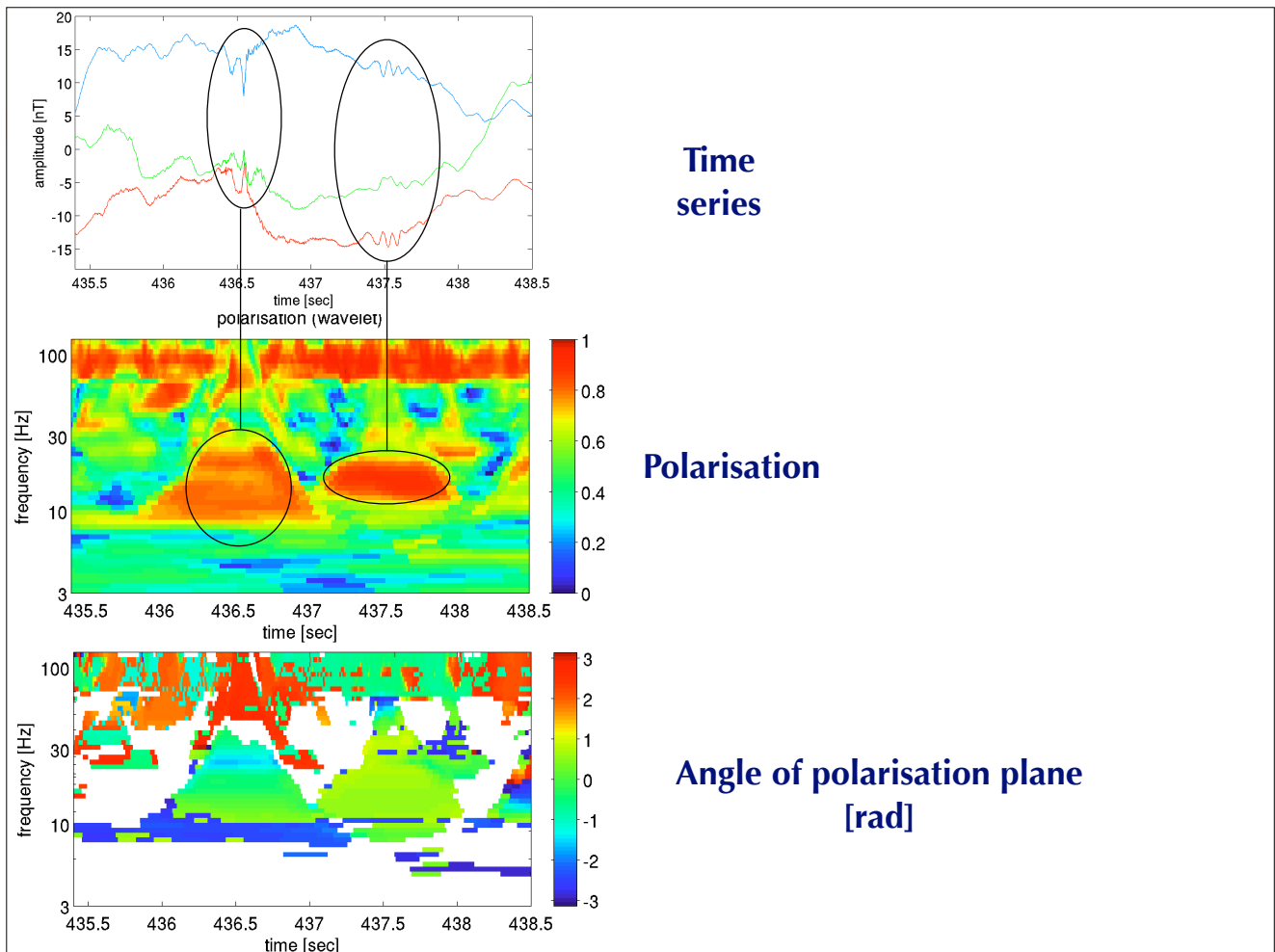
$$0 \leq p \leq 1$$

$$p = 1 - \sqrt{\frac{\lambda_{min}}{\lambda_{max}}}$$



## Fourier vs wavelet





## Conclusion



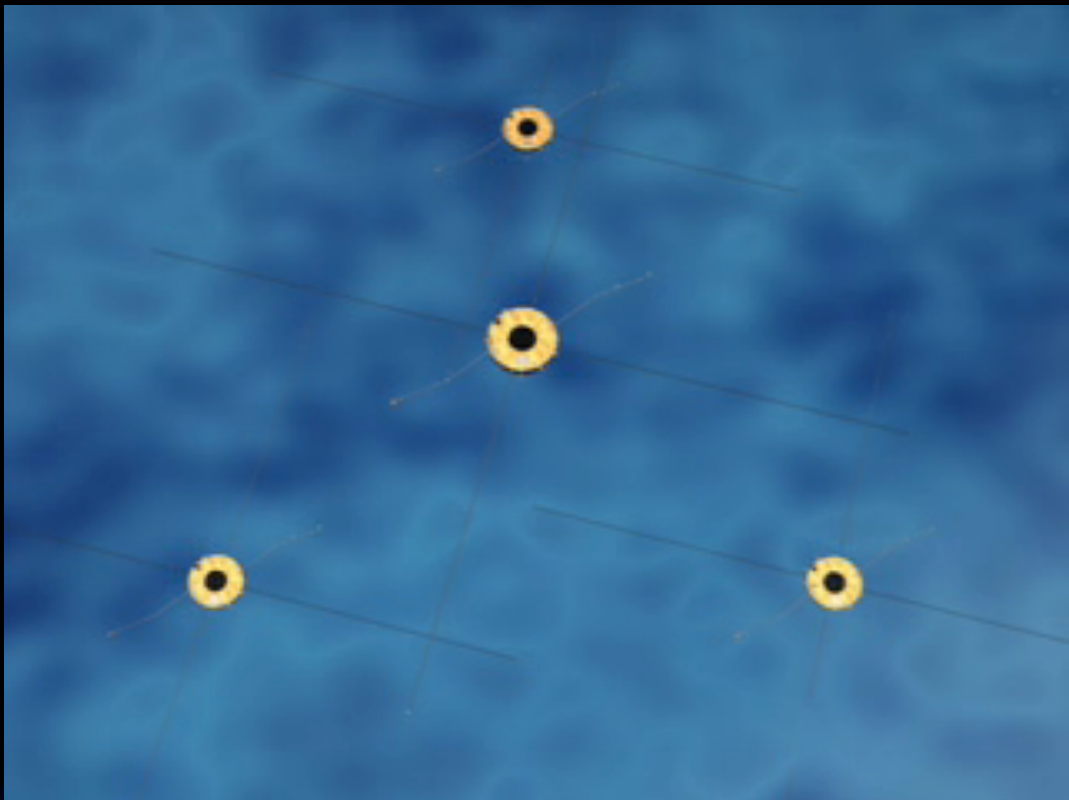
- Wavelets are better suited for the characterization of **wave packets**
- This is important for turbulent wavefields in which different types of waves may coexist
- But the compromise between time and frequency resolution cannot be avoided





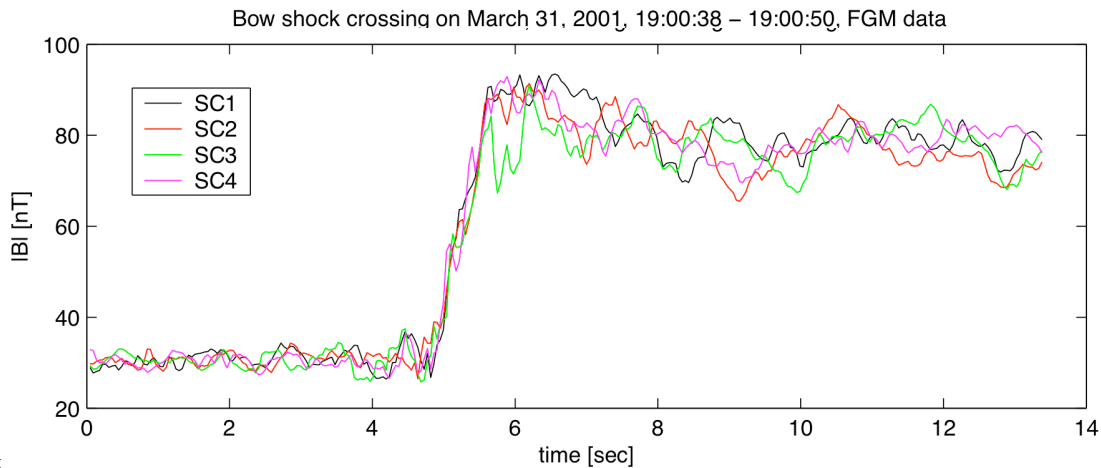
# Example

Automated event timing multipoint data



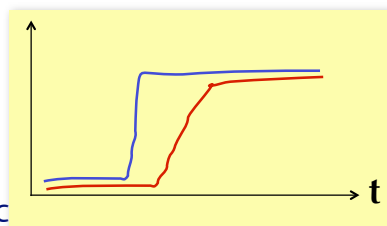
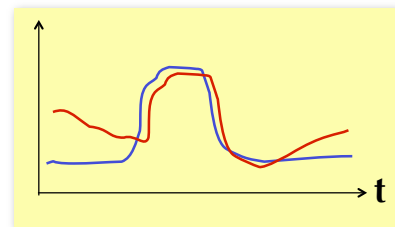
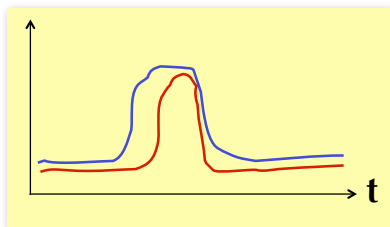
## Timing

- A major fraction of multisatellite data analysis involves timing of specific events on the 4 satellites
- Timing = when does a given event occur on the 4 spacecraft ?



## Timing

- But accurate timing can become quite difficult when one does not observe the same « pattern » on all four spacecraft
- Examples where there is no clear solution



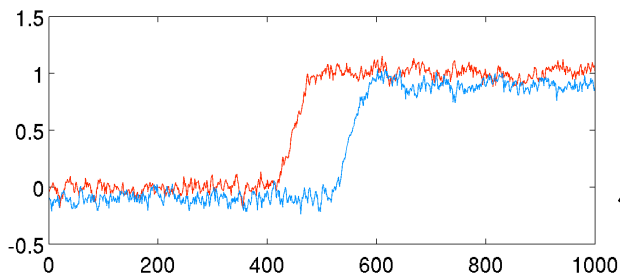
- This becomes a pattern identific

## Standard approach

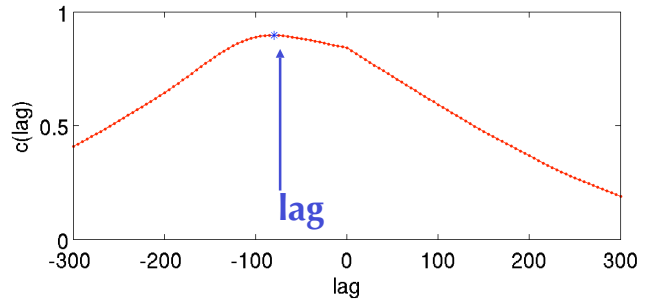


- The usual approach involves computation of the cross-correlation function

$$c(\tau) = \sum_{t=-T}^T \frac{[u_1(t) - \bar{u}_1][u_2(t - \tau) - \bar{u}_2]}{\sigma_{u_1} \sigma_{u_2}}$$



- And looking for the lag that maximizes  $c(\tau)$

 $c(\tau)$ 
 $\tau$ 


## Standard approach



$$c(\tau) = \sum_{t=-T}^T \frac{[u_1(t) - \bar{u}_1][u_2(t - \tau) - \bar{u}_2]}{\sigma_{u_1} \sigma_{u_2}}$$

This cross-correlation approach has many drawbacks

- The width  $T$  of the interval greatly influences the result. How should its value be chosen ?
- It is biased by large scales
- It is easily misled by offsets and trends

We use instead a multiresolution approach that has been developed in the frame of stereoscopic vision

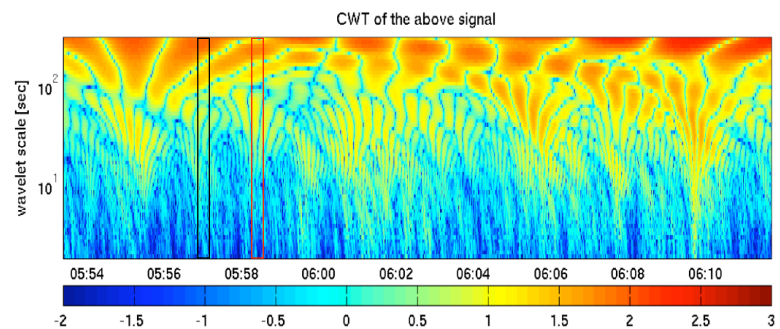
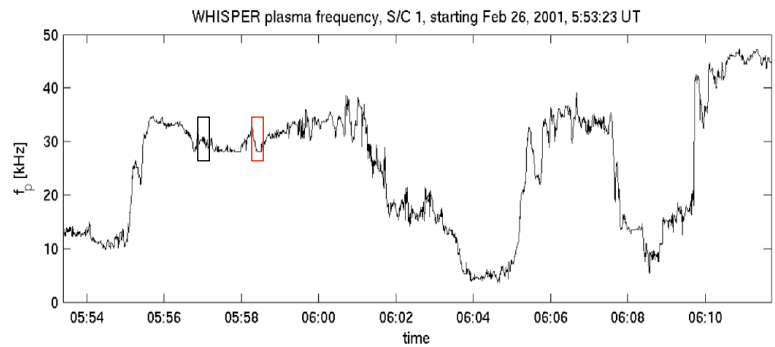
(Perrin & Torresani, 2001)

## Multiresolution approach

At each time  $t$ , the wavelet transform of  $u(t)$  uniquely describes  $u(t)$  in a local neighbourhood : amplitude, derivative, texture, ...

→ similar patterns should therefore have similar wavelet coefficients

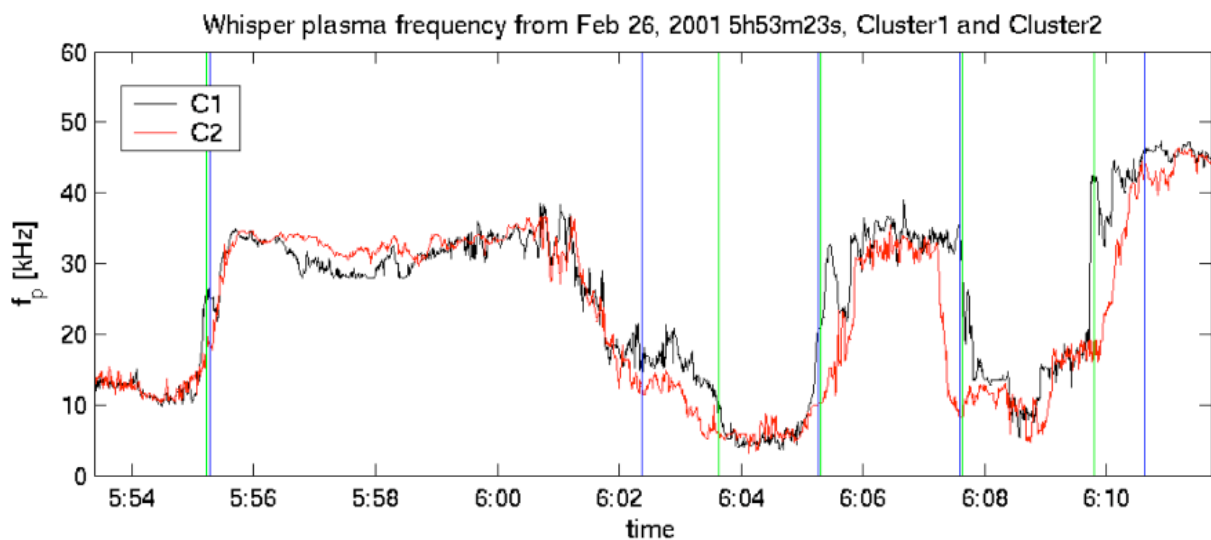
→ instead of comparing the time series, we compare their wavelet coefficients



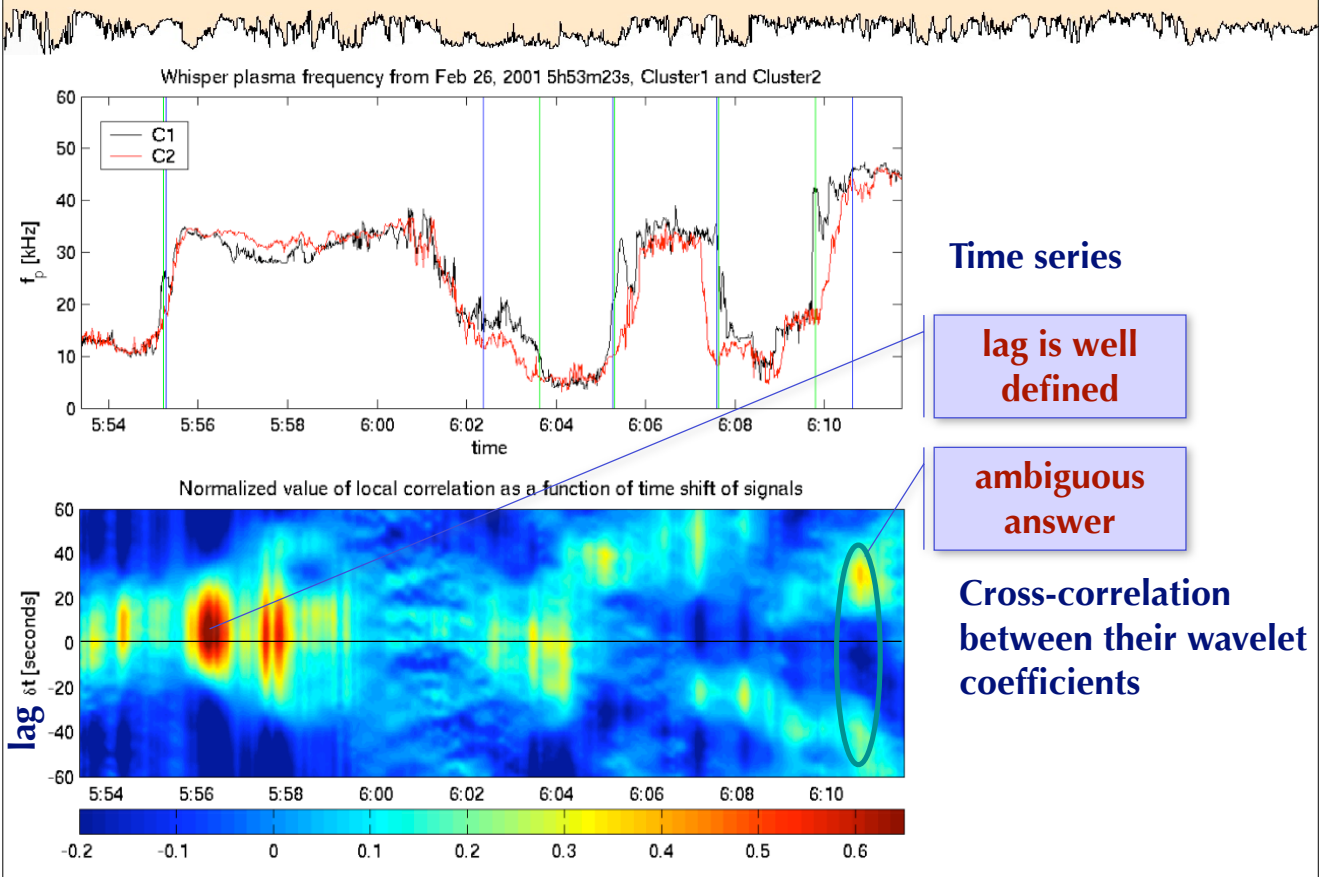
## Example

A test case : electron density measurements

(Soucek et al., 2004)



## Multiresolution approach



## Why is this approach better ?

Multiresolution analysis offers several advantages :

**good resolution** : its resolution can be better than the sampling period  
(for correlation analysis, it is fixed by the window width)

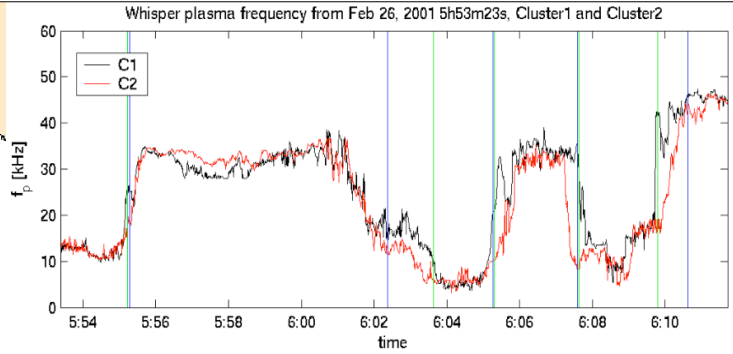
**robustness** : it can handle patterns that do not look exactly the same

it is **data-adaptive** : the method selects itself the scales at which the correlation is the highest (like in artificial vision)

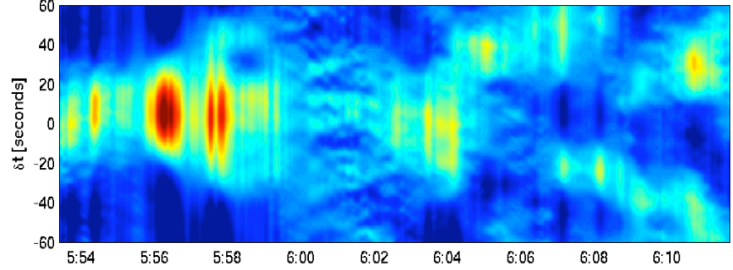
# Final result



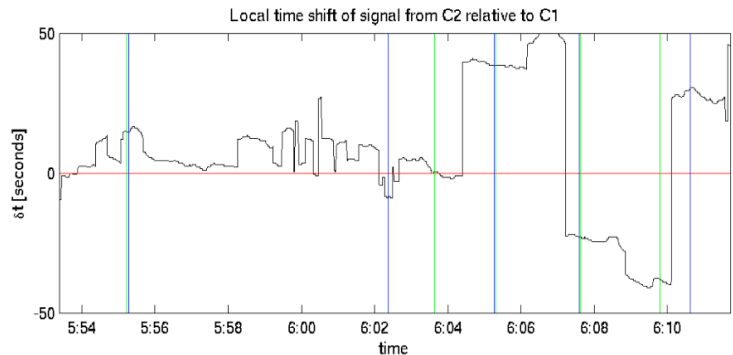
Time series



Cross-correlation between wavelet coefficients



Estimated lag



# Example : comparison of EIT images

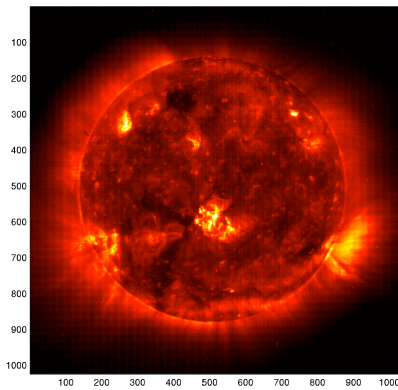


Image at 1:38 = right image

disparity due to solar rotation varies from 0 to 6 pixels

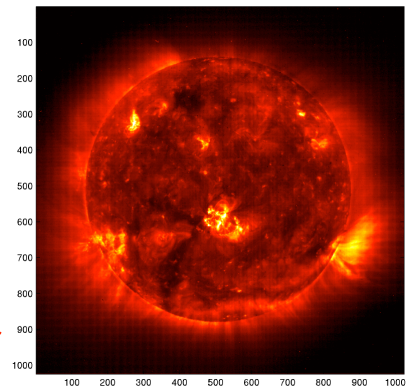
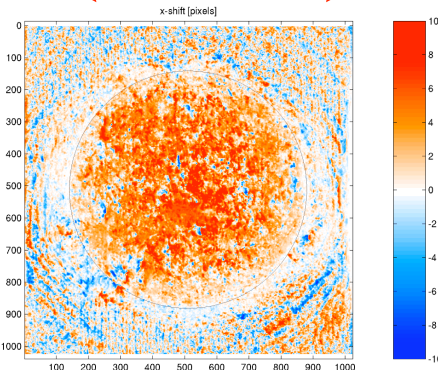


Image at 3:20 = left image

horizontal disparity between 2 images (in pixels)





# Example

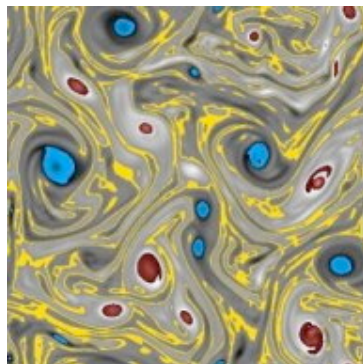
## Extracting coherent structures from turbulent flows

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## Turbulent flows

2 key issues in the analysis of 2D (or 3D) turbulence are:

- how to identify and isolate coherent structures ?



Vorticity field in 2D turbulence  
(Farge et al., 1985)

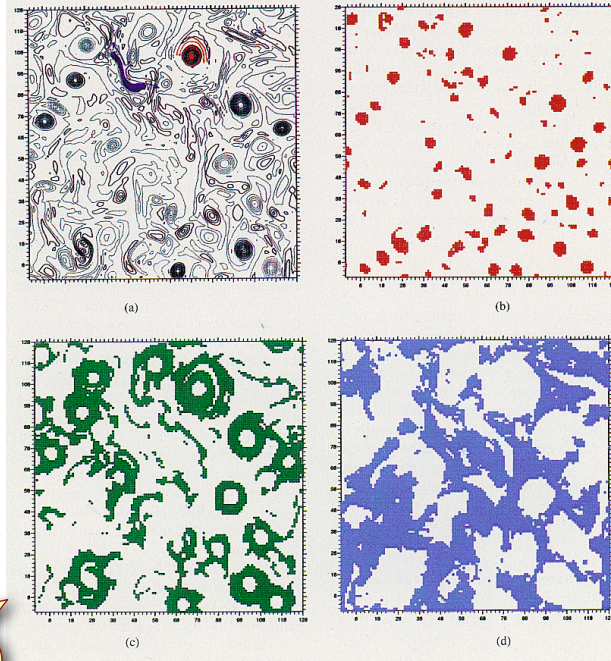
- can Navier-Stokes simulations be improved using spectral methods with better basis functions than Fourier modes ?

$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \vec{v} + \vec{f}$$



# Turbulent flows

2D numerical simulation of vorticity field (Farge et al., 1996)

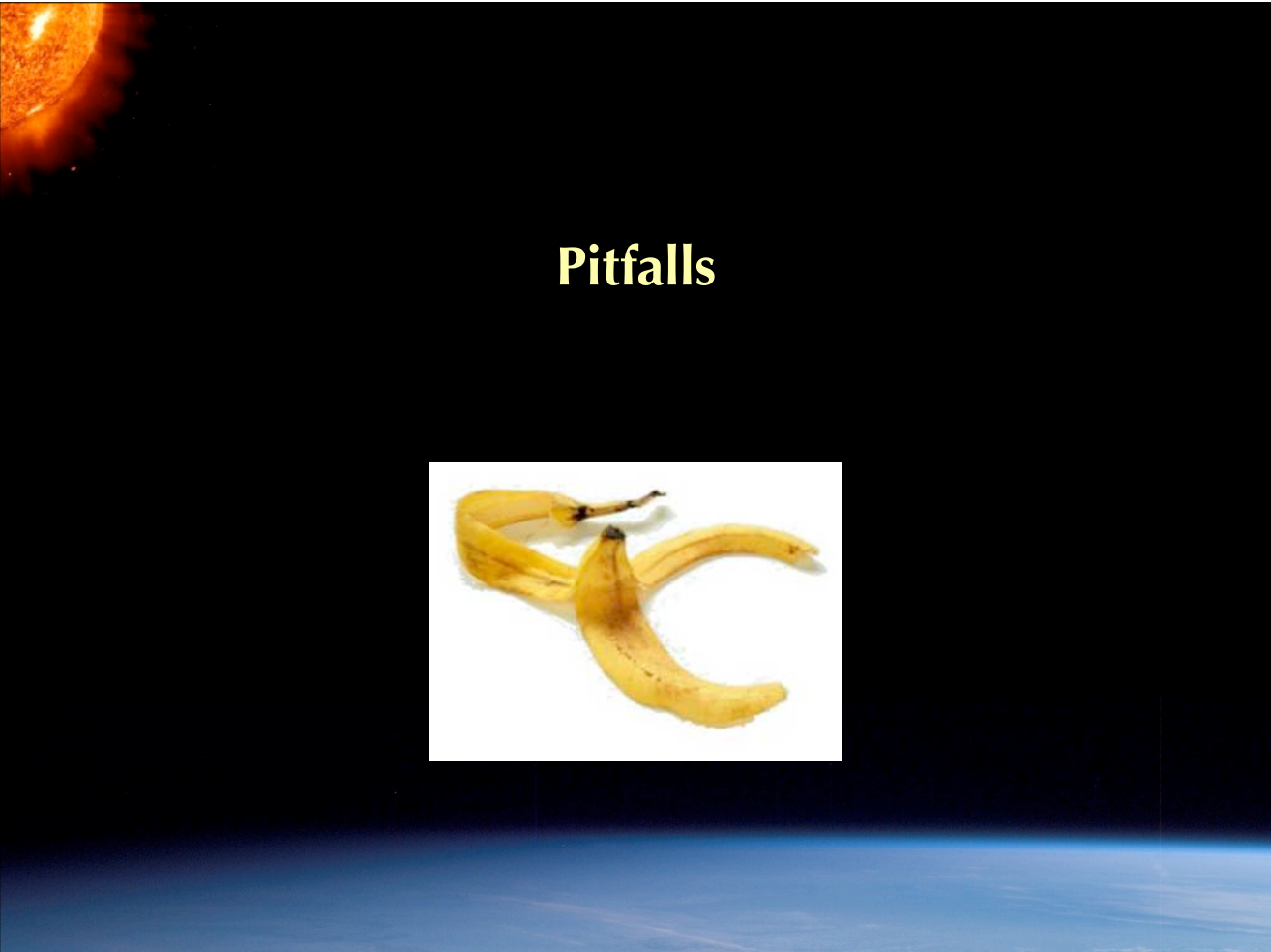


elliptic regions (dominated by rotation) : coherent vortices

shear layers

hyperbolic regions (dominated by deformation) : incoherent background

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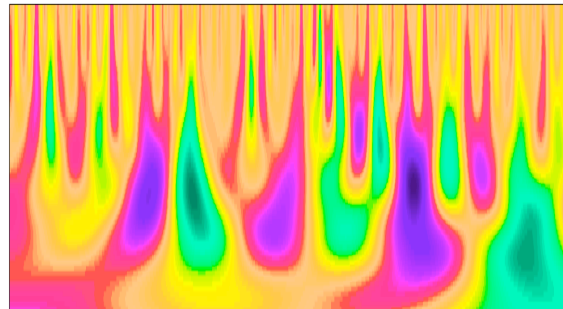
## Pitfalls



## Beware of pitfalls



**Example** : Arneodo et al. (Nature, 1991) claimed to have seen the Richardson cascade in fully developed turbulence, simply by looking at the intricate structure of the wavelet transform



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## Torrence

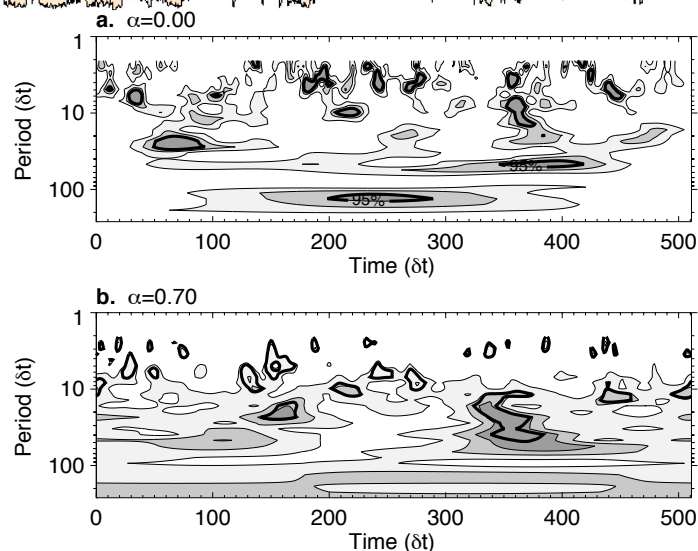
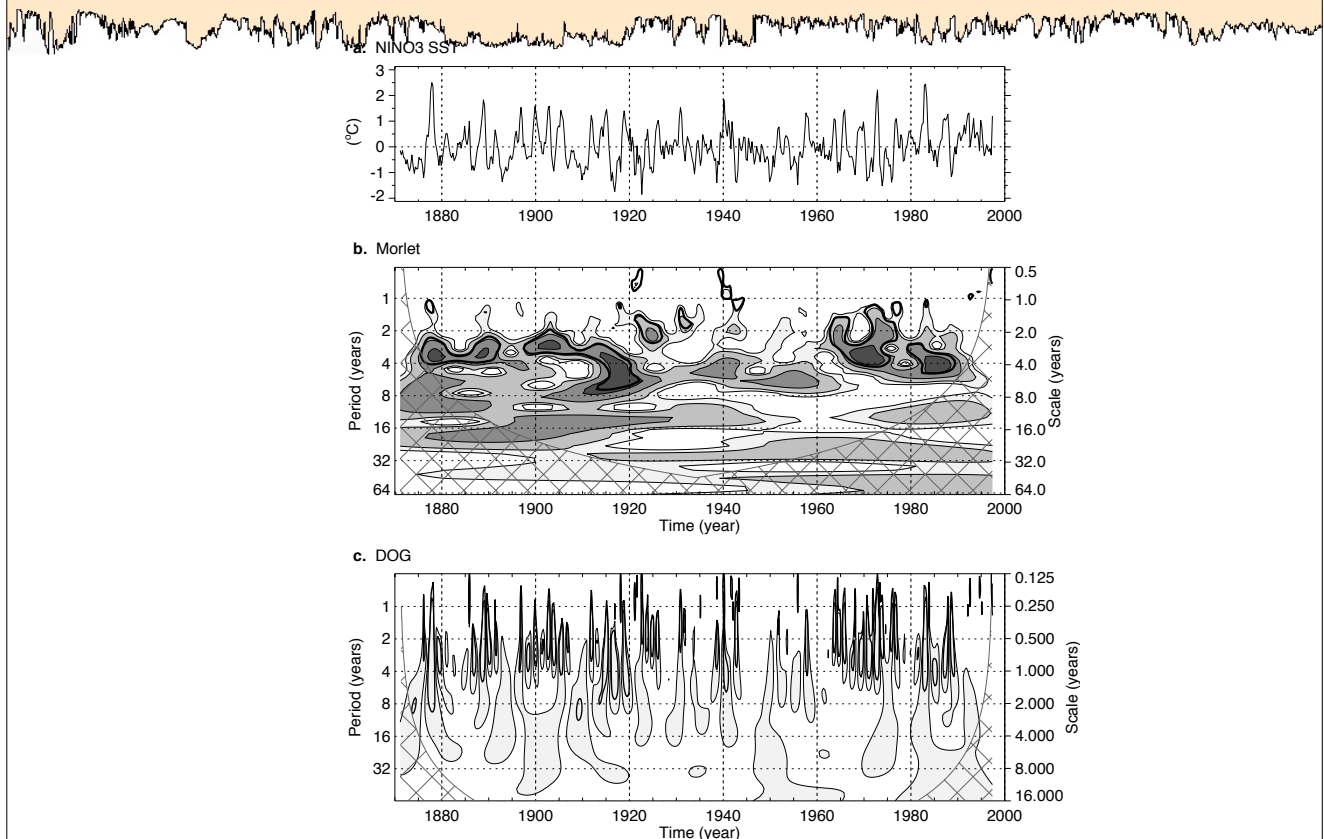


FIG. 4. (a) The local wavelet power spectrum for a Gaussian white noise process of 512 points, one of the 100 000 used for the Monte Carlo simulation. The power is normalized by  $1/\sigma^2$ , and contours are at 1, 2, and 3. The thick contour is the 95% confidence level for white noise. (b) Same as (a) but for a red-noise AR(1) process with lag-1 of 0.70. The contours are at 1, 5, and 10. The thick contour is the 95% confidence level for the corresponding red-noise spectrum.

Warwick, 2/2008



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## Some useful references

- B. Burke Hubbard, *The world according to wavelets: the story of a mathematical technique in the making* (Peters, 1998): a friendly introduction
- I. Daubechies, *Ten lectures on wavelets* (SIAM, 1992): more mathematical, a classic
- S. Mallat, *A wavelet tour of signal processing* (Academic Press, 2002) : excellent reference on wavelets
- J.-L. Starck & F. Murtagh, *Astronomical image and data analysis* (Springer, 2006): wavelets for astronomical images
- *Wavelets and turbulence*: lots of articles by M. Farge, K. Schneider, et al., see <http://wavelets.ens.fr/>
- G. Paschmann, et al. *Analysis Methods for Multi-Spacecraft Data* (ISSI, Bern, 2000), can be downloaded from [http://www.issi.unibe.ch/PDF-Files/analysis\\_methods\\_1\\_1a.pdf](http://www.issi.unibe.ch/PDF-Files/analysis_methods_1_1a.pdf)

Warwick, 2/2008