# DISTORTION AND ENTROPY FOR AUTOMORPHISMS OF FREE GROUPS 

Richard Sharp<br>University of Manchester


#### Abstract

Recently, several numerical invariants have been introduced to characterize the distortion induced by automorphisms of a free group. We unify these by interpreting them in terms of an entropy function of a kind familiar in thermodynamic ergodic theory. We draw an analogy between this approach and the Manhattan curve associated to a pair of hyperbolic surfaces.


## 0 . Introduction

Let $F$ be a free group on $k \geq 2$ generators and let $\mathcal{A}=\left\{a_{1}, \ldots, a_{k}\right\}$ be a free basis. We define the word length $|\cdot|=|\cdot|_{\mathcal{A}}$ (with respect to $\mathcal{A}$ ) by $|1|=0$ and, for $x \neq 1$,

$$
|x|=\min \left\{n: x=x_{0} x_{1} \cdots x_{n-1}, x_{i} \in \mathcal{A} \cup \mathcal{A}^{-1}\right\}
$$

where $\mathcal{A}^{-1}=\left\{a_{1}^{-1}, \ldots, a_{k}^{-1}\right\}$. Recall that any $x \neq 1$ may be written uniquely as

$$
x=x_{0} x_{1} \cdots x_{n-1},
$$

where $n=|x|, x_{i} \in \mathcal{A} \cup \mathcal{A}^{-1}, i=0, \ldots, n-1$, and $x_{i+1} \neq x_{i}^{-1}, i=0, \ldots, n-2$. We call such an expression a reduced word.

Let $\operatorname{Aut}(F)$ denote the group of automorphisms of $F$. An automorphism $\phi$ is said to be inner if it is a conjugation, i.e., $\phi(x)=y^{-1} x y$, for some $y \in F$. If an automorphism $\phi$ acts by permuting $\mathcal{A} \cup \mathcal{A}^{-1}$ then we call $\phi$ a permutation automorphism. Following the notation of [7], [12], we say that $\phi$ is simple if it is the product of an inner automorphism and a permutation automorphism.

Let $\partial F$ denote the boundary of $F$ in the sense of the theory of hyperbolic groups [5], [6]. This is a Cantor set and may be naturally identified with the one-sided shift space

$$
\Sigma^{+}=\left\{\left(x_{n}\right)_{n=0}^{\infty} \in\left(\mathcal{A} \cup \mathcal{A}^{-1}\right)^{\mathbb{Z}^{+}}: x_{n+1} \neq x_{n}^{-1}, n \geq 0\right\}
$$

There is a dynamical systems associated to this space, namely the shift map $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$, and its ergodic theory will play a key role in this paper.

Define a $\sigma$-invariant Borel probability measure $\mu_{0}$ on $\Sigma^{+}$by

$$
\mu_{0}\left(\left[x_{0}, x_{1}, \ldots, x_{n-1}\right]\right)=\frac{1}{2 k(2 k-1)^{n-1}}
$$

where $\left[x_{0}, x_{1}, \ldots, x_{n-1}\right]=\left\{\left(y_{n}\right)_{n=0}^{\infty} \in \Sigma^{+}: y_{i}=x_{i}, i=0, \ldots, n-1\right\}$. This measure allows us to describe generic behaviour of sequences in $F$, with respect to the basis $\mathcal{A}$.

In this paper, we shall be interested in various measures of the distortion of $F$ under an automorphism $\phi: F \rightarrow F$. The first of these quantifies the generic distortion or stretching. Following Kaimanovich, Kapovich and Schupp [7], define the generic stretching factor $\lambda(\phi)$ (with respect to $\mathcal{A}$ ) by

$$
\begin{aligned}
\lambda(\phi) & =\lim _{n \rightarrow+\infty} \int \frac{\left|\phi\left(x_{0} x_{1} \cdots x_{n-1}\right)\right|}{n} d \mu_{0}\left(\left(x_{n}\right)_{n=0}^{\infty}\right) \\
& =\lim _{n \rightarrow+\infty} \frac{\left|\phi\left(x_{0} x_{1} \cdots x_{n-1}\right)\right|}{n} \quad \text { for } \mu_{0} \text {-a.e. }\left(x_{n}\right)_{n=0}^{\infty} .
\end{aligned}
$$

Proposition 1 [7]. For any $\phi \in \operatorname{Aut}(F), \lambda(\phi) \geq 1$. Furthermore, $\lambda(\phi)=1$ if and only if $\phi$ is simple.

We now make a further definition, which measures the proportion of elements which are not stretched by a factor greater than one. Following Myasnikov and Shpilrain [12], define the curl of $\phi, \operatorname{Curl}(\phi)$, by

$$
\operatorname{Curl}(\phi)=\limsup _{n \rightarrow+\infty}\left(\frac{\#\{x \in F:|x| \leq n,|\phi(x)| \leq n\}}{\#\{x \in F:|x| \leq n\}}\right)^{1 / n}
$$

i.e., the growth rate of the proportion of points in the balls $\{x \in F:|x| \leq n\}$ which remain there under $\phi$.

Proposition 2 [12]. For any $\phi \in \operatorname{Aut}(F), 0<\operatorname{Curl}(\phi) \leq 1$. Furthermore, $\operatorname{Curl}(\phi)=1$ if and only if $\phi$ is simple.

Finally, we define a set introduced by Kapovich, which captures all possible distortions induced by $\phi$. Let $\mathcal{C}(F)$ denote the set of all non-trivial conjugacy classes in $F$ and note that, for $w \in \mathcal{C}(F), \phi(w) \in \mathcal{C}(F)$ is well-defined. Following Kapovich [8], we define the conjugacy distortion spectrum of $\phi$ to be the set

$$
\mathcal{D}_{\phi}=\left\{\frac{|\phi(w)|}{|w|}: w \in \mathcal{C}(F)\right\}
$$

where $|w|=\min \{|x|: x \in w\}$, and let $\overline{\mathcal{D}_{\phi}}$ denote the closure of $\mathcal{D}_{\phi}$. The structure of $\mathcal{D}_{\phi}$ was studied by Kapovich, who showed the following.
Proposition 3 [8]. $\overline{\mathcal{D}_{\phi}}$ is a closed interval with rational endpoints and $\mathcal{D}_{\phi}=\overline{\mathcal{D}_{\phi}} \cap \mathbb{Q}$. If $\phi$ is simple then $\mathcal{D}_{\phi}=\{1\}$. If $\phi$ is not simple then $\overline{\mathcal{D}_{\phi}}$ has non-empty interior which contains 1 and $\lambda(\phi)$.

These quantities may be related by the following theorem.

Theorem. Suppose that $\phi$ is not simple. Then there exists a strictly concave analytic function $\mathfrak{h}: \operatorname{int}\left(\overline{\mathcal{D}_{\phi}}\right) \rightarrow \mathbb{R}^{+}$such that, for each $\rho \in \operatorname{int}\left(\overline{\mathcal{D}_{\phi}}\right), 0<\mathfrak{h}(\rho) \leq \log (2 k-1)$ and, if $\rho$ is rational, then

$$
\begin{equation*}
\limsup _{n \rightarrow+\infty} \frac{1}{n} \log \#\left\{w \in \mathcal{C}(F):|w|=n, \frac{|\phi(w)|}{|w|}=\rho\right\}=\mathfrak{h}(\rho) \tag{0.1}
\end{equation*}
$$

Furthermore, $\mathfrak{h}(\rho)=\log (2 k-1)$ if and only if $\rho=\lambda(\phi)$ and $\operatorname{Curl}(\phi)=e^{\mathfrak{h}(1)} /(2 k-1)$.
This will be proved as Theorems 1, 2 and 3 below.
We shall now outline the contents of the paper. In section 1, we discuss the thermodynamic formalism associated to a class of dynamical systems called subshifts of finite. In section 2, we discuss the subshift associated to a free group, the relationship between periodic orbits and conjugacy classes and how to encode the quantity $|\phi(\cdot)|$ in terms of a function on this subshift. We also introduce the function $\mathfrak{h}$ and relate it to the generic stretch. In section 3, we study the conjugacy distortion spectrum via the periodic points of the shift map, proving equation (0.1). In sections 4 and 5 , we show how to obtain the relationship between $\operatorname{Curl}(\phi)$ and $\mathfrak{h}(1)$. In the final section, we recast our results in terms of a "Manhattan curve", analogous to that associated by Burger to a pair of hyperbolic surfaces.

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## 1. Thermodynamic Formalism

In this section, we shall describe the ergodic theory associated to the shift map $\sigma$ : $\Sigma^{+} \rightarrow \Sigma^{+}$. The theory we describe goes under the name of thermodynamic formalism; standard references are [13], [14, Appendix II], [23].

We begin with the definition of a subshift of finite type. Let $A$ be finite matrix, indexed by a set $\mathcal{I}$, with entries zero and one. We define the shift space

$$
\Sigma_{A}^{+}=\left\{\left(x_{n}\right)_{n=0}^{\infty} \in \mathcal{I}^{\mathbb{Z}^{+}}: A\left(x_{n}, x_{n+1}\right)=1 \forall n \in \mathbb{Z}^{+}\right\}
$$

and the (one-sided) subshift of finite type $\sigma: \Sigma_{A}^{+} \rightarrow \Sigma_{A}^{+}$by $(\sigma x)_{n}=x_{n+1}$. We give $\mathcal{I}$ the discrete topology, $\mathcal{I}^{\mathbb{Z}^{+}}$the product topology and $\Sigma_{A}^{+}$the subspace topology. A compatible metric is given by

$$
d\left(\left(x_{n}\right)_{n=0}^{\infty},\left(y_{n}\right)_{n=0}^{\infty}\right)=\sum_{n=0}^{\infty} \frac{1-\delta_{x_{n} y_{n}}}{2^{n}}
$$

where $\delta_{i j}$ is the Kronecker symbol.
We say that $A$ is irreducible if, for each $(i, j) \in \mathcal{I}^{2}$, there exists $n(i, j) \geq 1$ such that $A^{n(i, j)}(i, j)>0$ and aperiodic if there exists $n \geq 1$ such that, for each $(i, j) \in \mathcal{I}^{2}$, $A^{n}(i, j)>0$. The latter statement is equivalent to $\sigma: \Sigma_{A}^{+} \rightarrow \Sigma_{A}^{+}$being topologically mixing (i.e. that there exists $n \geq 1$ such that for any two non-empty open sets $U, V \subset \Sigma_{A}^{+}$, $\sigma^{-m}(U) \cap V \neq \varnothing$, for all $\left.m \geq n\right)$.

If $A$ is aperiodic then it has a positive simple eigenvalue $\beta$ which is strictly maximal in modulus (i.e. every other eigenvalue has modulus strictly less that $\beta$ ) and the topological entropy $h(\sigma)$ of $\sigma$ is equal to $\log \beta$.

If an ordered $n$-tuple $\left(x_{0}, x_{1}, \ldots, x_{n-1}\right) \in \mathcal{I}^{n}$ is such that $A\left(x_{m}, x_{m+1}\right)=1, m=$ $0,1, \ldots, n-2$ then we say that $\left(x_{0}, x_{1}, \ldots, x_{n-1}\right)$ is an allowed word of length $n$ in $\Sigma_{A}^{+}$; the set of these is denoted $W_{A}^{(n)}$. If $\sigma^{n} x=x$ then we say that $\left\{x, \sigma x, \ldots, \sigma^{n-1} x\right\}$ is a periodic orbit for $\sigma$. Clearly any such an $x$ is obtained by repeating a word $\left(x_{0}, x_{1}, \ldots, x_{n-1}\right) \in$ $W_{A}^{(n)}$ with the additional property that $A\left(x_{n-1}, x_{0}\right)=1$. Note that we regard the periodic orbits $\left\{x, \sigma x, \ldots, \sigma^{n-1} x\right\},\left\{x, \sigma x, \ldots, \sigma^{n-1} x\right\}, x, \ldots, \sigma^{n-1} x$, etc., as distinct objects (even though they are identical as point sets). If $\sigma^{n} x=x$ but $\sigma^{m} x \neq x$ for $0<m<n$ then we say that $\left\{x, \sigma x, \ldots, \sigma^{n-1} x\right\}$ is a prime periodic orbit.

It is sometimes convenient to replace $\sigma: \Sigma_{A}^{+} \rightarrow \Sigma_{A}^{+}$by a shift of finite type on a new space $\Sigma_{A_{N}}^{+}$, whose symbols are $W_{A}^{(N)}$ and for which $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{N-1}\right)$ may be followed by $\mathbf{y}=\left(y_{0}, y_{1}, \ldots, y_{N-1}\right)$ if and only if $y_{n}=x_{n+1}, n=0,1, \ldots, N-2$. (Since $\mathbf{x} \in W_{A}^{(N)}$, the latter condition automatically implies that $A\left(x_{N-1}, y_{N-1}\right)=1$.) We continue to denote the shift map by $\sigma: \Sigma_{A_{N}}^{+} \rightarrow \Sigma_{A_{N}}^{+}$. More formally, define a $W_{A}^{(N)} \times W_{A}^{(N)}$ zero-one matrix $A_{N}$ by

$$
A_{N}(\mathbf{x}, \mathbf{y})=\left\{\begin{array}{l}
1 \text { if } y_{n}=x_{n+1}, n=0, \ldots, N-2 \\
0 \text { otherwise }
\end{array}\right.
$$

Then

$$
\Sigma_{A_{N}}^{+}=\left\{\left(\mathbf{x}_{n}\right)_{n=0}^{\infty} \in\left(W_{A}^{(n)}\right)^{\mathbb{Z}^{+}}: A_{N}\left(\mathbf{x}_{n}, \mathbf{x}_{n+1}\right) \forall n \in \mathbb{Z}^{+}\right\} .
$$

Note that there is a natural period preserving bijection between periodic points for $\sigma$ : $\Sigma_{A}^{+} \rightarrow \Sigma_{A}^{+}$and $\sigma: \Sigma_{A_{N}}^{+} \rightarrow \Sigma_{A_{N}}^{+}$.

Let $\mathcal{M}$ denote the space of all Borel probability measures on $\Sigma_{A}^{+}$, equipped with the weak* topology, and let $\mathcal{M}_{\sigma}$ denote the subspace consisting of $\sigma$-invariant probability measures. For $\mu \in \mathcal{M}_{\sigma}$, write $h(\mu)$ for the measure theoretic entropy of $\mu$. There is a unique measure $\mu_{0} \in \mathcal{M}_{\sigma}$, called the measure of maximal entropy, for which

$$
h\left(\mu_{0}\right)=\sup _{\mu \in \mathcal{M}_{\sigma}} h(\mu)
$$

and this value coincides with the topological entropy $h(\sigma)$.
For a continuous function $f: \Sigma_{A}^{+} \rightarrow \mathbb{R}$, we define the pressure $P(f)$ of $f$ by the formula

$$
\begin{equation*}
P(f)=\sup _{\mu \in \mathcal{M}_{\sigma}}\left(h(\mu)+\int f d \mu\right) \tag{1.1}
\end{equation*}
$$

and call any measure for which the supremum is attained an equilibrium state for $f$. If $f$ is Hölder continuous (i.e. there exists $\alpha>0$ and $C(f, \alpha) \geq 0$ such that $|f(x)-f(y)| \leq$ $C(f, \alpha) d(x, y)^{\alpha}$, for all $\left.x, y \in \Sigma_{A}^{+}\right)$, then $f$ has a unique equilibrium state which we denote by $\mu_{f}$. The latter is fully supported and $h\left(\mu_{f}\right)>0$. The equilibrium state of the zero function is the measure of maximal entropy, so this is consistent with our earlier notation. The pressure of $f$ also has the following characterization in terms of periodic points:

$$
P(f)=\lim _{n \rightarrow+\infty} \frac{1}{n} \log \sum_{\sigma^{n} x=x} e^{f^{n}(x)}
$$

We say that two functions $f, g: \Sigma_{A}^{+} \rightarrow \mathbb{R}$ are (continuously) cohomologous if there is a continuous function $u: \Sigma_{A}^{+} \rightarrow \mathbb{R}$ such that $f=g+u \circ \sigma-u$. The cohomology class of a Hölder continuous function is determined by its values around periodic orbits. More precisely, writing $f^{n}=f+f \circ \sigma+\cdots+f \circ \sigma^{n-1}$, two Hölder continuous functions $f, g: \Sigma_{A}^{+} \rightarrow \mathbb{R}$ are cohomologous if and only if $f^{n}(x)=g^{n}(x)$ whenever $\sigma^{n} x=x$.

If $f$ and $g$ are cohomologous then $P(f)=P(g)$ and if $c$ is a real number then $P(f+c)=$ $P(f)+c$. Now suppose that $f$ is Hölder continuous and, for $t \in \mathbb{R}$, consider the function $t \mapsto P(t f)$ This function is convex and real analytic and

$$
P^{\prime}(t f)=\int f d \mu_{t f}
$$

Furthermore, if $f$ is not cohomologous to a constant then $P(t f)$ is strictly convex and $P^{\prime \prime}(t f)>0$ everywhere. (If $f$ is cohomologous to a constant $c$ then $P(t f)=h(\sigma)+t c$.)

Suppose that a Hölder continuous function $r: \Sigma_{A}^{+} \rightarrow \mathbb{R}$ is cohomologous to a strictly positive function. If $\delta>0$ satisfies $P(-\delta r)=0$ then, since $\mu_{-\delta r}$ attains the supremum in (1.1), we have the relation

$$
\begin{equation*}
\delta=\frac{h\left(\mu_{-\delta r}\right)}{\int r d \mu_{-\delta r}} \tag{1.2}
\end{equation*}
$$

For any continuous function $f: \Sigma_{A}^{+} \rightarrow \mathbb{R}$, we have

$$
I_{f}:=\left\{\int f d \mu: \mu \in \mathcal{M}_{\sigma}\right\}=\overline{\left\{\frac{f^{n}(x)}{n}: \sigma^{n} x=x\right\}}
$$

and $I_{f}$ is a closed interval. If $f$ is Hölder continuous then

$$
\operatorname{int}\left(I_{f}\right)=\left\{\int f d \mu_{t f}: t \in \mathbb{R}\right\}
$$

In particular, if $0 \in \operatorname{int}\left(I_{f}\right)$ then there exists a unique $\xi \in \mathbb{R}$ such that $\int f d \mu_{\xi f}=0$. Furthermore,

$$
P(\xi f)=h\left(\mu_{\xi f}\right)=\sup \left\{h(\mu): \mu \in \mathcal{M}_{\sigma} \text { and } \int f d \mu=0\right\}
$$

and $\mu_{\xi f}$ is the only measure for which the supremum is realized.
We say that a function $f: \Sigma_{A}^{+} \rightarrow \mathbb{R}$ is locally constant if there exists $N \geq 0$ such that if $x=\left(x_{n}\right)_{n=0}^{\infty}, y=\left(x_{n}\right)_{n=0}^{\infty}$ have $x_{n}=y_{n}$ for all $n \geq N$ then $f(x)=f(y)$. In other words, $f$ may be regarded as a function on $W_{A}^{(N)}$. Clearly, if $f$ is locally constant then $f$ is Hölder continuous (for any choice of exponent $\alpha>0$ ).

We may also regard such a locally constant function as a function $f: \Sigma_{A_{N}}^{+} \rightarrow \mathbb{R}$, which depend on only one co-ordinate. As we noted above, there is a natural correspondence between periodic points for $\sigma: \Sigma_{A}^{+} \rightarrow \Sigma_{A}^{+}$and $\sigma: \Sigma_{A_{N}}^{+} \rightarrow \Sigma_{A_{N}}^{+}$. One easily sees that the value of $f^{n}(x)$ is the same for corresponding periodic orbits for the two shift maps.

If $f$ is locally constant then $P(t f)$ may be given in terms of a matrix. For $t \in \mathbb{R}$, define a $W_{A}^{(N)} \times W_{A}^{(N)}$ matrix $A_{f}(t)$ by

$$
A_{f}(t)(\mathbf{x}, \mathbf{y})=A_{N}(\mathbf{x}, \mathbf{y}) e^{t f(\mathbf{x})}
$$

where $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{N-1}\right), \mathbf{y}=\left(y_{0}, y_{1}, \ldots, y_{N-1}\right) \in W_{A}^{(N)}$. One can easily see that $A_{f}(t)$ is non-negative and aperiodic and thus has a simple positive eigenvalue which is strictly greater in modulus that all the other eigenvalues of $A_{f}(t)$; in fact, this eigenvalue is equal to $e^{P(t f)}$.

## 2. Shifts and Free Groups

In this section we shall consider the shift of finite type $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$associated to the free group $F$ and free basis $\mathcal{A}=\left\{a_{1}, \ldots, a_{k}\right\}$. Setting $\mathcal{I}=\mathcal{A} \cup \mathcal{A}^{-1}$, it is clear that $\Sigma^{+}=\Sigma_{A}^{+}$, where $A(i, j)=1$ unless $j$ is the inverse of $i$, and that $A$ is aperiodic. We may also think of $\Sigma^{+}$as the space of infinite reduced words in $\mathcal{A} \cup \mathcal{A}^{-1}$ and, for a free group, this may be identified with the Gromov boundary of $F$. Furthermore, $W^{(n)}=W_{A}^{(n)}$ may be identified with the set $\{x \in F:|x|=n\}$.

A simple calculation shows that the topological entropy of $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$is $h(\sigma)=$ $\log (2 k-1)$ and the measure of maximal entropy is the measure $\mu_{0}$ defined in the introduction.

Recall that $\mathcal{C}(F)$ denotes the set of non-trivial conjugacy classes in $F$. A conjugacy class $w \in \mathcal{C}(F)$ contains a cyclically reduced word in $\mathcal{A} \cup \mathcal{A}^{-1}$, i.e., a reduced word $x_{0} x_{1} \cdots x_{n-1}$ such that $x_{n-1} \neq x_{0}^{-1}$. The only other cyclically reduced elements of $w$ are obtained from this by cyclic permutation (and also have word length $n$ ) and non-cyclically reduced elements of $w$ have word length greater than $n$. Therefore it is natural to define the length of $w$ (with respect to $\mathcal{A}$ ) to be

$$
|w|:=n=\min _{x \in w}|x| .
$$

It is immediate from the definition that, for $m \geq 1$,

$$
\left|w^{m}\right|=m|w|
$$

where $w^{m}$ is the conjugacy class $\left\{x^{m}: x \in w\right\}$. Furthermore, it is clear from the preceding discussion that there is a natural bijection between $\mathcal{C}(F)$ and the set of periodic orbits of $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$(and between primitive conjugacy classes and prime periodic orbits), such that, if $x, \sigma x, \ldots, \sigma^{n-1} x\left(\sigma^{n} x=x\right)$ corresponds to $w \in \mathcal{C}(F)$ then $|w|=n$.

In order to represent elements of $F$ as elements of a shift space, it is convenient to augment $\Sigma^{+}$by adding an extra "dummy" symbol 0 . Introduce a square matrix $A_{0}$, with rows indexed by $\mathcal{A} \cup \mathcal{A}^{-1} \cup\{0\}$, such that

$$
A_{0}(i, j)= \begin{cases}A(i, j) & \text { if } i, j \in \mathcal{A} \cup \mathcal{A}^{-1} \\ 1 & \text { if } i \in \mathcal{A} \cup \mathcal{A}^{-1} \cup\{0\} \text { and } j=0 \\ 0 & \text { if } i=0 \text { and } j \in \mathcal{A} \cup \mathcal{A}^{-1}\end{cases}
$$

We then define $\Sigma_{0}^{+}=\Sigma_{A_{0}}^{+}$. In other words, an element of $\Sigma_{0}^{+}$is either an element of $\Sigma^{+}$ or an element of $W^{(n)}$, for some $n \geq 0$, followed by an infinite string of 0 s, which we shall denote by $\dot{0}$. (Of course, this procedure does not introduce any extra periodic points.)

Let $\phi: F \rightarrow F$ be an automorphism. We wish to encode the quantity $|\phi(\cdot)|$ in terms of a function $r: \Sigma^{+} \rightarrow \mathbb{R}$. The following lemma is easily seen.

Lemma 2.1. Let $S$ be any finite subset of $F$. Then there exists an integer $M \geq 1$ such that for any reduced word $y_{0} y_{1} \cdots y_{m-1}, m \geq M$, and any $s \in S$

$$
\left|s y_{0} y_{1} \cdots y_{m-1}\right|-m=\left|s y_{0} y_{1} \cdots y_{M-1}\right|-M
$$

Proof. Let $M=\max \{|s|: s \in S\}$ and suppose $m \geq M$. Then $s y_{0} y_{1} \cdots y_{M-1}$ may be written as a reduced word $s^{\prime}$ such that $s^{\prime} y_{M} \cdots y_{m-1}$ is also a reduced word and, in particular, $\left|\left(s y_{0} y_{1} \cdots y_{M-1}\right) y_{M} \cdots y_{m-1}\right|=\left|s y_{0} y_{1} \cdots y_{M-1}\right|+m-M$.

Noting that there exists $C>1$ such that, for all $x \in F$,

$$
C^{-1}|x| \leq|\phi(x)| \leq C|x|
$$

the following may be deduced from Lemma 2.1. (It may be compared with Proposition 4 of [16].)
Lemma 2.2. There exists an integer $N \geq 1$ such that if $x_{0} x_{1} \cdots x_{n-1}$ is a reduced word and $n \geq N$ then

$$
\left|\phi\left(x_{0} x_{1} \cdots x_{n-1}\right)\right|-\left|\phi\left(x_{1} \cdots x_{n-1}\right)\right|=\left|\phi\left(x_{0} x_{1} \cdots x_{N-1}\right)\right|-\left|\phi\left(x_{1} \cdots x_{N-1}\right)\right| .
$$

Proof. We shall apply Lemma 2.1. Let $S=\left\{\phi\left(a_{1}\right)^{ \pm 1}, \ldots, \phi\left(a_{k}\right)^{ \pm 1}\right\}$ and, as before, $M=$ $\max \{|s|: s \in S\}$. For each $p \geq 1$, we have $\left|\phi\left(x_{1} \cdots x_{p}\right)\right| \geq C^{-1} p$. Choose $N=[C M]+1$; then $n \geq N$ implies $\left|\phi\left(x_{1} \cdots x_{n-1}\right)\right| \geq C^{-1}[C M] \geq M$, from which the result follows.

We now introduce a function on $\Sigma_{0}^{+}$which the values $|\phi(x)|, x \in F$, may be recovered. Define $r: \Sigma_{0}^{+} \rightarrow \mathbb{R}$ by

$$
r\left(\left(x_{n}\right)_{n=0}^{\infty}\right)=\left|\phi\left(x_{0} \cdots x_{N-1}\right)\right|-\left|\phi\left(x_{1} \cdots x_{N-1}\right)\right| .
$$

The following result is immediate from the definition.
Lemma 2.3. Suppose that $x_{0} x_{1} \cdots x_{n-1}$ is reduced word. Then

$$
\left|\phi\left(x_{0} x_{1} \cdots x_{n-1}\right)\right|=r^{n}\left(x_{0}, x_{1}, \cdots, x_{n-1}, \dot{0}\right)
$$

Remark. This construction has been used by Lalley [10], Bourdon [2] and Pollicott and Sharp [17],[18]. Apparently, it goes back to Eichler [4].

Clearly, $r$ is a locally constant function. To make this explicit, for $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{N-1}\right)$ $\in W^{(N)}$, we also write

$$
r(\mathbf{x})=\left|\phi\left(x_{0} \cdots x_{N-1}\right)\right|-\left|\phi\left(x_{1} \cdots x_{N-1}\right)\right|
$$

and, as above, define

$$
A_{r}(t)(\mathbf{x}, \mathbf{y})=A_{N}(\mathbf{x}, \mathbf{y}) e^{\operatorname{tr}(\mathbf{x})}
$$

Define a function $\mathfrak{p}: \mathbb{R} \rightarrow \mathbb{R}$ by $\mathfrak{p}(t)=P(t r)$ and recall that $e^{\mathfrak{p}(t)}$ may be characterized as the (simple) maximal eigenvalue of $A_{r}(t)$. We know that $\mathfrak{p}(t)$ is a convex real-analytic function of $t$. We define an associated concave function $\mathfrak{h}: \operatorname{int}\left(I_{r}\right) \rightarrow \mathbb{R}$ by

$$
\mathfrak{h}(\rho)=\sup \left\{h(\mu): \mu \in \mathcal{M}_{\sigma} \text { and } \int r d \mu=\rho\right\} .
$$

Recalling that $\operatorname{int}\left(I_{r}\right)=\left\{\mathfrak{p}^{\prime}(t): t \in \mathbb{R}\right\}$, we have $\mathfrak{h}(\rho)=\mathfrak{p}(\xi)-\xi \rho$, where $\xi \in \mathbb{R}$ is chosen to be the unique value with $\mathfrak{p}^{\prime}(\xi)=\rho$ and, in particular, $\mathfrak{h}$ is real-analytic. (In the language of convex analysis, $-\mathfrak{h}$ is the Legendre transform of $\mathfrak{p}$ [19].)

The next result relates the generic stretch $\lambda(\phi)$ to the functions $\mathfrak{p}(t)$ and $\mathfrak{h}(\rho)$.
Theorem 1. Suppose that $\phi$ is not simple.
(i)

$$
\lambda(\phi)=\int r d \mu_{0}=\mathfrak{p}^{\prime}(0) .
$$

(ii) For $\rho \in \operatorname{int}\left(I_{r}\right), 0<\mathfrak{h}(\rho) \leq 2 k-1$ and $\mathfrak{h}(\rho)=2 k-1$ if and only if $\rho=\lambda(\phi)$.

Proof.
(i) For $x=\left(x_{n}\right)_{n=0}^{\infty} \in \Sigma^{+}$,

$$
r^{n}(x)=r^{n}\left(x_{0}, x_{1}, \cdots, x_{n-1}, \dot{0}\right)+O(1)
$$

Thus, by Lemma 2.3,

$$
\lim _{n \rightarrow+\infty} \frac{r^{n}(x)}{n}=\lim _{n \rightarrow+\infty} \frac{\left|\phi\left(x_{0} x_{1} \cdots x_{n-1}\right)\right|}{n}
$$

(provided the limit exists). By the ergodic theorem, the Left Hand Side exists $\mu_{0}-$ a.e. and the limit is equal to $\int r d \mu_{0}$.
(ii) It is clear that $\mathfrak{h}(\rho) \leq 2 k-1$. Using the formula $\mathfrak{h}(\rho)=\mathfrak{p}(\xi)-\xi \rho$ and the definition of $\mathfrak{p}$, we see that $\mathfrak{h}(\rho)=h\left(\mu_{\xi r}\right)>0$. By part (i), we have $\mathfrak{h}(\lambda(\phi))=2 k-1$ while, for $\rho \neq \lambda(\phi), \xi \neq 0$, so $\mathfrak{h}(\rho)<2 k-1$.

## 3. Conjugacy Distortion Spectrum

The construction in the preceding section is particularly well suited to studying the action of $\phi$ on conjugacy classes. As shown in the next lemma, the length of the resulting conjugacy classes is simply obtained by summing $r$ around the corresponding periodic orbit. Fortunately, there is a well developed theory of periodic orbit sums which we will then able to use.

Lemma 3.1. Let $\sigma^{n} x=x$ correspond to the conjugacy class $w$ containing the cyclically reduced word $x_{0} x_{1} \cdots x_{n-1}$. Then

$$
|\phi(w)|=r^{n}(x) .
$$

Proof. Let $x^{(m)}$ denote the reduced word obtained from the $m$-fold concatenation of $x_{0} x_{1} \cdots x_{n-1}$. Since $r$ is locally constant, there exists $N \geq 1$ such that

$$
\left|r^{m n}(x)-r^{m n}\left(x^{(m)}, 0,0, \ldots\right)\right| \leq 2 N\|r\|_{\infty}
$$

Noting that $r^{m n}(x)=m r^{n}(x)$, the above estimate gives us that

$$
r^{n}(x)=\lim _{m \rightarrow+\infty} \frac{1}{m} r^{m n}\left(x^{(m)}, 0,0, \ldots\right)=\lim _{m \rightarrow+\infty} \frac{1}{m}\left|\phi\left(x^{(m)}\right)\right|
$$

Thus it remains to show that this last quantity is equal to $|\phi(w)|$. We shall do this by proving inequalities in both directions.

First observe that $x^{(m)}$ is a cyclically reduced word in $w^{m}$. Therefore

$$
m|\phi(w)|=\left|\phi(w)^{m}\right|=\left|\phi\left(w^{m}\right)\right| \leq\left|\phi\left(x^{(m)}\right)\right|
$$

and so $|\phi(w)| \leq \lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(x^{(m)}\right)\right|$.
Now suppose that $v \in w$. It is clear that

$$
\lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(v^{m}\right)\right|=\lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(x^{(m)}\right)\right| .
$$

On the other hand, since $\left\{\left|\phi\left(v^{m}\right)\right|\right\}_{m \geq 1}$ is a subadditive sequence,

$$
\lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(v^{m}\right)\right|=\inf _{m \geq 1} m^{-1}\left|\phi\left(v^{m}\right)\right|
$$

and so $\lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(x^{(m)}\right)\right| \leq|\phi(v)|$. Hence

$$
\lim _{m \rightarrow+\infty} m^{-1}\left|\phi\left(v^{m}\right)\right| \leq \inf _{v \in w}|\phi(v)|=|\phi(w)| .
$$

Suppose that $w \in \mathcal{C}(F)$. If $x, y \in w$ then $\phi(x), \phi(y)$ are conjugate, so the conjugacy class $\phi(w)$ is well defined. As in the introduction, we define the conjugacy distortion spectrum of $\phi$ to be the set

$$
\mathcal{D}_{\phi}=\left\{\frac{|\phi(w)|}{|w|}: w \in \mathcal{C}(F)\right\}
$$

and let $\overline{\mathcal{D}_{\phi}}$ be the closure of $\mathcal{D}_{\phi}$. The structure of $\mathcal{D}_{\phi}$ was studied by Kapovich [8], who showed the following.

## Proposition 3.1 [8].

(i) $\overline{\mathcal{D}_{\phi}}$ is a closed interval with rational endpoints and $\mathcal{D}_{\phi}=\overline{\mathcal{D}_{\phi}} \cap \mathbb{Q}$. (In other words, if $\rho \in \overline{\mathcal{D}_{\phi}}$ is rational then there exists $w \in \mathcal{C}(F)$ such that $|\phi(w)|=\rho|w|$.)
(ii) If $\phi$ is not simple then $\overline{\mathcal{D}_{\phi}}$ has non-empty interior which contains 1 and $\lambda(\phi)$.

We shall strengthen Kapovich's result by showing that the set conjugacy classes with a given distortion has exponential growth and that the growth rate is given by the function $\mathfrak{h}$.
Theorem 2. Suppose that $\phi$ is not simple.
(i) $\overline{\mathcal{D}_{\phi}}=I_{r}$.
(ii) If $\rho \in \overline{\mathcal{D}_{\phi}}$ is rational then

$$
\limsup _{n \rightarrow+\infty} \frac{1}{n} \log \#\left\{w \in \mathcal{C}(F):|w|=n, \frac{|\phi(w)|}{|w|}=\rho\right\}=\mathfrak{h}(\rho)
$$

(iii) If $\rho \in \operatorname{int}\left(\overline{\mathcal{D}_{\phi}}\right)$ is rational then there exists an integer $d(\rho) \geq 1$ and a constant $C(\rho)>0$ such that

$$
\#\left\{w \in \mathcal{C}(F):|w|=d(r) n, \frac{|\phi(w)|}{|w|}=\rho\right\} \sim C(\rho) \frac{e^{d(\rho) n \mathfrak{h}(\rho)}}{\sqrt{d(r) n}}, \quad \text { as } n \rightarrow+\infty
$$

Corollary. For each $\rho$, the set $\{w \in \mathcal{C}(F):|\phi(w)| /|w|=\rho\}$ has zero density in in $\mathcal{C}(F)$, i.e.,

$$
\lim _{n \rightarrow+\infty} \frac{\#\{w \in \mathcal{C}(F):|w|=n,|\phi(w)| /|w|=\rho\}}{\#\{w \in \mathcal{C}(F):|w|=n\}}=0
$$

We shall prove this theorem by rephrasing it in terms of periodic points for the shift map $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$(or, more precisely, $\sigma: \Sigma_{A_{N}}^{+} \rightarrow \Sigma_{A_{N}}^{+}$). In view of the correspondence between $\mathcal{C}(F)$ and periodic points for $\sigma$ and Lemma 3.1, we have that

$$
\left\{w \in \mathcal{C}(F):|w|=n, \frac{|\phi(w)|}{|w|}=\rho\right\}=\left\{x \in \operatorname{Fix}_{n}: \frac{r^{n}(x)}{n}=\rho\right\}
$$

The theorem will follow from the main result of [15], stated in the lemma below, once we have checked the hypotheses.
Lemma 3.2 [15]. Let $A$ be an aperiodic zero-one matrix. Suppose that $f: \Sigma_{A}^{+} \rightarrow \mathbb{Z}$ is a locally constant function depending on two coordinates, i.e., $f(x)=f\left(x_{0}, x_{1}\right)$, satisfying the following conditions:
(1) there exists $\xi \in \mathbb{R}$ such that

$$
\int f d \mu_{\xi f}=0
$$

(2) $\bigcup_{n \geq 1}\left\{f^{n}(x): \sigma^{n} x=x\right\}$ is not contained in any proper subgroup of $\mathbb{Z}$.

Then there exists an integer $d(f) \geq 1$ and a constant $C(f)>0$ such that

$$
\left\{x \in F i x_{d n}: f^{d n}(x)=0\right\} \sim C(f) \frac{e^{d(f) n h\left(\mu_{\xi f}\right)}}{\sqrt{d(f) n}}, \quad \text { as } n \rightarrow+\infty
$$

Remark. The requirement that $f$ depends on only two coordinates is only a technical simplification and may be easily removed.
Proof of Theorem 2. Write $\rho=p / q$ in lowest terms. Note that $r^{n}(x) / n=\rho$ if and only if $(q r-p)^{n}(x)=q r^{n}(x)-n p=0$. Thus we will apply the above result to the function $f_{\rho}=q r-p$. We shall show that $f_{\rho}$ satisfies conditions (1) and (2).

First observe that $I_{f_{\rho}}=q I_{r}-p$. We know that $p / q \in \operatorname{int}\left(I_{r}\right)$, so $0 \in \operatorname{int}\left(I_{f_{\rho}}\right)$. Hence condition (1) is satisfied. Furthermore, as $I_{f_{\rho}}$ has interior, $f_{\rho}$ is not cohomologous to a constant (necessarily zero), so $\bigcup_{n \geq 1}\left\{f^{n}(x): \sigma^{n} x=x\right\}$ generates a group $a \mathbb{Z}$, where $a \geq 1$ is an integer. Replacing $f_{\rho}$ by $f_{\rho} / a$, condition (2) is satisfied. This proves part (iii) and hence part (ii) of Theorem 2 .

Remark. One may read off upper and lower bounds on the number of $w \in \mathcal{C}(F)$ with $|w|=n$ and $|\phi(w)| /|w|=\rho$ from the paper [11]. One has

$$
\#\left\{w \in \mathcal{C}(F):|w|=n, \frac{|\phi(w)|}{|w|}=\rho\right\} \leq(n+1)^{2 k(2 k-1)^{N}+1} e^{n \mathfrak{h}(\rho)}
$$

The lower bound is more involved. Given $\epsilon>0$, there exists $\delta>0$ and $m \in \mathbb{N}$ such that, if $d(\rho) \mid n$ and $n$ is sufficiently large then

$$
\#\left\{w \in \mathcal{C}(F):|w|=n, \frac{|\phi(w)|}{|w|}=\rho\right\} \geq \delta n^{-m} e^{n(\mathfrak{h}(\rho)-\epsilon)}
$$

## 4. Counting Group Elements

Our aim in the next two sections is to relate the number $\operatorname{Curl}(\phi)$ to the function $\mathfrak{h}$. This is complicated by the fact that $\operatorname{Curl}(\phi)$ is defined in terms of the action of $\phi$ on elements of $F$ rather than conjugacy classes. Group elements do not have such a nice correspondence to elements of $\Sigma^{+}$as occurs in the correspondence between conjugacy classes and periodic orbits for $\sigma: \Sigma^{+} \rightarrow \Sigma^{+}$. However, if we consider the larger shift space $\Sigma_{0}^{+}$then we may identify $F$ with a set of pre-images of a point under $\sigma: \Sigma_{0}^{+} \rightarrow \Sigma_{0}^{+}$.

Suppose that $\phi$ is not simple. As in the previous section, write $f=f_{1}=r-1$ and define $\xi$ by

$$
h\left(\mu_{\xi f}\right)=\mathfrak{h}(1)=\sup \left\{h(\mu): \mu \in \mathcal{M}_{\sigma} \text { and } \int f d \mu=0\right\} .
$$

We know $f$ is a locally constant function and may be regarded as a function on $W_{A_{0}}^{(N)}$. To avoid too many subscripts, we shall write $B=A_{0}$. As in section 1 , we write

$$
B_{f}(\xi+i t)(\mathbf{x}, \mathbf{y})=B_{N}(\mathbf{x}, \mathbf{y}) e^{(\xi+i t) f(\mathbf{x})}
$$

where $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{N-1}\right), \mathbf{y}=\left(y_{0}, y_{1}, \ldots, y_{N-1}\right) \in W_{B}^{(N)}$. (Note that the argument is now $\xi+i t$.) The matrices $A_{f}(\xi+i t)$ are defined similarly and have the same non-zero spectrum as $B_{f}(\xi+i t)$. In particular, $B_{f}(\xi)$ has a simple maximal eigenvalue equal to $e^{P(\xi f)}=e^{\mathfrak{p}(\xi)-\xi}=e^{\mathfrak{h}(1)}$. For small values of $|t|$, this eigenvalue persists and $B_{f}(\xi+i t)$ has a simple eigenvalue $\beta(t)$, depending analytically on $t$, with $\beta(0)=e^{\mathfrak{h}(1)}$.

A simple calculation shows that

$$
\begin{equation*}
\#\{x \in F:|x|=n\}=\sum_{\mathbf{x} \in W_{A}^{(N)}} \sum_{a \in \mathcal{A} \cup \mathcal{A}^{-1}} B^{n}(\mathbf{x}, \mathbf{a}), \tag{4.1}
\end{equation*}
$$

where $\mathbf{a}=(a, 0, \ldots, 0)$.
Lemma 4.1.

$$
\begin{aligned}
\#\{x \in F:|\phi(x)|=|x|=n\} & =\sum_{|x|=n} \frac{1}{2 \pi} \int_{-\pi}^{\pi} e^{(\xi+i t)(|\phi(x)|-|x|)} d t \\
& =\sum_{\mathbf{x} \in W_{A}^{(N)}} \sum_{a \in \mathcal{A} \cup \mathcal{A}^{-1}} \frac{1}{2 \pi} \int_{-\pi}^{\pi} B_{f}(\xi+i t)(\mathbf{x}, \mathbf{a}) d t
\end{aligned}
$$

Proof. By orthogonality,

$$
\frac{1}{2 \pi} \int_{-\pi}^{\pi} e^{i t(|\phi(x)|-|x|)} d t= \begin{cases}1 & \text { if }|\phi(x)|=|x| \\ 0 & \text { otherwise. }\end{cases}
$$

In the former case, the factor $e^{\xi(|\phi(x)|-|x|)}=1$, so the first equality holds. The second follows from Lemma 2.3 and (4.1).

This may be handled as in [22, pp.899-900] (where the roles of $A$ and $B$ are switched) to show the following.

Proposition 4.1. There exists an integer $d \geq 1$ and a real number $C>0$ such that

$$
\#\{x \in F:|\phi(x)|=|x|=d n\} \sim C \frac{e^{d n \mathfrak{h}(1)}}{\sqrt{n}}, \quad \text { as } n \rightarrow+\infty
$$

## Corollary.

$$
\operatorname{Curl}(\phi) \geq \frac{e^{\mathfrak{h}(1)}}{2 k-1}
$$

Proof. Since

$$
\#\{x \in F:|\phi(x)|=|x|=d n\} \leq \#\{x \in F:|x| \leq d n,|\phi(x)| \leq d n\}
$$

this follows from Proposition 4.1 and the definition of $\operatorname{Curl}(\phi)$.

## 5. The Formula for $\operatorname{Curl}(\phi)$

We now establish our formula for $\operatorname{Curl}(\phi)$ by obtaining an upper bound. It is clear that, for each $n \geq 1$ and $0 \leq p \leq 1$,

$$
\#\{x \in F:|x| \leq n,|\phi(x)| \leq n\} \leq \#\{x \in F:(1-p)|x|+p \mid \phi(x) \leq n\}
$$

so

$$
\begin{equation*}
\operatorname{Curl}(\phi) \leq \frac{1}{2 k-1} \inf _{0 \leq p \leq 1}\left(\limsup _{n \rightarrow+\infty}(\#\{x \in F:(1-p)|x|+p|\phi(x)| \leq n\})^{1 / n}\right) \tag{5.1}
\end{equation*}
$$

We shall show that there exists a value of $p$ for which the above limsup is equal to $\mathfrak{h}(1)$. Combining this with the corollary to Proposition 4.1 will give the desired formula.

For $0 \leq p \leq 1$, write

$$
\alpha(p)=\limsup _{n \rightarrow+\infty} \frac{1}{n} \#\{x \in F:(1-p)|x|+p|\phi(x)| \leq n\}
$$

Then $\alpha(p)$ is the abscissa of convergence of the Dirichlet series

$$
\sum_{x \in F} e^{-t((1-p)|x|+p|\phi(x)|)}
$$

To analyse this series, define locally constant functions $r_{p}: \Sigma_{0}^{+} \rightarrow \mathbb{R}$ by $r_{p}=(1-p)+p r$ and introduce matrices

$$
B_{p}(t):=B_{r_{p}}(t)=B_{N}(\mathbf{x}, \mathbf{y}) e^{t r_{p}(\mathbf{x})}
$$

Then

$$
\sum_{x \in F} e^{-t((1-p)|x|+p|\phi(x)|)}=\sum_{n=0}^{\infty} \sum_{\mathbf{x} \in W_{A}^{(N)}} \sum_{a \in \mathcal{A} \cup \mathcal{A}^{-1}} B_{p}^{n}(-t)(\mathbf{x}, \mathbf{a}),
$$

so $\alpha(p)$ is the value of $t$ for which $B_{p}(-t)$ has spectral radius equal to one. Thus, $\alpha(p)$ is determined by the equation

$$
P\left(-\alpha(p) r_{p}\right)=0
$$

which is equivalent to

$$
\sup _{m \in \mathcal{M}_{\sigma}}\left(h(m)-\alpha(p) \int r_{p} d m\right)=0
$$

or

$$
\begin{equation*}
\alpha(p)(1-p)=\sup _{m \in \mathcal{M}_{\sigma}}\left(h(m)-\alpha(p) p \int r d m\right) \tag{5.2}
\end{equation*}
$$

The supremum in (5.2) is uniquely attained at $\mu_{-\alpha(p) p r}$, the equilibrium state of $-\alpha(p) p r$.

Lemma 5.1. There exists $0<p^{*}<1$ such that

$$
\int r d \mu_{\alpha\left(p^{*}\right) p^{*} r}=1
$$

Proof. First note that $R(p):=\int r d \mu_{\alpha(p) p r}$ depends continuously on $p$. We shall show that $R(0)>1$ and $R(1)<1$; hence, by the Intermediate Value Theorem, there will be a value $p^{*}$ with $0<p^{*}<1$ and $R\left(p^{*}\right)=1$.

The first inequality is straightforward since

$$
R(0)=\int r d \mu_{0}=\lambda(\phi)>1
$$

Since $|\phi(x)|=|x|_{\phi^{-1}(\mathcal{A})}$, one has $\alpha(1)=\log (2 k-1)=h(\sigma)$. On the other hand, by equation (1.2),

$$
\alpha(1)=\frac{h\left(\mu_{-\alpha(1) r}\right)}{\int r d \mu_{-\alpha(1) r}}=\frac{h\left(\mu_{-h(\sigma) r}\right)}{R(1)} .
$$

Since $r$ is not cohomologous to a constant, $\mu_{-h(\sigma) r} \neq \mu_{0}$, so $h\left(\mu_{-h(\sigma) r}\right)<h(\sigma)$. Consequently, $R(1)<1$.

Consider (5.2) for the value $p^{*}$. Since the supremum is attained at $\mu_{-\alpha\left(p^{*}\right) p^{*} r}$ and $\int r d \mu_{-\alpha\left(p^{*}\right) p^{*} r}=1$, we may rewrite (5.2) as

$$
\begin{aligned}
\alpha\left(p^{*}\right)\left(1-p^{*}\right) & =\sup \left\{h(m)-\alpha\left(p^{*}\right) p^{*} \int r d m: m \in \mathcal{M}_{\sigma}, \int r d m=1\right\} \\
& =\sup \left\{h(m): m \in \mathcal{M}_{\sigma}, \int r d m=1\right\}-\alpha\left(p^{*}\right) p^{*} \\
& =\mathfrak{h}(1)-\alpha\left(p^{*}\right) p^{*}
\end{aligned}
$$

The terms $-\alpha\left(p^{*}\right) p^{*}$ cancel on each side to give

$$
\alpha\left(p^{*}\right)=\mathfrak{h}(1) .
$$

Substituting this in inequality (5.1) gives

$$
\operatorname{Curl}(\phi) \leq \frac{e^{\mathfrak{h}(1)}}{2 k-1}
$$

Combining this with the corollary to Proposition 4.1 gives the following theorem.
Theorem 3. If $\phi$ is not simple then

$$
\operatorname{Curl}(\phi)=\frac{e^{\mathfrak{h}(1)}}{2 k-1}
$$

## 6. A Manhattan Curve

In this final section, we draw a parallel between the functions introduced in this paper and Burger's "Manhattan Curve", associated to a pair of hyperbolic Riemann surfaces [3].

We recall Burger's definition. Let $\Sigma_{1}$ and $\Sigma_{2}$ be two compact hyperbolic Riemann surfaces, which are assumed to be homeomorphic, and let $\psi: \Sigma_{1} \rightarrow \Sigma_{2}$ be a homeomorphism between them. In particular, $\psi$ allow us to identify free homotopy classes on the surfaces. For a free homotopy class $\gamma$, we write $l_{1}(\gamma)$ and $l_{2}(\gamma)$, respectively, for the lengths of the corresponding closed geodesics on the two surfaces.

The Manhattan curve $\mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ is defined to be the boundary of the convex set

$$
\left\{(a, b) \in \mathbb{R}^{2}: \sum_{\gamma} e^{-a l_{1}(\gamma)-b l_{2}(\gamma)}<+\infty\right\}
$$

This was studied in [3] and [21], where the following results were proved.
Proposition 6.1 [3],[21].
(i) $\mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ is a straight line if and only if $\Sigma_{1}$ and $\Sigma_{2}$ are isometric.
(ii) $\mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ is real analytic.
(iii) $\mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ has asymptotes whose normals have slopes equal to the maximum and minimum geodesic stretch between $\Sigma_{1}$ and $\Sigma_{2}$.
(iv) $\mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ passes through $(1,0)$, where its normal has slope equal to the intersection $i\left(\Sigma_{1}, \Sigma_{2}\right)$. (It also passes through $(0,1)$, where its normal has slope equal to $1 / i\left(\Sigma_{2}, \Sigma_{1}\right)$.
(v) There is a unique point $(a, b) \in \mathfrak{M}\left(\Sigma_{1}, \Sigma_{2}\right)$ where the normal has slope 1 and $a+b$ is equal to the correlation number of $\Sigma_{1}$ and $\Sigma_{2}$.

See [1],[3] for the definitions of max and min geodesic stretch and intersection. The correlation number of $\Sigma_{1}$ and $\Sigma_{2}$ is defined to be the exponential growth rate of

$$
\#\left\{\gamma: l_{1}(\gamma), l_{2}(\gamma) \in(T, T+\epsilon)\right\}
$$

for fixed $\epsilon>0$, as $T \rightarrow+\infty$ [20].
In our setting, let us define a Manhattan curve $\mathfrak{M}_{\phi}$, associated to $\phi$, to be the boundary of the set

$$
\left\{(a, b) \in \mathbb{R}^{2}: \sum_{w \in \mathcal{C}(F)} e^{-a|w|-b|\phi(w)|}<+\infty\right\}
$$

By writing this in terms of periodic points for the shift map, this may be described as the set

$$
\left\{(a, b) \in \mathbb{R}^{2}: P(-a-b r)=0\right\}
$$

or, equivalently,

$$
\left\{(a, b) \in \mathbb{R}^{2}: \mathfrak{p}(-b)=a\right\}
$$

Let us define $\mathfrak{q}(s)$ implicitly by

$$
\begin{equation*}
\mathfrak{p}(-\mathfrak{q}(s))=s \tag{6.1}
\end{equation*}
$$

then $\mathfrak{M}_{\phi}$ is the graph of $\mathfrak{q}$. Since $\mathfrak{p}^{\prime}(-t)=-\int r d \mu_{-t r} \neq 0$, the Implicit Function Theorem gives that $\mathfrak{q}$ is real analytic. (This parallels the analysis of [21].)

Let us examine the slope of the normal at a point $(a, b)=(s, \mathfrak{q}(s))$ on $\mathfrak{M}_{\phi}$. At this point, the normal has slope $-1 / \mathfrak{q}^{\prime}(s)$. Now, differentiating (6.1),

$$
1=\frac{d}{d s} \mathfrak{p}(-\mathfrak{q}(s))=\left(-\int r d \mu_{-\mathfrak{q}(s) r}\right) \mathfrak{q}^{\prime}(s)
$$

so the normal to $\mathfrak{M}_{\phi}$ at $(s, \mathfrak{q}(s))$ has slope

$$
\frac{-1}{\mathfrak{q}^{\prime}(s)}=\int r d \mu_{-\mathfrak{q}(s) r} .
$$

Therefore the set of normals to $\mathfrak{M}_{\phi}$ is equal to

$$
\left\{\int r d \mu_{t}: t \in \mathbb{R}\right\}=\operatorname{int}\left(I_{r}\right)=\operatorname{int}\left(\overline{\mathcal{D}_{\phi}}\right)
$$

and it is easy to recover the following statements.

## Theorem 4.

(i) $\mathfrak{M}_{\phi}$ is a straight line if and only if $\phi$ is simple.
(ii) $\mathfrak{M}_{\phi}$ is real analytic.
(iii) $\mathfrak{M}_{\phi}$ has asymptotes whose normals have slopes equal to the $\max \mathcal{D}_{\phi}$ and $\min \mathcal{D}_{\phi}$.
(iv) $\mathfrak{M}_{\phi}$ passes through $(\log (2 k-1), 0)$, where its normal has slope equal to the generic stretch $\lambda(\phi)$
(v) There is a unique point $(a, b) \in \mathfrak{M}_{\phi}$ where the normal has slope 1 and $a+b=\mathfrak{h}(1)$.

Proof. The only statement which requires proof is (v). If the normal has slope 1 at $(a, b)=$ $(s, \mathfrak{q}(s))$ then

$$
\int r d \mu_{-\mathfrak{q}(s) r}=1
$$

and $-\mathfrak{q}(s)=\xi$, where

$$
\mathfrak{h}(1)=\mathfrak{p}(\xi)-\xi=\mathfrak{p}(-\mathfrak{q}(s))+\mathfrak{q}(s)=s+\mathfrak{q}(s)=a+b
$$

Remark. Kaimanovich, Kapovich and Schupp have shown how to identify the generic stretching factor $\lambda(\phi)$ as an intersection of currents [7], [9].

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School of Mathematics, University of Manchester, Oxford Road, Manchester M13 9PL, U.K.

