Local Defects in the Thomas-Fermi-von Weizsäcker Theory of Crystals

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 $\mathcal{I} = \inf\{\mathcal{E}(U_{\Lambda}) \mid U_{\Lambda} \in \mathcal{U}^{1,2}\}, \quad \mathcal{I}^{d} = \inf\{\mathcal{E}^{d}(U_{\Lambda}) \mid U_{\Lambda} \in \mathcal{U}^{1,2}\}.$ (1)

- [Cances, Ehrlacher (2011)] Local defects are always neutral in the Thomas-Fermi-von Weizsäcker theory of crystals, Arch. Rational Mech. Anal.
- [Ehrlacher, Ortner, Shapeev (2013)] Analysis of Boundary Conditions for Crystal Defect Atomistic Simulations, *Preprint*.
- [Catto, Le Bris, Lions (1998)] The Mathematical Theory of Thermodynamic Limits: Thomas–Fermi Type Models, Oxford Mathematical Monographs.

- Introduction
- Exponential Estimates
- Applications of the Exponential Estimates
- Renormalising the Energy Differences
- Outlook

The TFW Model

Let $u, m : \mathbb{R}^3 \to \mathbb{R}_{\geq 0}$ denote

- u square root of the electron density,
- m nuclear distribution,
- ϕ potential generated by u and m, which solves

$$-\Delta\phi=4\pi(m-u^2).$$

The Thomas-Fermi-von Weizsäcker functional is then defined by

$$E^{TFW}(u,m) = \int |\nabla u|^2 + \int u^{10/3} + \frac{1}{2}D(m-u^2,m-u^2),$$

where

$$D(f,g) = \int \int_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{f(x)g(y)}{|x-y|} dx dy.$$

For finite systems, with $u \in H^1(\mathbb{R}^3)$ and $m \in L^{6/5}(\mathbb{R}^3)$, this is well-defined.

Existence and Uniqueness of Solutions

Theorem (Catto, Le Bris, Lions (1998))

Let m be a non-negative function satisfying:

(H1)
$$\sup_{x \in \mathbb{R}^3} \int_{B_1(x)} m(z) \, dz < \infty$$
,
(H2) $\lim_{R \to +\infty} \inf_{x \in \mathbb{R}^3} \frac{1}{R} \int_{B_R(x)} m(z) \, dz = +\infty$,

then there exists a unique distributional solution (u, ϕ) to

$$-\Delta u + \frac{5}{3}u^{7/3} - \phi u = 0,$$
 (2)
$$-\Delta \phi = 4\pi (m - u^2),$$
 (3)

satisfying $u \in H^2_{unif}(\mathbb{R}^3) \cap L^{\infty}(\mathbb{R}^3)$, inf $u > 0, \phi \in L^2_{unif}(\mathbb{R}^3)$.

where $||f||_{L^2_{unif}(\mathbb{R}^3)} = \sup_{x \in \mathbb{R}^3} ||f||_{L^2(B_1(x))}, ||f||_{H^2_{unif}(\mathbb{R}^3)} = \sup_{x \in \mathbb{R}^3} ||f||_{H^2(B_1(x))}.$

Lattice Displacements

Let $\Lambda = \mathbb{Z}^3$ and e_1, e_2, e_3 denote the standard normal basis of \mathbb{R}^3 .

$$\begin{aligned} \mathcal{U} &= \mathcal{U}^{1,2} = \{ U_{\Lambda} : \Lambda \to \mathbb{R}^3 \, | \, \| \nabla U_{\Lambda} \|_{\ell^2(\Lambda)} < \infty \}, \\ \| \nabla U_{\Lambda} \|_{\ell^2(\Lambda)} &= \left(\sum_{l \in \Lambda} \sum_{i=1}^3 | U_{\Lambda}(l+e_i) - U_{\Lambda}(l) |^2 \right)^{1/2} \\ &\sim \| \nabla U_{\Lambda} \|_{\gamma} = \left(\sum_{l \in \Lambda} \sum_{\rho \in \Lambda \smallsetminus \{0\}} e^{-\gamma |\rho|} | U_{\Lambda}(l+\rho) - U_{\Lambda}(l) |^2 \right)^{1/2}. \end{aligned}$$

Useful embeddings,

$$\mathcal{U}^{1,2} \hookrightarrow \ell^6(\Lambda) \hookrightarrow \ell^\infty(\Lambda),$$

by the Gagliardo-Nirenberg-Sobolev embedding theorem.

Also, define

$$\mathcal{U}^{c} = \{ V_{\Lambda} \in \mathcal{U}^{1,2} \mid V_{\Lambda} \text{ has compact support} \}.$$

Nuclear Configurations

Let $\eta \in C^2_c(B_{1/4}(0))$ be non-negative and radial.





$$m_{per}(x) = \sum_{l \in \Lambda} \eta(x - l),$$

$$-\Delta u_{per} + \frac{5}{3} u_{per}^{7/3} - \phi_{per} u_{per} = 0,$$

$$-\Delta \phi_{per} = 4\pi (m_{per} - u_{per}^2),$$

$$m_U(x) = \sum_{l \in \Lambda} \eta(x - l - U_{\Lambda}(l)),$$

$$\begin{aligned} -\Delta u_U + \frac{5}{3} u_U^{7/3} - \phi_U u_U &= 0, \\ -\Delta \phi_U &= 4\pi (m_U - u_U^2). \end{aligned}$$

Homogeneous Energy Difference

For $U_{\Lambda} \in \mathcal{U}^{1,2}$, define

$$\begin{aligned} \mathcal{E}(U_{\Lambda}) &= E^{TFW}(u_U, m_U) - E^{TFW}(u_{per}, m_{per}) \\ &= \int |\nabla u_U|^2 + \int u_U^{10/3} + \frac{1}{2} \int \phi_U(m_U - u_U^2) \\ &- \int |\nabla u_{per}|^2 - \int u_{per}^{10/3} - \frac{1}{2} \int \phi_{per}(m_{per} - u_{per}^2). \end{aligned}$$

Since $m_{per}, m_U \in C^2(\mathbb{R}^3)$, by elliptic regularity

$$u_U, u_{per}, \phi_U, \phi_{per} \in W^{4,\infty}(\mathbb{R}^3).$$

Direct estimates give

$$\begin{aligned} |\mathcal{E}(U_{\Lambda})| &\leq C(\|u_{U} - u_{per}\|_{W^{1,1}(\mathbb{R}^{3})} + \|\phi_{U} - \phi_{per}\|_{L^{1}(\mathbb{R}^{3})} \\ &+ \|m_{U} - m_{per}\|_{L^{1}(\mathbb{R}^{3})}). \end{aligned}$$
(4)

Nuclear Configurations

Let $\rho_{def}^{nuc} \in C_c^2(\mathbb{R}^3)$ represent the density of a defect.



$$\begin{split} m_{per}(x) &= \sum_{l \in \Lambda} \eta(x - l), \qquad m_{U,d}(x) = m_U(x) + \rho_{def}^{nuc}(x - U_{\Lambda}(0)) \ge 0, \\ &- \Delta u_{per} + \frac{5}{3} u_{per}^{7/3} - \phi_{per} u_{per} = 0, \qquad - \Delta u_{U,d} + \frac{5}{3} u_{U,d}^{7/3} - \phi_{U,d} u_{U,d} = 0, \\ &- \Delta \phi_{per} = 4\pi (m_{per} - u_{per}^2), \qquad - \Delta \phi_{U,d} = 4\pi (m_{U,d} - u_{U,d}^2). \end{split}$$

Defective Energy Difference

For $U_{\Lambda} \in \mathcal{U}^{1,2}$, define

$$\mathcal{E}^{d}(U_{\Lambda}) = E^{TFW}(u_{U,d}, m_{U,d}) - E^{TFW}(u_{per}, m_{per})$$

= $\int |\nabla u_{U,d}|^{2} + \int u_{U,d}^{10/3} + \frac{1}{2} \int \phi_{U,d}(m_{U,d} - u_{U,d}^{2})$
- $\int |\nabla u_{per}|^{2} - \int u_{per}^{10/3} - \frac{1}{2} \int \phi_{per}(m_{per} - u_{per}^{2}).$

Similarly,

$$\begin{aligned} |\mathcal{E}^{d}(U_{\Lambda})| &\leq C(\|u_{U,d} - u_{per}\|_{W^{1,1}(\mathbb{R}^{3})} + \|\phi_{U,d} - \phi_{per}\|_{L^{1}(\mathbb{R}^{3})} \\ &+ \|m_{U,d} - m_{per}\|_{L^{1}(\mathbb{R}^{3})}). \end{aligned}$$

Translational Invariance

For any
$$c \in \mathbb{R}^3$$
, $m_{U+c} = m_U(\cdot - c)$, $m_{U+c,d} = m_{U,d}(\cdot - c)$,
 $\Longrightarrow \mathcal{E}(U_{\Lambda} + c) = \mathcal{E}(U_{\Lambda})$, $\mathcal{E}^d(U_{\Lambda} + c) = \mathcal{E}^d(U_{\Lambda})$. (5)

Energy Differences

Questions:

- 1. For $U_{\Lambda} \in \mathcal{U}^{1,2}$, how do we compare (u_U, ϕ_U) and (u_{per}, ϕ_{per}) ?
- 2. Are $\mathcal{E}(U_{\Lambda}), \mathcal{E}^{d}(U_{\Lambda})$ well-defined for all $U_{\Lambda} \in \mathcal{U}^{1,2}$?
- 3. Regularity of $\mathcal{E}, \mathcal{E}^d$?
- 4. Do minimisers decay?

Strategy:

- ► For 1, show exponential estimates.
- Answer 2 and 3 for $U_{\Lambda} \in \mathcal{U}^{c}$.
- ► Utilise lattice symmetries when U_Λ = 0 and translational invariance for all U_Λ ∈ U^{1,2}.
- Find a suitable renormalisation, use the density of \mathcal{U}^{c} in $\mathcal{U}^{1,2}$.
- Answer 2 and 3 for $U_{\Lambda} \in \mathcal{U}^{1,2}$.

Exponential Estimates

Theorem 1

Suppose that m_1, m_2 are nuclear arrangements satisfying (H1), (H2) and $m_2 - m_1 = R_{nuc} \in W^{1,\infty}(\mathbb{R}^3)$. Then there exist unique solutions to the following systems

$$\begin{aligned} -\Delta u_1 + \frac{5}{3} u_1^{7/3} - \phi_1 u_1 &= 0, & -\Delta u_2 + \frac{5}{3} u_2^{7/3} - \phi_2 u_2 &= 0, \\ -\Delta \phi_1 &= 4\pi (m_1 - u_1^2), & -\Delta \phi_2 &= 4\pi (m_1 - u_2^2) + R_{nuc}, \end{aligned}$$
(6)

with $u_1, u_2, \phi_1, \phi_2 \in W^{3,\infty}(\mathbb{R}^3)$. Then there exists $C, \tilde{\gamma} > 0$ such that for all $0 < \gamma \leq \tilde{\gamma}$ and all $y \in \mathbb{R}^3$

$$\int_{\mathbb{R}^3} \sum_{|\alpha| \le 3} \left(|\partial^{\alpha} (u_1 - u_2)(x)|^2 + |\partial^{\alpha} (\phi_1 - \phi_2)(x)|^2 \right) e^{-2\gamma|x-y|} dx$$
$$\leq C \int_{\mathbb{R}^3} \sum_{|\beta| \le 1} |\partial^{\beta} R_{nuc}(x)|^2 e^{-2\gamma|x-y|} dx.$$
(7)

Corollary 2

Suppose the conditions for Theorem 1 are satisfied and in addition $R_{nuc} \in H^1(\mathbb{R}^3)$, then

$$\|u_1 - u_2\|_{H^3(\mathbb{R}^3)} + \|\phi_1 - \phi_2\|_{H^3(\mathbb{R}^3)} \le C \|R_{nuc}\|_{H^1(\mathbb{R}^3)}.$$
(8)

Remark: The constants C from Theorem 1 and Corollary 2 depend on

$$C = C(\inf(u_1 + u_2)^{-1}, \max\{\|u_1\|_{W^{1,\infty}}, \|u_2\|_{W^{1,\infty}}\}, \\ \max\{\|\phi_1\|_{W^{1,\infty}}, \|\phi_2\|_{W^{1,\infty}}\}).$$

Sketch Proof of Theorem 1

Define $w = u_1 - u_2, \psi = \phi_1 - \phi_2$. Key steps:

1. Initial estimates: fix $\xi \in H^1(\mathbb{R}^3)$, then

$$\int_{\mathbb{R}^{3}} |\nabla w|^{2} \xi^{2} + \int_{\mathbb{R}^{3}} w^{2} \xi^{2} + \int_{\mathbb{R}^{3}} |\nabla \psi|^{2} \xi^{2}$$
$$\leq C \left(\int_{\mathbb{R}^{3}} R_{nuc} \psi \xi^{2} + \int_{\mathbb{R}^{3}} \left(w^{2} + \psi^{2} \right) |\nabla \xi|^{2} \right).$$
(9)

2. Higher order estimates:

$$\int_{\mathbb{R}^{3}} \psi^{2} \xi^{2} + \int_{\mathbb{R}^{3}} \sum_{2 \leq |\alpha| \leq 3} \left(|\partial^{\alpha} w|^{2} + |\partial^{\alpha} \psi|^{2} \right) \xi^{2}$$
$$\leq C \left(\int_{\mathbb{R}^{3}} \sum_{|\beta| \leq 1} |\partial^{\beta} R_{nuc}|^{2} \xi^{2} + \int_{\mathbb{R}^{3}} \left(w^{2} + \psi^{2} \right) |\nabla \xi|^{2} \right).$$
(10)

Proof of Theorem 1

3. Combine (9) and (10)

$$\int_{\mathbb{R}^{3}} \sum_{|\alpha| \leq 3} \left(|\partial^{\alpha} w|^{2} + |\partial^{\alpha} \psi|^{2} \right) \xi^{2}$$

$$\leq C^{*} \left(\int_{\mathbb{R}^{3}} \sum_{|\beta| \leq 1} |\partial^{\beta} R_{nuc}|^{2} \xi^{2} + \int_{\mathbb{R}^{3}} \left(w^{2} + \psi^{2} \right) |\nabla \xi|^{2} \right). \tag{11}$$

Let $y \in \mathbb{R}^3$ and $0 < \gamma \le \widetilde{\gamma} = \frac{1}{\sqrt{2C^*}}$, then choose

$$\xi(x) = e^{-\gamma|x-y|} \implies |\nabla\xi(x)|^2 \le \frac{1}{2C^*}\xi^2(x).$$

Finally, this gives

$$\int_{\mathbb{R}^3} \sum_{|\alpha| \leq 3} (|\partial^{\alpha} w|^2 + |\partial^{\alpha} \psi|^2) \xi^2 \leq C \int_{\mathbb{R}^3} \sum_{|\beta| \leq 1} |\partial^{\beta} R_{nuc}|^2 \xi^2.$$

Application I.I - Local Defects

Let $U_{\Lambda} \in \mathcal{U}^{1,2}$, consider



 $R_m = \rho_{def}^{nuc} \in C_c^2(B_R(U_\Lambda(0))).$

Lemma (Application I.I)

Let $U_{\Lambda} \in \mathcal{U}^{1,2}$, then there exists $C(U_{\Lambda}, \rho_{def}^{nuc}), \gamma(U_{\Lambda}, \rho_{def}^{nuc}) > 0$ such that

$$\frac{|(u_{U,d} - u_U)(y)| + |\nabla(u_{U,d} - u_U)(y)|}{|(\phi_{U,d} - \phi_U)(y)| + |\nabla(\phi_{U,d} - \phi_U)(y)|} \le Ce^{-\gamma|y - U_{\Lambda}(0)|}.$$
 (12)

Proof: Let
$$w = u_V - u_{per}, \psi = \phi_V - \phi_{per}$$
.
Applying Theorem 1 gives

$$\begin{split} \int_{\mathbb{R}^3} \sum_{|\alpha| \leq 3} \left(|\partial^{\alpha} w(x)|^2 + |\partial^{\alpha} \psi(x)|^2 \right) e^{-2\gamma|x-y|} \, \mathrm{d}x \\ &\leq C \int_{B_{R'}(z)} \left(|R_m(x)|^2 + |\nabla R_m(x)|^2 \right) \, e^{-2\gamma|x-y|} \, \mathrm{d}x \\ &\leq C \, \|R_m\|_{H^1(B_{R'}(U_{\Lambda}(0)))}^2 \, e^{-2\gamma|y-z|}. \end{split}$$

Application I.I - Local Defects

Idea: Restrict the integral to $B_1(y)$ and use the embedding $H^3(B_1(y)) \hookrightarrow C^{1,\alpha}(B_1(y))$.

$$\begin{split} & (|w(y)|^{2} + |\nabla w(y)|^{2} + |\psi(y)|^{2} + |\nabla \psi(y)|^{2}) \\ & \leq C \left(\|w\|_{C^{1,\alpha}(B_{1}(y))}^{2} + \|\psi\|_{C^{1,\alpha}(B_{1}(y))}^{2} \right) \\ & \leq C \left(\|w\|_{H^{3}(B_{1}(y))}^{2} + \|\psi\|_{H^{3}(B_{1}(y))}^{2} \right) \\ & \leq C e^{2\gamma} \int_{B_{1}(y)} \sum_{|\alpha| \leq 3} \left(|\partial^{\alpha} w(x)|^{2} + |\partial^{\alpha} \psi(x)|^{2} \right) e^{-2\gamma|x-y|} dx \\ & \leq C \int_{\mathbb{R}^{3}} \sum_{|\alpha| \leq 3} \left(|\partial^{\alpha} w(x)|^{2} + |\partial^{\alpha} \psi(x)|^{2} \right) e^{-2\gamma|x-y|} dx \\ & \leq C \|R_{m}\|_{H^{1}(B_{R'}(U_{\Lambda}(0)))}^{2} e^{-2\gamma|y-U_{\Lambda}(0)|}. \end{split}$$

Application I.II - Compact Displacements

Let $V_{\Lambda} \in \mathcal{U}^c$ with $\operatorname{spt}(V_{\Lambda}) \subset B_R(z)$.



$$\begin{aligned} -\Delta u_{per} + \frac{5}{3} u_{per}^{7/3} - \phi_{per} u_{per} = 0, & -\Delta u_V + \frac{5}{3} u_V^{7/3} - \phi_V u_V = 0, \\ -\Delta \phi_{per} = 4\pi (m_{per} - u_{per}^2), & -\Delta \phi_V = 4\pi (m_V - u_V^2), \end{aligned}$$

$$R_m = m_V - m_{per} \in C_c^2(B_{R'}(z)).$$

Lemma (Application I.II)

Let $V_{\Lambda} \in \mathcal{U}^{c}$ with spt $(V_{\Lambda}) \subset B_{R}(z)$. Then there exists $C(V_{\Lambda}), \gamma(V_{\Lambda}) > 0$ such that

$$\frac{|(u_V - u_{per})(y)| + |\nabla(u_V - u_{per})(y)|}{|(\phi_V - \phi_{per})(y)| + |\nabla(\phi_V - \phi_{per})(y)|} \le Ce^{-\gamma|y-z|}.$$
 (13)

Proof: This follows immediately from the proof of Application I.I.

Energy Differences for Compact Displacments

Proposition 3

 $\mathcal{E}, \mathcal{E}^{d} : \mathcal{U}^{c} \to \mathbb{R}$ are well-defined.

Proof: Let $V_{\Lambda} \in \mathcal{U}^{c}$ with spt $(V_{\Lambda}) \subset B_{R}(z)$, then by Applications I.I and I.II,

$$\begin{split} m_V - m_{per}, m_{V,d} - m_{per} \in C^2_c(\mathbb{R}^3) \subset L^1(\mathbb{R}^3), \\ |\nabla(u_V - u_{per})|, |u_V - u_{per}|, |\phi_V - \phi_{per}| \leq Ce^{-\gamma|\cdot - z|} \in L^1(\mathbb{R}^3), \\ |\nabla(u_{V,d} - u_{per})|, |u_{V,d} - u_{per}|, |\phi_{V,d} - \phi_{per}| \\ \leq C(e^{-\gamma|\cdot - z|} + e^{-\gamma|\cdot - U_\Lambda(0)|}) \in L^1(\mathbb{R}^3), \end{split}$$

hence by direct estimates (4),

$$|\mathcal{E}(V_{\Lambda})| < \infty, \qquad |\mathcal{E}^{d}(V_{\Lambda})| < \infty.$$

Application II - Change of Variables Estimates

Let $U_{\Lambda} \in \mathcal{U}^{1,2}$, now consider



Application II - Change of Variables Estimates

We can not estimate $u_U - u_{per}$ directly, instead we use (u_{per}, ϕ_{per}) and U_{Λ} to construct predictor variables $(\tilde{u}_U, \tilde{\phi}_U)$.

Application II

Let $U_{\Lambda} \in \mathcal{U}^{1,2}$, then there exists $(\widetilde{u}_U, \widetilde{\phi}_U) \in W^{4,\infty}(\mathbb{R}^3)$ satisfying

$$\|u_U - \widetilde{u}_U\|_{H^3(\mathbb{R}^3)} + \|\phi_U - \widetilde{\phi}_U\|_{H^3(\mathbb{R}^3)} \le C \|\nabla U_{\Lambda}\|_{\ell^2(\Lambda)}.$$
(14)

Proof: Let $U_{\Lambda} \in \mathcal{U}^{1,2}$. We interpolate U_{Λ} to \mathbb{R}^3 to find $\mathbf{U}, \mathbf{Y} \in C^{\infty}(\mathbb{R}^3, \mathbb{R}^3)$, satisfying:

$$\begin{aligned} \mathbf{Y} \text{ is invertible,} \\ \mathbf{Y}(x) &= x + \mathbf{U}(x), \\ \|\nabla \mathbf{U}\|_{L^2(\mathbb{R}^3)} + \|\nabla^2 \mathbf{U}\|_{L^2(\mathbb{R}^3)} \leq C \|\nabla U_{\Lambda}\|_{\ell^2(\Lambda)}. \end{aligned}$$
(15)

The predictor variables are defined by

$$\widetilde{u}_U = u_{\textit{per}} \circ \mathbf{Y}^{-1}, \qquad \widetilde{\phi}_U = \phi_{\textit{per}} \circ \mathbf{Y}^{-1}$$

Application II - General estimates

Also,

$$m_{per} \circ \mathbf{Y}^{-1} = m_U + R_{nuc},$$

where $R_{nuc} \in C_c^{\infty}(B_R(0))$. The predictors $(\tilde{u}_U, \tilde{\phi}_U)$ satisfy

$$\begin{aligned} -\Delta \widetilde{u}_U + \frac{5}{3} \widetilde{u}_U^{7/3} - \widetilde{\phi}_U \widetilde{u}_U &= R_1, \\ -\Delta \widetilde{\phi}_U &= 4\pi (m_U - \widetilde{u}_U^2) + R_2 + R_{nuc}, \\ &-\Delta \phi_U &= 4\pi (m_U - u_U^2). \end{aligned}$$

The residual terms R_1, R_2, R_{nuc} satisfy

$$\begin{split} \sum_{|\beta| \leq 1} \left(|\partial^{\beta} R_1(x)| + |\partial^{\beta} R_2(x)| + |\partial^{\beta} R_{nuc}(x)| \right) \\ \leq C \left(|\nabla \mathbf{U}(x)| + |\nabla^2 \mathbf{U}(x)| \right) \in L^2(\mathbb{R}^3). \end{split}$$

By Corollary 1, we have

$$\begin{aligned} \|u_U - \widetilde{u}_U\|_{H^3(\mathbb{R}^3)} + \|\phi_U - \widetilde{\phi}_U\|_{H^3(\mathbb{R}^3)} \\ &\leq C \|\nabla \mathbf{U}\|_{L^2(\mathbb{R}^3)} + \|\nabla^2 \mathbf{U}\|_{L^2(\mathbb{R}^3)} \leq C \|\nabla U_{\Lambda}\|_{\ell^2(\Lambda)}. \quad \Box \end{aligned}$$

Application III

Let $U_{\Lambda} \in \mathcal{U}^{1,2}$ and suppose

 $|\nabla U_{\Lambda}(I)| \leq C(1+|I|)^{-j}$

Then

$$\frac{|(u_U - \widetilde{u}_U)(y)| + |\nabla(u_U - \widetilde{u}_U)(y)|}{|(\phi_U - \widetilde{\phi}_U)(y)| + |\nabla(\phi_U - \widetilde{\phi}_U)(y)|} \le C(1 + |y|)^{-j}.$$
 (16)

Proof:

$$ert
abla U_{\Lambda}(I) ert \leq C(1+ert I))^{-j}$$

 $\Longrightarrow ert
abla U(x) ert + ert
abla^2 U(x) ert \leq C(1+ert x))^{-j}$

Application III - Decay estimates

Using Theorem 1,

$$\begin{split} \int_{\mathbb{R}^3} \sum_{|\alpha| \leq 3} & \left(|\partial^{\alpha} (u_U - \widetilde{u}_U)(x)|^2 + |\partial^{\alpha} (\phi_U - \widetilde{\phi}_U)(x)|^2 \right) e^{-2\gamma|x-y|} dx \\ & \leq C \int_{\mathbb{R}^3} \left(|\nabla \mathbf{U}(x)|^2 + |\nabla^2 \mathbf{U}(x)|^2 \right) e^{-2\gamma|x-y|} dx \\ & \leq C \int_{\mathbb{R}^3} (1 + |x|)^{-2j} e^{-2\gamma|x-y|} dx \\ & \leq C (1 + |y|)^{-2j}. \end{split}$$

Arguing as in Application I.I gives

$$egin{aligned} |(u_U - \widetilde{u}_U)(x)| + |
abla(u_U - \widetilde{u}_U)(x)| \ |(\phi_U - \widetilde{\phi}_U)(x)| + |
abla(\phi_U - \widetilde{\phi}_U)(x)| \ &\leq C(1 + |x|)^{-j}. \quad \Box \end{aligned}$$

Variations of the Energy Difference

For
$$U_{\Lambda} \in \mathcal{U}^{1,2}, V_{\Lambda}, V_1, V_2 \in \mathcal{U}^c$$

 $\langle \delta \mathcal{E}(U_{\Lambda}), V_{\Lambda} \rangle$ is well-defined,
 $\langle \delta^2 \mathcal{E}(U_{\Lambda}) V_1, V_2 \rangle = \sum_{i,j \in \Lambda} V_1(i)^T H_{i,j}(U_{\Lambda}) V_2(j),$

where for $i \neq j$

$$\begin{aligned} H_{i,j}(U_{\Lambda}) &= -\int \Psi_{U,i}(x) \cdot \nabla \eta (x-j-U_{\Lambda}(j)) \, \mathrm{d}x, \\ |H_{i,j}(U_{\Lambda})| &\leq C \, e^{-\gamma |i+U_{\Lambda}(i)-i-U_{\Lambda}(j)|}. \end{aligned}$$

Translational Invariance:

$$\sum_{i\in\Lambda}H_{i,j}(U_{\Lambda})=\sum_{j\in\Lambda}H_{i,j}(U_{\Lambda})=0.$$

Consequently

$$\langle \delta^2 \mathcal{E}(U_{\Lambda}) V_1, V_2 \rangle = \sum_{i,j \in \Lambda} V_1(i)^T H_{i,j}(U_{\Lambda}) \left(V_2(j) - V_2(i) \right).$$

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Renormalisation Sketch (*)

Lattice symmetries: for $U_{\Lambda}=0, V_{\Lambda}, V_1, V_2 \in \mathcal{U}^{1,2}$

$$egin{aligned} \mathcal{E}(0) &= 0, \quad \langle \delta \mathcal{E}(0), V_{\mathsf{A}}
angle &= 0, \ &|\langle \delta^2 \mathcal{E}(0) V_1, V_2
angle| \leq C \|
abla V_1 \|_{\ell^2(\mathsf{A})} \|
abla V_2 \|_{\ell^2(\mathsf{A})}, \end{aligned}$$

using the lattice Fourier transform.

Changes of Variables Estimate:

$$\begin{aligned} |H_{i,j}(U_{\Lambda}) - H_{i,j}(0)| &= |\int \Psi_{U,i}(x) \cdot \nabla \eta(x - j - U_{\Lambda}(j)) \, \mathrm{d}x \\ &- \int \Psi_{0,i}(x) \cdot \nabla \eta(x - j) \, \mathrm{d}x| \\ &\approx \left| \int (\Psi_{U,i}(x) - \widetilde{\Psi}_{U,i}(x)) \cdot \nabla \eta(x - j - U_{\Lambda}(j)) \, \mathrm{d}x \right| \\ &\approx F_i(U_{\Lambda}) e^{-\gamma |x - i|}. \end{aligned}$$

 $F(U_{\Lambda}) : \Lambda \to \mathbb{R}$ is approximately

$$F_i(U_{\Lambda}) \approx \left(\int |\nabla \mathbf{U}(x)|^2 e^{-2\gamma |x-i|} \, \mathrm{d}x \right)^{1/2},$$
$$\|F(U_{\Lambda})\|_{\ell^2(\Lambda)} \leq C \|\nabla U_{\Lambda}\|_{\ell^2(\Lambda)}.$$

Renormalisation Sketch (*)

For $V_1, V_2 \in \mathcal{U}^c$

$$\begin{aligned} |\langle \left(\delta^{2} \mathcal{E}(U_{\Lambda}) - \delta^{2} \mathcal{E}(0) \right) V_{1}, V_{2} \rangle | \\ &\leq \sum_{i,j \in \Lambda} |V_{1}(i)| F_{i}(U_{\Lambda}) e^{-\gamma |i-j|} |V_{2}(j) - V_{2}(i)| \\ &\leq \sum_{l \in \Lambda} \sum_{\rho \in \Lambda \smallsetminus \{0\}} \left(|V_{1}(l)| F_{l}(U_{\Lambda}) e^{-\gamma / 2 |\rho|} \right) \left(e^{-\gamma / 2 |\rho|} |D_{\rho} V_{2}(l)| \right) \\ &\leq \left(\sum_{l \in \Lambda} \sum_{\rho \in \Lambda \smallsetminus \{0\}} |V_{1}(l)|^{2} F_{l}(U_{\Lambda})^{2} e^{-\gamma |\rho|} \right)^{1/2} \\ &\quad \cdot \left(\sum_{l \in \Lambda} \sum_{\rho \in \Lambda \smallsetminus \{0\}} e^{-\gamma |\rho|} |D_{\rho} V_{2}(l)|^{2} \right)^{1/2} \\ &\leq C \|V_{1}\|_{\ell^{\infty}(\Lambda)} \|F(U_{\Lambda})\|_{\ell^{2}} \|\nabla V_{2}\|_{\gamma} \\ &\leq C \|\nabla U_{\Lambda}\|_{\ell^{2}(\Lambda)} \|\nabla V_{1}\|_{\ell^{2}(\Lambda)} \|\nabla V_{2}\|_{\ell^{2}(\Lambda)}. \end{aligned}$$

Renormalisation Sketch (*)

Hence for $V_1, V_2 \in \mathcal{U}^{1,2}$

$$\begin{split} |\langle \delta^2 \mathcal{E}(U_{\Lambda}) V_1, V_2 \rangle| &\leq |\langle \delta^2 (\mathcal{E}(U_{\Lambda}) - \mathcal{E}(0)) V_1, V_2 \rangle| + |\langle \delta^2 \mathcal{E}(0) V_1, V_2 \rangle| \\ &\leq C (1 + \|\nabla U_{\Lambda}\|_{\ell^2(\Lambda)}) \|\nabla V_1\|_{\ell^2(\Lambda)} \|\nabla V_2\|_{\ell^2(\Lambda)}. \end{split}$$

Renormalisation: for $U_{\Lambda} \in \mathcal{U}^{c}$

$$\begin{split} \mathcal{E}(U_{\Lambda}) &= \mathcal{E}(0) + \langle \delta \mathcal{E}(0), U_{\Lambda} \rangle + \int_{0}^{1} (1-t) \langle \delta^{2} \mathcal{E}(tU_{\Lambda}) U_{\Lambda}, U_{\Lambda} \rangle \, \mathrm{d}t \\ &= \int_{0}^{1} (1-t) \langle \delta^{2} \mathcal{E}(tU_{\Lambda}) U_{\Lambda}, U_{\Lambda} \rangle \, \mathrm{d}t. \end{split}$$

So

$$|\mathcal{E}(U_{\Lambda})| \leq C \left(1 + \|
abla U_{\Lambda}\|_{\ell^{2}(\Lambda)}
ight) \|
abla U_{\Lambda}\|_{\ell^{2}(\Lambda)}^{2},$$

hence \mathcal{E} can be continuously extended to $\mathcal{U}^{1,2}$. Similarly, this can be used to show \mathcal{E}^d can also be extended to $\mathcal{U}^{1,2}$.

Exponential estimates.

- Estimates for compact displacements.
- Change of variables, global L^2 estimates.
- Decay of displacement implies decay of correctors.
- Exploiting lattice symmetries for U_Λ = 0, translational invariance for U_Λ ∈ U^{1,2}, equivalence of norms and embeddings of U^{1,2}.
- Renormalisation implies the minimisation problems (1) are well-defined.

Outlook:

- Approximating the global minimisation problem by a finite problem. Numerical simulations.
- Dislocations.

Thank you for your attention!