

PART I: Modeling of Production of Aluminium

PART II: The one-fluid MHD equations - Analysis and Numerics

PART III: The two-fluid MHD equations - Analysis and Numerics

PART IV: Simulation of Aluminium Electrolysis

PART V: Summary

# Production of Aluminium: Modeling, Analysis and Numerics

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... the industrial problem

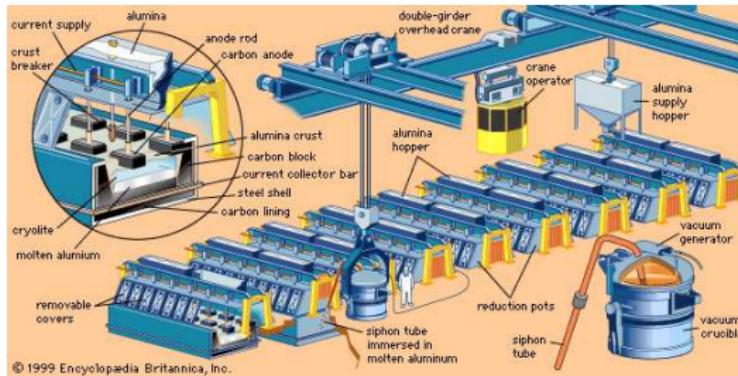
Modeling of the problem

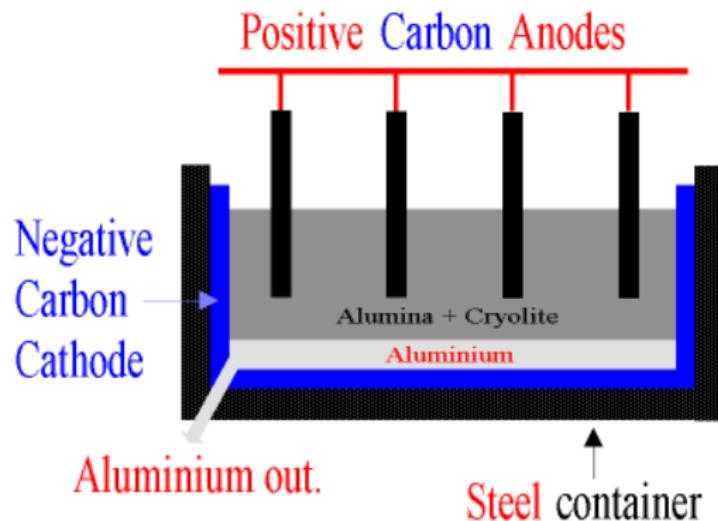
Continuum Model

# Part I: Modeling of Production of Aluminium

# Motivation: Production of Aluminium

- ▶ electrolysis: reduce aluminium oxide to aluminium
- ▶ two non-miscible, conducting, incompressible fluids
- ▶ high temperatures & high currents: no experimental data
- ▶ industrial challenge: reduce electric power waste

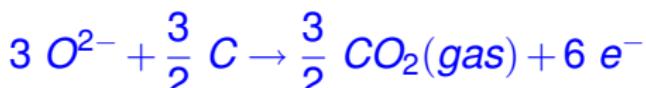




- ▶ anod-metal distance: controlled movement of interface

## ... the industrial problem

- ▶ chem. reaction A: at surface of carbon anod



- ▶ chem. reaction B: at interface between the two fluids



- ▶ Global balance:



- ▶ needed:  $\approx 1000 C^0$ , high intensity of current
- ▶ "the higher intensity, the higher production"

- ▶ Goals/Problem: reduce power waste
  - reduce anod-metal distance: small distance between anod and surface of aluminium layer ('a few centimeters')
  - strong Lorentz forces: motion of the interface; instabilities
  - avoid short -circuits: no touching of fluid-fluid interface and anod
- ▶ Questions: how stabilize the position of the interface?
- ▶ Control parameter: height of anod, intensity of current, geometry
- ▶ experimental observations difficult to obtain

## Modeling of the problem

- ▶ Physical phenomena and simplifying assumptions:
  - magnetohydrodynamics (MHD)
  - moving interface
  - electrochemistry (concentration of chemical species not homogeneous throughout liquids)
  - three-phase flows (bubbles of carbon oxydes at surface of anodes)
  - solidification process at boundaries
  - temperature effects (influence of physical parameters)

# Continuum Model

- ▶ Maxwell equations coupled to multi-fluid Navier-Stokes equations

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0,$$

$$\partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) - \operatorname{div}\left(2\eta \boldsymbol{\varepsilon}(\mathbf{u})\right) + \nabla p = \rho \mathbf{f} + \frac{1}{\mu} \operatorname{curl} \mathbf{B} \times \mathbf{B}$$

$$\operatorname{div} \mathbf{u} = 0 \quad \operatorname{div} \mathbf{B} = 0$$

$$\partial_t \mathbf{B} + \operatorname{curl} \mathbf{E} = 0$$

$$-\partial_t(\varepsilon \mathbf{E}) + \operatorname{curl} \frac{\mathbf{B}}{\mu} = \mathbf{j} \quad \text{for} \quad \mathbf{j} = \sigma(\mathbf{E} + \mathbf{u} \times \mathbf{B}),$$

$$\rho \text{ density of fluids} \quad \boldsymbol{\varepsilon}(\mathbf{u}) = \frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^T)$$

►  $\mathbf{B}$  magnetic field  $\sigma$  electric conductivity

$$\mathbf{j} \text{ electric current} \quad \eta \text{ viscosity}$$

► Assumption: 'low frequency hypothesis': neglect  $\partial_t(\varepsilon \mathbf{E})$ .

$$\partial_t \mathbf{B} + \operatorname{curl} \left( \frac{1}{\mu_0 \sigma} \operatorname{curl} \mathbf{B} \right) = \operatorname{curl} (\mathbf{u} \times \mathbf{B}),$$

$$\operatorname{div} \mathbf{B} = 0$$

$$\Rightarrow \partial_t \mathbf{B} + \operatorname{curl} \left( \frac{1}{\mu_0 \sigma} \operatorname{curl} \mathbf{B} \right) = \operatorname{curl} (\mathbf{u} \times \mathbf{B}),$$

$$\operatorname{div} \mathbf{B} = 0$$

► magnetic boundary conditions on  $\partial\Omega$ :

$$\mathbf{B} \cdot \mathbf{n}, \quad \operatorname{curl} \mathbf{B} \times \mathbf{n}$$

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Analytical results/Numerical Strategies

Numerical Strategies to solve one-fluid MHD equation

## PART II: The one-fluid MHD equations – Analysis and Numerics

# Analytical results/Numerical Strategies

- ▶ M. Sermange & R. Temam [1983]:
  - ▶ global weak solutions for  $\Omega \subset \mathbb{R}^3$
  - ▶ local strong solutions for  $\Omega \subset \mathbb{R}^3$
- ▶ Goal: Develop schemes that approximate weak solutions
- ▶ for simplicity here: only temporal discretization

- Scheme A: Let  $n \geq 1$ . Find  $(\mathbf{u}^n, p^n, \mathbf{b}^n, r^n)$  such that

$$d_t \mathbf{u}^n - \Delta \mathbf{u}^n + (\mathbf{u}^{n-1} \cdot \nabla) \mathbf{u}^n + \frac{1}{2} (\operatorname{div} \mathbf{u}^{n-1}) \mathbf{u}^n$$

$$+ \mathbf{b}^{n-1} \times \operatorname{curl} \mathbf{b}^n + \nabla p^n = \mathbf{g}^n$$

$$\operatorname{div} \mathbf{u}^n = 0 \quad \operatorname{div} \mathbf{b}^n = 0$$

$$d_t \mathbf{b}^n + \operatorname{curl}(\operatorname{curl} \mathbf{b}^n) - \operatorname{curl}(\mathbf{u}^n \times \mathbf{b}^{n-1}) - \nabla r^n = 0$$

# Numerical Strategies to solve one-fluid MHD equation

- discrete energy law: “multiply 1st eqn. by  $\mathbf{u}^n$ , and 3rd eqn. by  $\mathbf{b}^n$ ”

$$\frac{1}{2} d_t \left[ \|\mathbf{u}^n\|^2 + \|\mathbf{b}^n\|^2 \right] + \frac{k}{2} \left[ \|d_t \mathbf{u}^n\|^2 + \|d_t \mathbf{b}^n\|^2 \right] \\ + \|\nabla \mathbf{u}^n\|^2 + \|\mathbf{curl} \mathbf{b}^n\|^2 = (\mathbf{g}^n, \mathbf{u}^n)$$

- ▶ Result [P'08]: Subconvergence ( $k, h \rightarrow 0$ ) to weak solution of one-fluid MHD system
  - drawback of scheme: system coupled
- +/- iterative decoupling strategy: restrictive mesh-constraint  
 $k \leq Ch^4$  needed for convergence

► Scheme B: decoupled scheme

$$\begin{aligned}
 & d_t \mathbf{u}^n - \Delta \mathbf{u}^n + (\mathbf{u}^{n-1} \cdot \nabla) \mathbf{u}^n + \frac{1}{2} (\operatorname{div} \mathbf{u}^{n-1}) \mathbf{u}^n \\
 & \quad + \mathbf{b}^{n-1} \times \operatorname{curl} \mathbf{b}^{n-1} - \nabla p^n = \mathbf{g}^n \\
 & d_t \mathbf{b}^n + \operatorname{curl}(\operatorname{curl} \mathbf{b}^n) - \operatorname{curl}(\mathbf{u}^{n-1} \times \mathbf{b}^{n-1}) - \nabla r^n = 0
 \end{aligned}$$

- Property: A perturbed energy law holds *in a fully discrete setting*, for  $k \leq Ch^3$ .
- Result [P. '08]: Subsequence convergence ( $k, h \rightarrow 0$ ) to weak solution of one-fluid MHD-system.

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The Model

## PART III: The two-fluid MHD equations – Analysis and Numerics

# The Model

- ▶ Navier-Stokes with variable density and viscosity with Maxwell's equation

$$\begin{aligned}
 (\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) - \operatorname{div}(\eta(\rho) \mathbf{D}(\mathbf{u})) &= -\nabla p + \mathbf{g} + \frac{1}{\bar{\mu}} \operatorname{curl} \mathbf{b} \times \mathbf{b}, \\
 \operatorname{div} \mathbf{u} &= 0, \\
 \rho_t + \operatorname{div}(\rho \mathbf{u}) &= 0, \\
 \mathbf{b}_t + \frac{1}{\bar{\mu}} \operatorname{curl} \left( \frac{1}{\xi(\rho)} \operatorname{curl} \mathbf{b} \right) &= \operatorname{curl}(\mathbf{u} \times \mathbf{b}),
 \end{aligned} \tag{1}$$

div  $\mathbf{b}$  = 0,

together with IC's & BC's

► Assumptions:

1.  $0 < \bar{\eta}_- \leq \eta \leq \bar{\eta}_+$ ,  $0 < \bar{\xi}_- \leq \xi \leq \bar{\xi}_+$ .
- 2.

$$\rho_0 = \begin{cases} \bar{\rho}_1 > 0, & \text{constant on } \Omega_1, \\ \bar{\rho}_2 > 0, & \text{constant on } \Omega_2, \end{cases} \quad \text{with } \bar{\Omega}_1 \cup \bar{\Omega}_2 = \bar{\Omega}, \quad \text{meas } (\Omega_i) > 0$$

► Properties:

1. non-negativity, boundedness of  $\rho$

$$\bar{\rho}_1 \leq \rho \leq \bar{\rho}_2 \quad \text{in } \Omega_T$$

2. energy law:

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\Omega} \left[ \frac{\rho |\mathbf{u}|^2}{2} + \frac{|\mathbf{b}|^2}{\bar{\mu}} \right] dx + \int_{\Omega} \left[ \eta(\rho) |\boldsymbol{\epsilon}(\mathbf{u})|^2 + \frac{1}{\bar{\mu}^2 \xi(\rho)} |\mathbf{curl} \mathbf{b}|^2 \right] dx \\ = \int_{\Omega} \rho \mathbf{g} \cdot \mathbf{u} dx \end{aligned}$$

# Analytical results

- ▶ weak formulation: For
  - ▶  $\Omega \subset \mathbb{R}^3$  polyhedral domain
  - ▶  $\mathbf{u}_0, \mathbf{b}_0 \in \mathbf{H} := \left\{ \boldsymbol{\xi} \in L^2(\Omega) : \operatorname{div} \boldsymbol{\xi} = \text{weakly in } \Omega, \boldsymbol{\xi} \cdot \mathbf{n} = 0 \text{ on } \partial\Omega \right\}$
  - ▶  $\rho_0$  as above
  - ▶  $\mathbf{f} \in L^2(0, T; L^2)$   
exists a weak solution  
 $\mathbf{u} \in L^\infty(0, T; \mathbf{H}) \cap L^2(0, T; \mathbf{J})$   
 $\mathbf{b} \in L^\infty(0, T; \mathbf{H}) \cap L^2(0, T; \mathbf{X})$   
 $\rho \in L^\infty(\Omega_T) \cap C([0, T]; L^p)$  for all  $p \geq 1$   
which satisfies an energy inequality.

## ... Tools employed

- ▶ general Galerkin method
- ▶ Aubin-Lions compactness result
- ▶ Compactness result by J. Di Perna & P.L. Lions [’89]: Solvability of

$$\rho_t + \operatorname{div}(\mathbf{u}\rho) = f \quad \text{in } \Omega_T, \quad \rho(0, \cdot) = \rho_0 \in L^\infty(\Omega) \quad (2)$$

Let  $\{\rho_k\}_{k \geq 0} \subset L^\infty(0, T; L^\infty(\Omega))$  solve

$$\begin{aligned} (\rho_k)_t + \operatorname{div}(\mathbf{u}_k \rho_k) + [\operatorname{div} \mathbf{u}_k] \rho_k &= f_k \quad \text{in } \Omega_T \\ (\rho_k)(0, \cdot) &= (\rho_k) \quad \text{in } \Omega \end{aligned}$$

Assume:

- (i)  $\{\mathbf{u}_k\}_k \subset L^1(0, T; \mathbf{W}_0^{1,2})$ , and  $\mathbf{u} \in L^1(0, T; \mathbf{W}_0^{1,2})$ , such that  
 $\mathbf{u}_k \rightarrow \mathbf{u}$  in  $L^1(0, T; \mathbf{L}^2)$   
 $\operatorname{div} \mathbf{u}_k \rightarrow \operatorname{div} \mathbf{u}$  in  $L^1(0, T; L^2)$
- (ii)  $f_k \rightarrow f$  in  $L^1(0, T; L^2)$   
 $(\rho_k)_0 \rightarrow \rho_0$  in  $L^2(\Omega)$

Then

$$\rho_k \rightarrow \rho \text{ in } L^2(0, T; L^2)$$

where  $\rho : \Omega \rightarrow \mathbb{R}$  is unique solution of (2).

# Problems to construct a convergent Finite Element Discretization

- **Discrete Energy Law:** In continuous setting: multiply  
 $(1_1)$  with  $\mathbf{u}$   
 $(1_3)$  with  $\frac{1}{2}|\mathbf{u}|^2$

**Observation:**  $\frac{1}{2}|\mathbf{u}|^2$  no admissible test function in  
FE-discretization

**Idea (N. Walkington ['07]):** reformulation

$$(\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) = \frac{1}{2} \left\{ \rho \mathbf{u}_t + [\rho \mathbf{u} \cdot \nabla] \mathbf{u} + (\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) \right\}$$

- ▶ (uniform) positivity and  $L^\infty$ - boundedness of discrete densities:

**Idea:** M-matrix property of stiffness matrix related to

$$(d_t \rho^n, \chi)_h + (\mathbf{U}^n \cdot \nabla \rho^n, \chi) + \frac{1}{2} ([\operatorname{div} \mathbf{U}^n] \rho^n, \chi) + \alpha h^\alpha (\nabla \rho^n, \nabla \chi) = 0$$

### Tools:

- ▶  $\alpha > 0$ : M-matrix property gives  $0 < \rho_1 \leq \rho^n \leq \rho_2 < \infty$
- ▶ numerical integration
- ▶ regularization term to NSE  $\beta_2 h^{\beta_2} (\nabla d_t \mathbf{U}^n, \nabla \mathbf{W}^n)$        $\beta_2 \geq 0$
- ▶ discrete version of compactness result by R. DiPerna & P.L. Lions: Idea (N. Walkington [’07])

Tool: to validate  $\text{div} \mathcal{U} \rightarrow \text{div} \mathbf{u}$  in  $L^1(0, T; L^2)$   
 we add the regularization term to NSE

$$\beta_1 k^{\beta_1} (\text{div} \mathbf{U}^n, \text{div} \mathbf{W}^n) , \quad \beta_1 > 0$$

- ▶ a compactness result of J.L. Lions to control temporal changes of iterates: There exists  $\kappa > 0$

$$\int_0^T \left[ \rho_1 \|\mathcal{U}^+ - \mathcal{U}^-\|^2 + \|\mathcal{B}^+ - \mathcal{B}^-\|^2 \right] ds \leq Ck^\kappa$$

Tool:

- ▶ inverse estimates
- ▶ mesh-constraints:  $F(k, h) \geq 0$
- ▶ Problems already experienced for one-fluid MHD equation:  
 $\mathbf{H}(\mathbf{curl}), \mathbf{H}(\text{div})$ ,

Tool: discrete compactness result by F. Kikuchi [’89]

## Result: A stable, convergent FE–based discretization

- ▶ Let  $\{(\rho^n, \mathbf{U}^n, \mathbf{B}^n)\}$  be solution of FE –based fully discrete scheme
- ▶ L. Banas, A.P. [’08]: Let
  - ▶  $\mathcal{T}$  be strongly acute
  - ▶  $F(k, h; \alpha, \beta_1, \beta_2, d) \geq 0$  be valid
  - ▶  $\rho^0 \rightarrow \rho_0$  in  $L^2$ ,  $(\mathbf{U}^0, \mathbf{B}^0) \rightharpoonup (\mathbf{u}_0, \mathbf{b}_0)$  in  $[L^2]^2$

For  $(k, h) \rightarrow 0$  exist a convergent subsequence, and  $(\rho, \mathbf{u}, \mathbf{b})$ , s.t.

$\mathbf{U} \rightharpoonup \mathbf{u}$  in  $L^\infty(0, T; \mathbf{L}^2)$ ,  $\mathbf{B} \rightharpoonup \mathbf{b}$  in  $L^\infty(0, T; \mathbf{L}^2)$ ,  $\sigma \xrightarrow{*} \rho$  in  $L^\infty(0, T; \mathbf{L}^\infty)$

Where  $(\mathbf{u}, \mathbf{b}, \rho)$  is weak solution to (1).

# From Discretization to a fully practical Scheme

- ▶ Status:
  - + a convergent discretization of (1)
  - a nonlinear algebraic problem
- ▶ Tool: A fixed point strategy for every  $n \geq 1$ 
  - + linear, decoupled problems for  $\{(\rho^{n,l}, \mathbf{u}^{n,l}, \mathbf{b}^{n,l})\}_{n,l}$
  - + a **thresholding** criterion
  - + contraction property for restrictive  $F(k, h) \geq 0$
  - + perturbed discrete energy law, and
  - + uniform upper and lower  $L^\infty$ -bounds for  $\{\rho^{n,l}\}_{n,l}$

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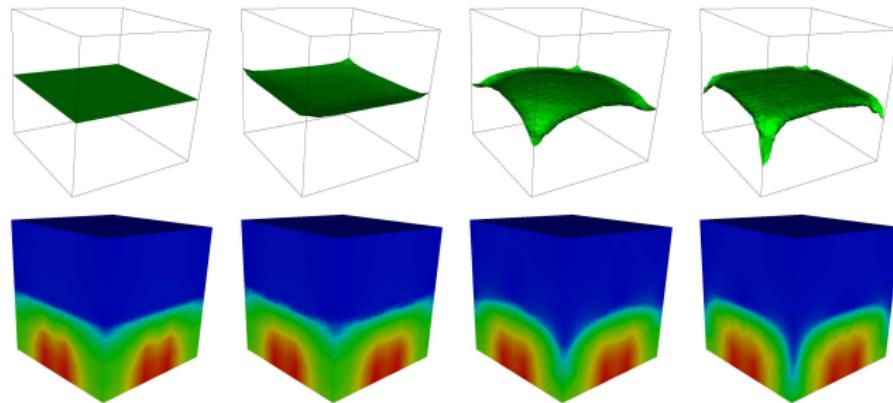
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## PART IV: Simulation of Aluminium Electrolysis

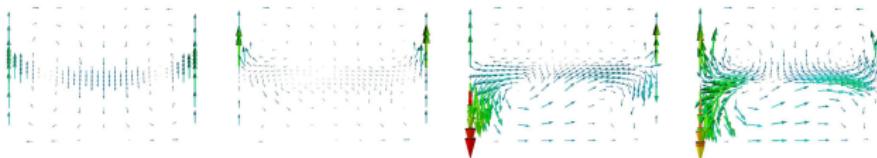
# Aluminium Electrolysis – Two-Fluid MHD

- ▶ evolution of interface between two conducting fluids
  - ▶ top: lighter cryolite, bottom: heavier liquid aluminium

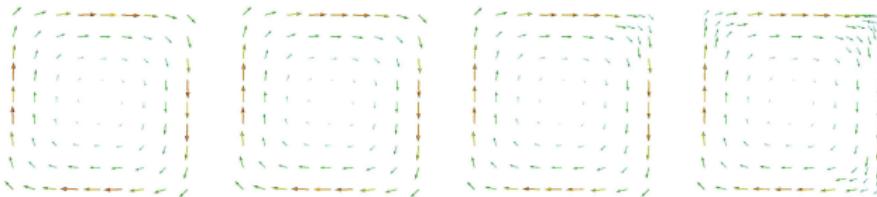


- ▶ magnitude of magnetic field

- ▶ 2D Velocity profile



- ▶ 2D Magnetic field profile



## PART V: Summary

- ▶ Aluminium production: Two Fluid–MHD problem
- ▶ **Goal/Problem:** Develop convergent FE–based discretization such that computed iterates satisfy relevant properties
- ▶ **Problems:** Discrete energy law, upper and lower bounds for density, compactness result by R. DiPerna & P.L. Lions
- ▶ **Tools:** reformulation, M–matrix property, numerical integration, regularization terms, mesh constraints

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