



# ***Adaptive Resolution Molecular Dynamics Simulation***

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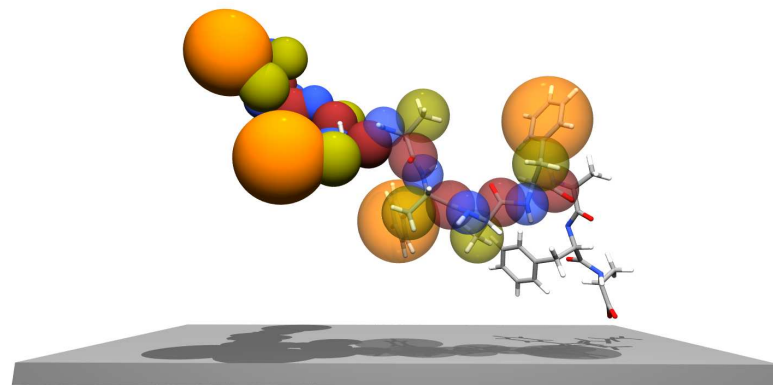
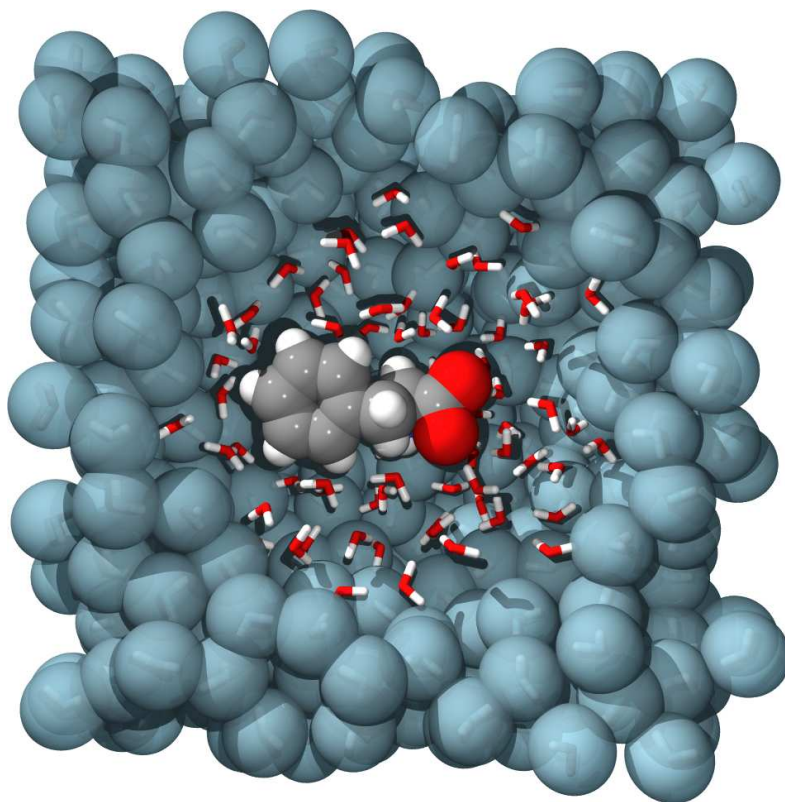
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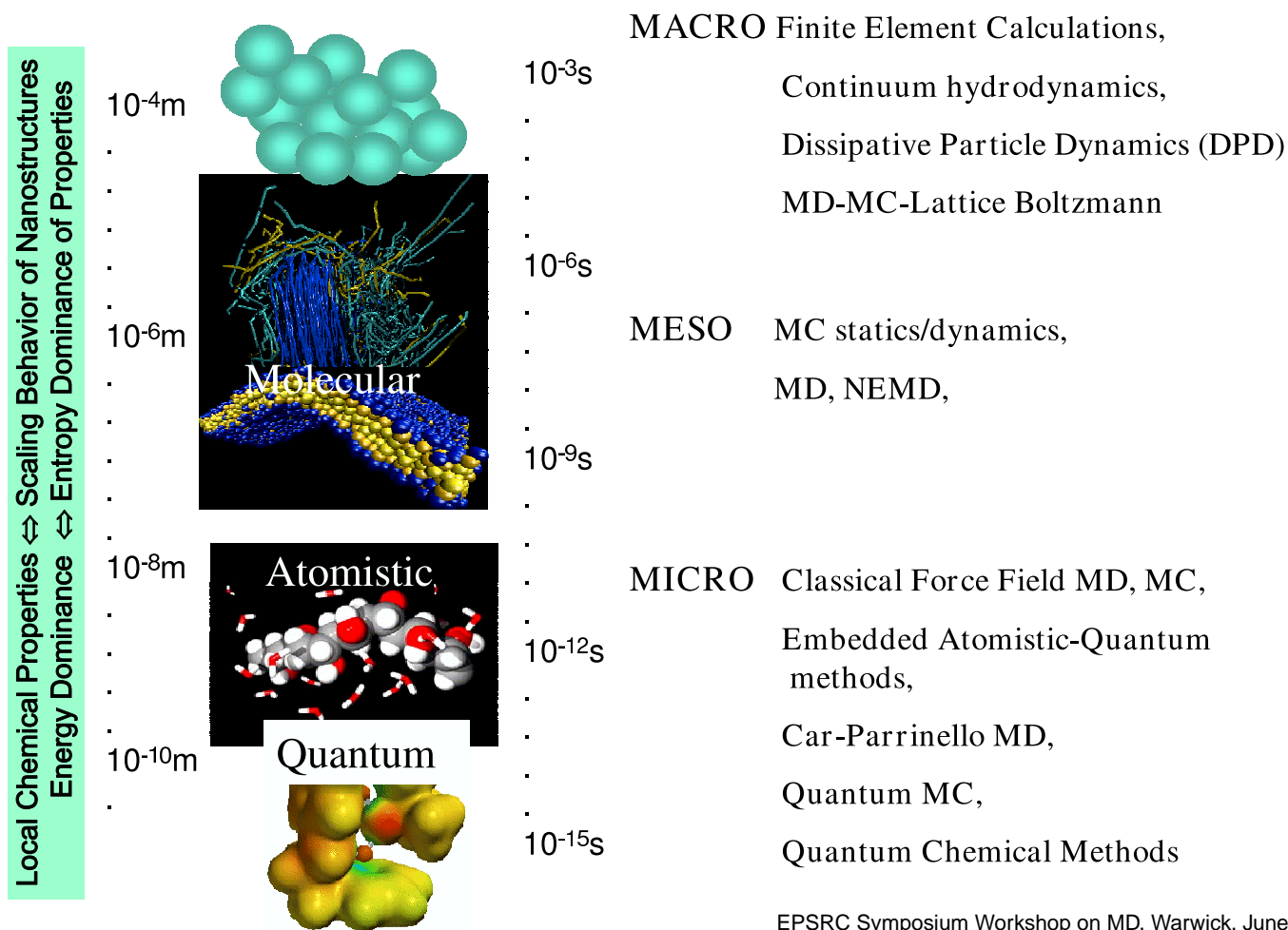
**Silvina Matysiak**, Rice University, Houston, Texas

**Brad Lambeth**, Rice University, Houston, Texas

# Multiscale Phenomena



# Multiscale Modeling



# Adaptive resolution simulation



## ⑥ Motivation:

- △ to simplify the model to the largest extent possible while keeping all the necessary details where this is required.

## ⑥ Method: *AdResS (Adaptive Resolution Scheme)*

- △ allows for an dynamical switching between the atomistic and mesoscopic levels of detail  $\implies$  on-the-fly changing of the number of DOFs
- △ tailor-made for molecular systems where spatially localized domains with the required atomistic resolution exchange particles with the remainder of the system sufficiently described on the mesoscopic scale.

## ⑥ Results:

- △ accurately reproduces the statistical properties of the reference all-atom system.



## **All-Atom MD simulation:**

- ⑥ allows to study processes at the atomic level of detail
- ⑥ is often incapable to bridge a gap between a wide range of length and time scales involved in molecular systems

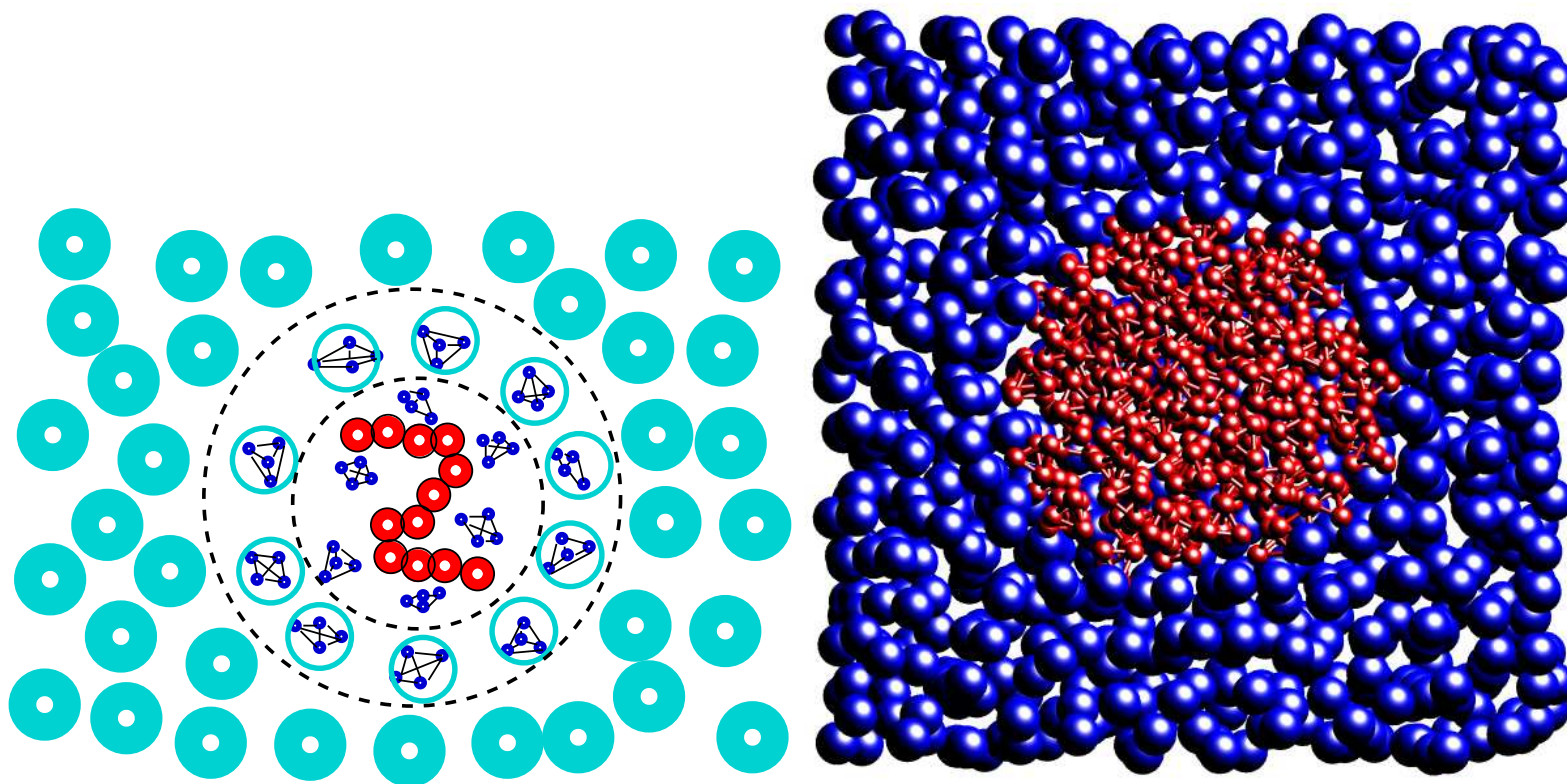
## **Mesoscopic MD simulation:**

- ⑥ reduces the number of DOFs by retaining only those that are relevant for the property of interest  $\implies$  longer length and time scales can be reached
- ⑥ specific chemical details are usually lost in the coarse-graining procedure

## **Combining the best from both approaches:**

- ⑥ Hybrid Adaptive MD Schemes

# *Hybrid atomistic/mesososcopic liquid*



MP, L. Delle Site, K. Kremer, J. Chem. Phys. **126**, 134902 (2007).

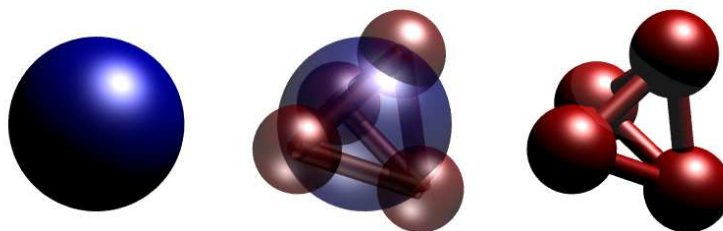
MP, L. Delle Site, K. Kremer, Phys. Rev. E **73**, 066701 (2006).





# Changing Number of Degrees of Freedom

- ⑥ A tetrahedral molecule has a defined spatial orientation and  $3N = 12$  DOFs:
  - △ 3 translational
  - △ 3 rotational
  - △  $3N - 6 = 6$  vibrational
- ⑥ One particle mesoscopic molecule has no defined spatial orientation and only 3 translational DOFs.



MP, L. Delle Site, K. Kremer, J. Chem. Phys. **123**, 224106 (2005).

MP, L. Delle Site, K. Kremer, Annu. Rev. Phys. Chem. **59**, 545 (2008).





# ***Geometrically induced phase transition***

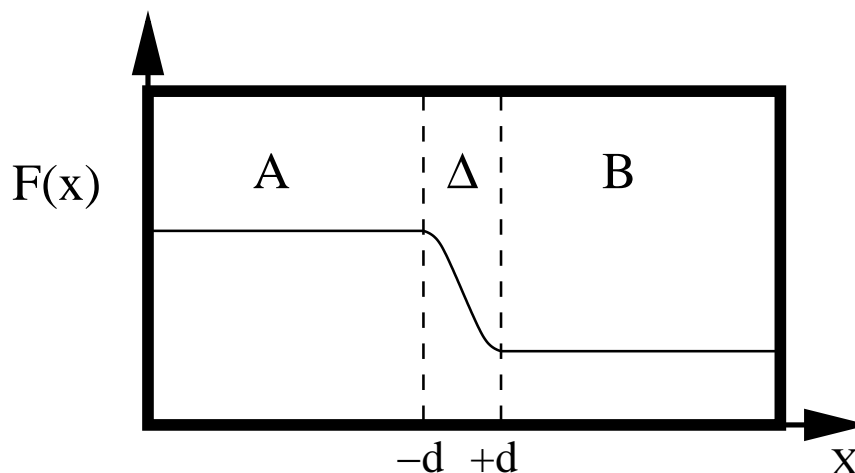


In thermodynamical equilibrium, boundary conditions analogous to two-phase coexistence must be satisfied:

$$\mu_{ex} = \mu_{cg}, \quad p_{ex} = p_{cg}, \quad T_{ex} = T_{cg}.$$

The rotational and vibrational parts of the free energy can be viewed as the latent heat, which is supplied or taken by the thermostat, at this transition.

# Transition region



- ⑥ Molecules in A and B are **physically identical but differently represented**.
- ⑥ The number of DOFs is  $n = n(x)$  with:  $n_A = \text{const}_A$ ;  $n_B = \text{const}_B$ ; and  $n_\Delta = n(x)$

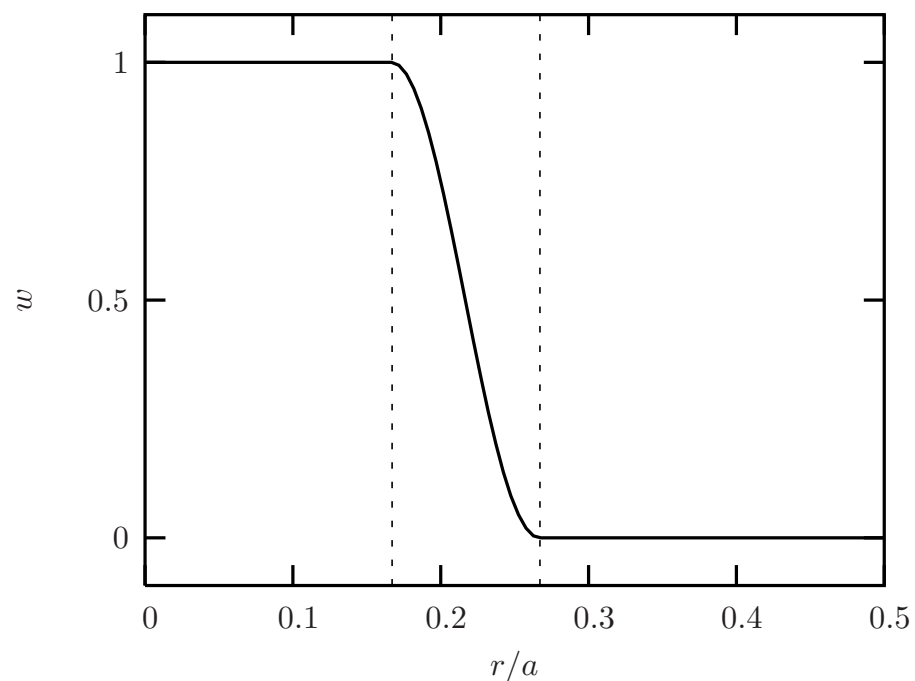
- ⑥ The system is in equilibrium which implies:

$$\lim_{x \rightarrow d^-} \frac{\partial F_A(x)}{\partial x} = \lim_{x \rightarrow d^+} \frac{\partial F_B(x)}{\partial x} = 0 \implies$$

$$\lim_{x \rightarrow d^-} \frac{\partial n_A(x)}{\partial x} = \lim_{x \rightarrow d^+} \frac{\partial n_B(x)}{\partial x} = 0$$



# Weighting Function = Order Parameter



- ⑥ The values  $w = 1$  and  $w = 0$  correspond to the atomistic and coarse-grained regions, respectively, while the values  $0 < w < 1$  correspond to the transition (*hyb*) regime.



# AdResS: Linking Atomic and Mesoscopic Length Scales



AdResS consists of two main steps:

1. Derive the effective pair potential  $U^{cm}$  between coarse-grained molecules on the basis of the reference all-atom system.
2. Couple the atomistic and mesoscopic scales:

$$\mathbf{F}_{\alpha\beta} = w(X_{\alpha})w(X_{\beta})\mathbf{F}_{\alpha\beta}^{atom} + [1 - w(X_{\alpha})w(X_{\beta})]\mathbf{F}_{\alpha\beta}^{cm},$$

where

$$\mathbf{F}_{\alpha\beta}^{atom} = \sum_{i\alpha, j\beta} \mathbf{F}_{i\alpha j\beta}^{atom}$$

is the sum of all pair interactions between explicit atoms of molecules  $\alpha$  and  $\beta$  and

$$\begin{aligned}\mathbf{F}_{i\alpha j\beta}^{atom} &= -\frac{\partial U^{atom}}{\partial \mathbf{r}_{i\alpha j\beta}}, \\ \mathbf{F}_{\alpha\beta}^{cm} &= -\frac{\partial U^{cm}}{\partial \mathbf{R}_{\alpha\beta}}.\end{aligned}$$

*May the Force be with you*



One must interpolate the **forces** and not the interaction potentials  
if the **Newton's Third Law** is to be satisfied!

MP, K. Kremer, L. Delle Site, J. Phys. A: Math. Theor. **40**, F281, 2007.



# The extension of equipartition theorem to non-integer DOFs



- ⑥ For the fractional quadratic DOF  $\Theta$  with the weight  $w = \alpha$  we can write the partition function as:

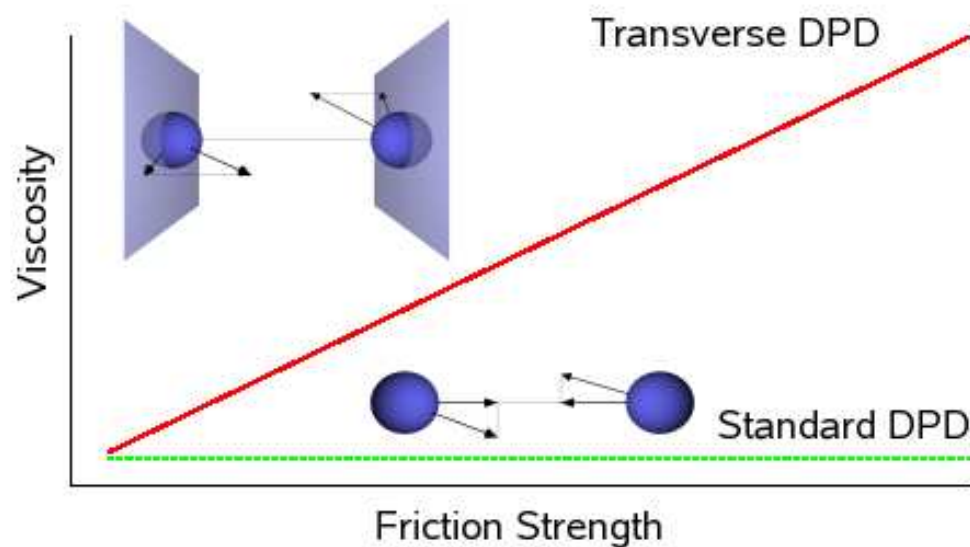
$$\begin{aligned}\exp(-\beta F_\alpha) &= C \int \exp(-\beta f(\alpha) p_\Theta^2/2) dV_\alpha = \\ &= 2C \int_0^\infty \exp(-\beta f(\alpha) p_\Theta^2/2) |p_\Theta|^{\alpha-1} \frac{dp_\Theta}{\Gamma(\alpha)} = \\ &= \frac{2^{\alpha/2} C \Gamma(\alpha/2)}{\Gamma(\alpha)} f(\alpha)^{-\alpha/2} \beta^{-\alpha/2} \sim \beta^{-\alpha/2}.\end{aligned}$$

- ⑥  $\langle K_\alpha \rangle = \frac{d(\beta F_\alpha)}{d\beta} = \frac{\alpha}{2\beta} = \frac{\alpha k_B T}{2}.$

- ⑥ In equilibrium  $T_A = T_B = T_\Delta = T$  and thus:  $n_\alpha \sim \alpha.$

MP, K. Kremer, L. Delle Site, Phys. Rev. E **75**, 017701 (2007).

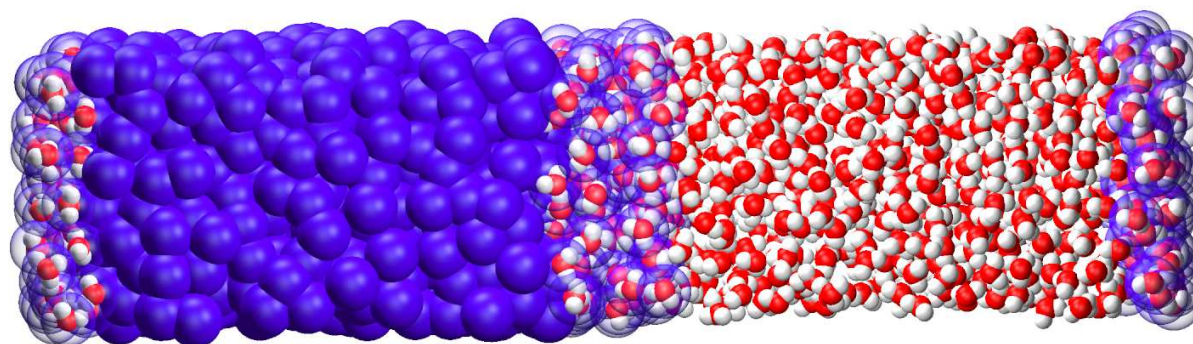
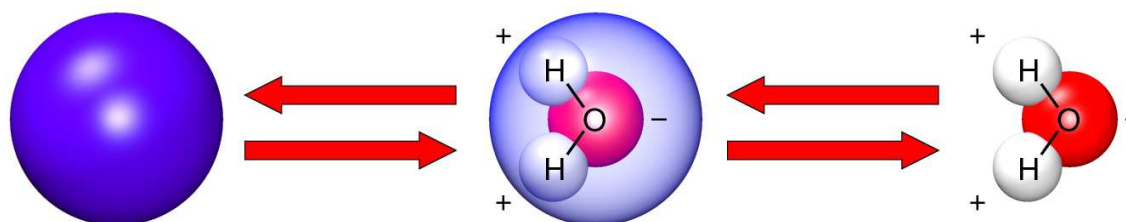
# Transverse DPD Thermostat



- ⑥ The variation of the **dissipative particle dynamics (DPD)** thermostat includes the damping of the **perpendicular** components of the relative velocity, yet keeping the advantages of conserving **Galilei invariance** and within our error bar also **hydrodynamics**. It allows for **controlling transport properties** of molecular fluids.

C. Junghans, MP, K. Kremer, *Soft Matter* **4**, 156 (2008).





- ⑥ MP, S. Matysiak, L. Delle Site, K. Kremer, C. Clementi, J. Phys.: Condens. Matter **19**, 292201 (2007).
- ⑥ S. Matysiak, C. Clementi, MP, K. Kremer, L. Delle Site, J. Chem. Phys. **128**, 024503 (2008).

# Electrostatics: Reaction field method



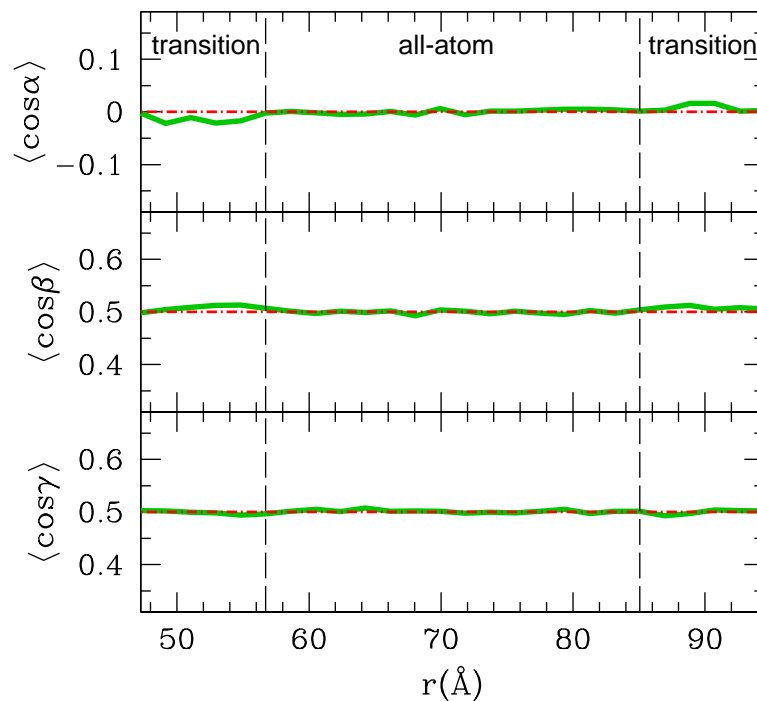
The electrostatic forces interactions are described using the Reaction field (RF) method:

$$\mathbf{F}_{C_{i\alpha j\beta}}^{atom}(\mathbf{r}_{i\alpha j\beta}) = \frac{e_{i\alpha} e_{j\beta}}{4\pi\epsilon_0} \left[ \frac{1}{r_{i\alpha j\beta}^3} - \frac{1}{R_c^3} \frac{2(\epsilon_{RF} - 1)}{1 + 2\epsilon_{RF}} \right] \mathbf{r}_{i\alpha j\beta}.$$

The RF is suitable to be used with AdResS because:

- ⑥ it is pairwise
- ⑥ like AdResS it must also be applied with a thermostat

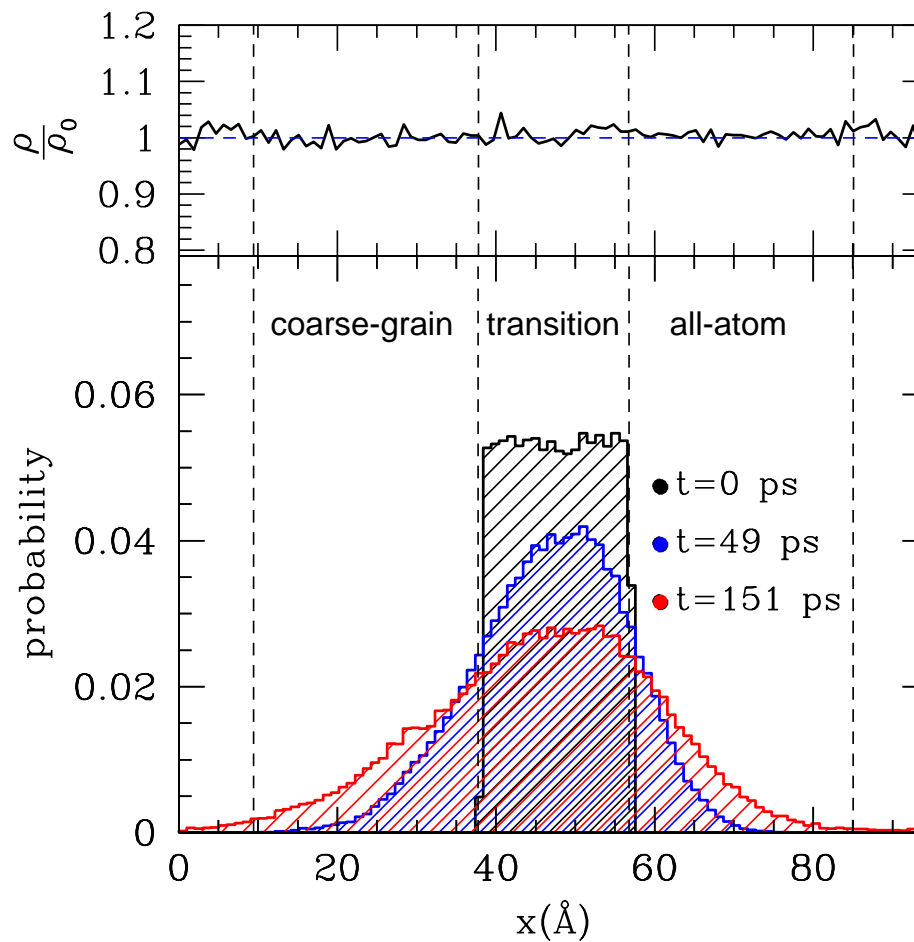
# Interface effect of the cg water



The transition regime neutralizes the interface effect of the cg water  $\implies$  the structure of water in the explicit regime is the same as in the bulk.



# Diffusion across the transition regime



# AdResS: Conclusions



## ⑥ Adaptive Resolution MD simulation:

- △ Changing resolution is formally equivalent to a phase transition → latent heat.
- △ For a smooth variation of the resolution we introduce a transition regime.
- △ The temperature in the transition region can be obtained by extending the equipartition theorem to non-integer dimensions.

## ⑥ Hybrid method AdResS:

- △ Allows for a dynamical switching of the spatial resolution.
- △ We treat only as many DOFs as absolutely necessary for the problem considered.
- △ AdResS was so far applied to MD simulations of a simple tetrahedral liquid, a macromolecule in an explicit solvent, and liquid water at standard conditions.

## ⑥ Future work:

- △ Application to different soft matter systems and molecular liquids.

# Concurrent Triple-Scale Simulation



## ⑥ Motivation:

- △ to perform MD of **open domains** with relatively large molecules, either in the **grand canonical ensemble** or under **non-equilibrium conditions**.

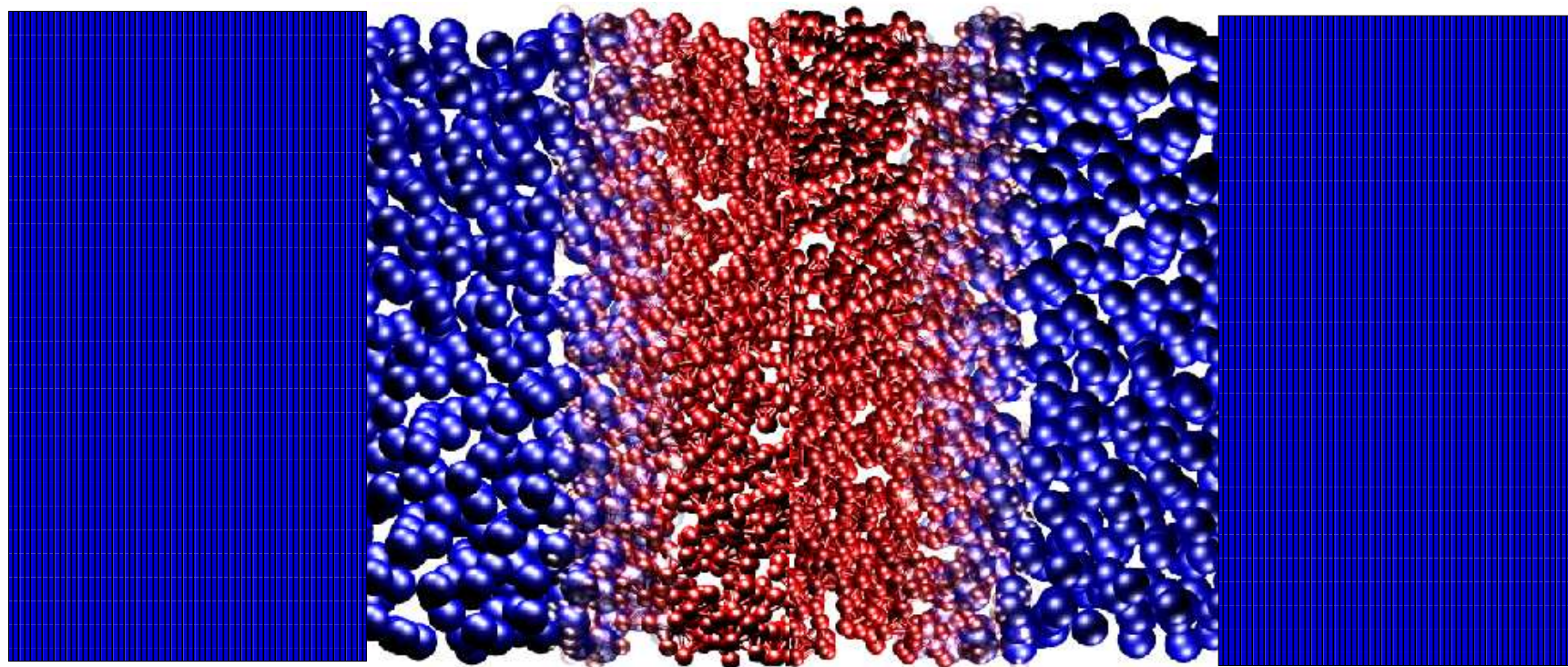
## ⑥ Method: *Triple-scale AdResS-HybridMD scheme*

- △ is a combination of two dual-scale models: a particle-based **Adaptive Resolution Scheme (AdResS)**, which couples the **atomic** and **mesoscopic** scales, and a **hybrid continuum-molecular dynamics scheme (HybridMD)**
- △ covers the **length-scales** ranging from the **micro-** to **macro-scale**
- △ successfully sorts out the problem of **large molecule insertion** in the hybrid particle-continuum simulations of molecular liquids
- △ opens up the possibility to perform efficient **grand-canonical molecular dynamics simulations** of truly open molecular liquid systems

## ⑥ Results:

- △ the **structural** and **dynamical** properties of the liquid are **accurately captured**

# *Triple-Scale Model*



R. Delgado Buscalioni, K. Kremer, MP, J. Chem. Phys. **128**, 114110 (2008).





# HybridMD: Coupling Particle-Based and Continuum Descriptions

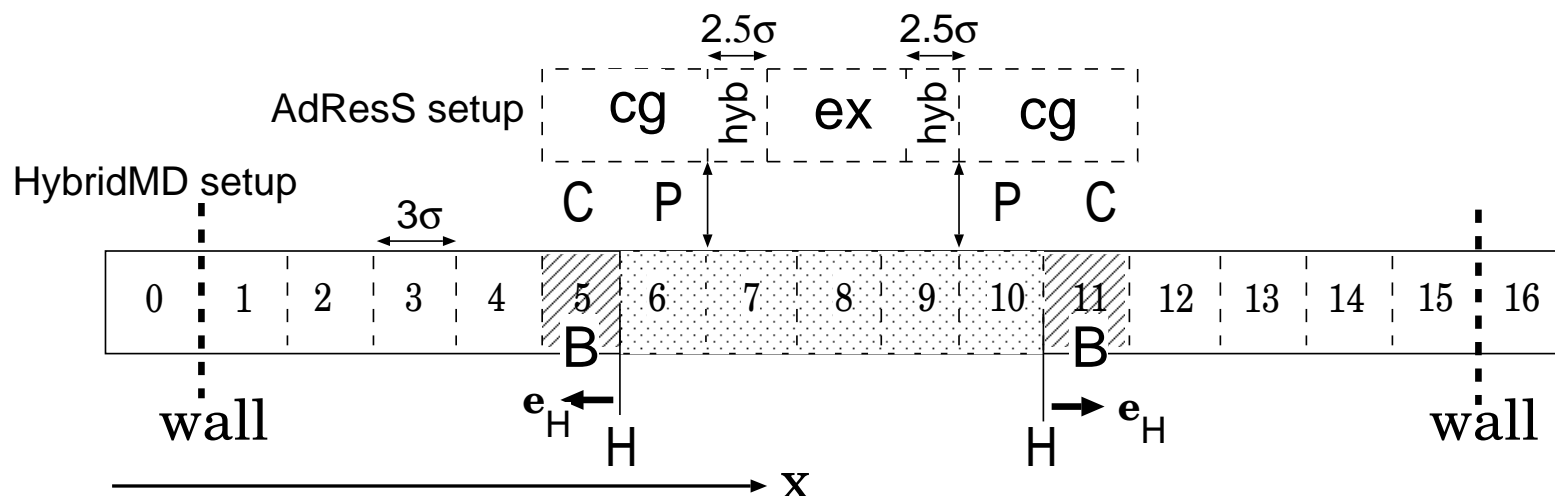


- ⑥ The hybrid particle-continuum scheme (HybridMD) is designed to connect the dynamics of a “**molecular domain**” with that obtained from a **continuum description** of the surrounding fluid flow.
- ⑥ The method is based on **flux-exchange**.
- ⑥ The system is divided in (at least) two domains, described via classical **molecular dynamics (MD)** and **continuum fluid dynamics (CFD)**, i.e., solving the **Navier-Stokes** equations.
- ⑥ The MD and CFD domains share one unique “hybrid interface”,  $H$ : Flux balance implies the **conservation of mass and momentum** across  $H$ .

G. De Fabritiis, R. Delgado Buscalioni, P. Coveney, Phys. Rev. Lett **97**, 134501 (2006).  
R. Delgado Buscalioni, G. De Fabritiis, Phys. Rev. E **76**, 036709 (2007).



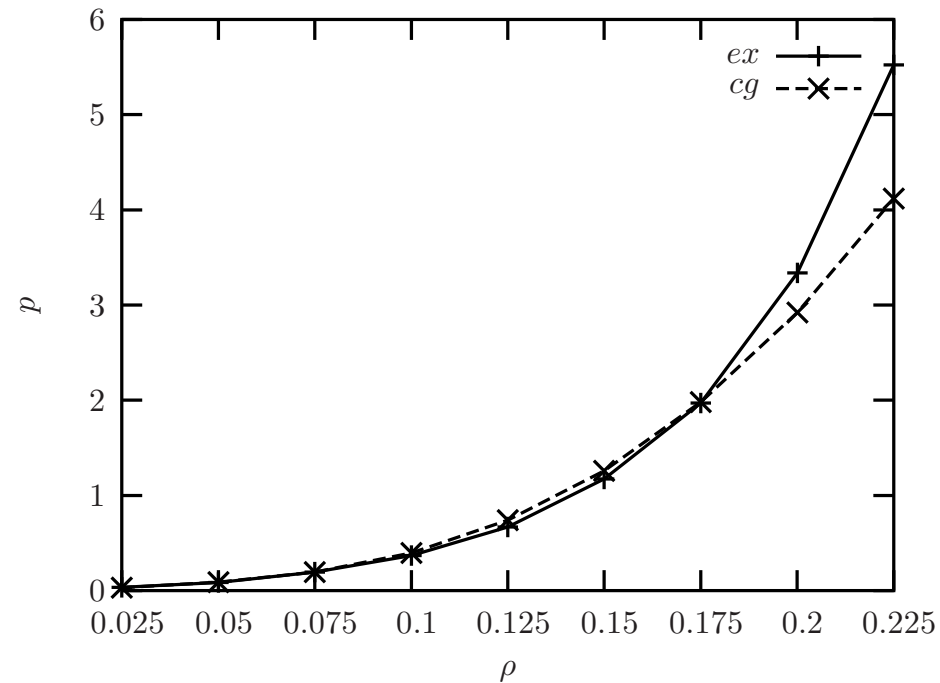
# AdResS-HybridMD: Combined Scheme



- 6 Domain decomposition of the combined scheme. The top part of the figure shows the location of the fluid model layers (*cg*, *hyb* and *ex*) within the HybridMD setup. The bottom part of the figure shows the set of control cells used in the HybridMD setup.

R. Delgado Buscalioni, K. Kremer, MP, J. Chem. Phys. **128**, 114110 (2008).

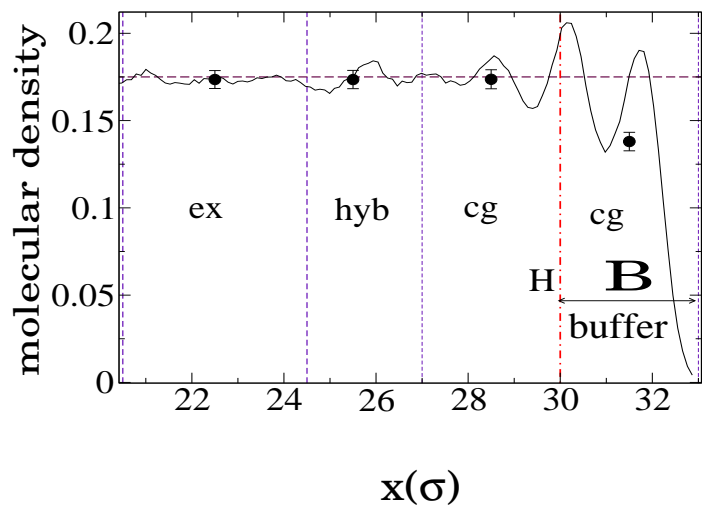
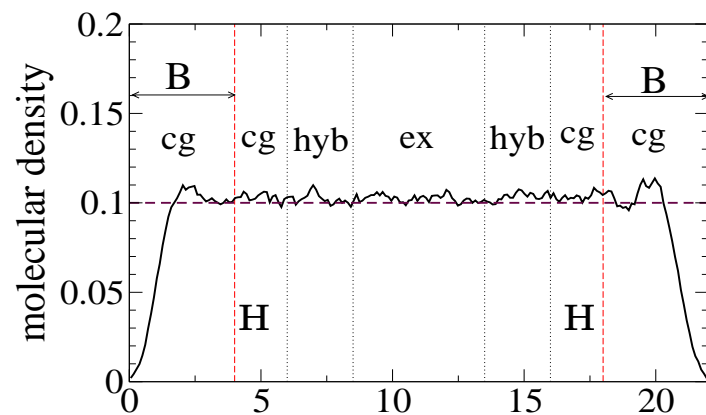
# Equation of state



The pressure tensor:

$$\mathbf{J} = p \mathbf{I} + \rho \mathbf{v} \mathbf{v} + \mathbf{\Pi}$$

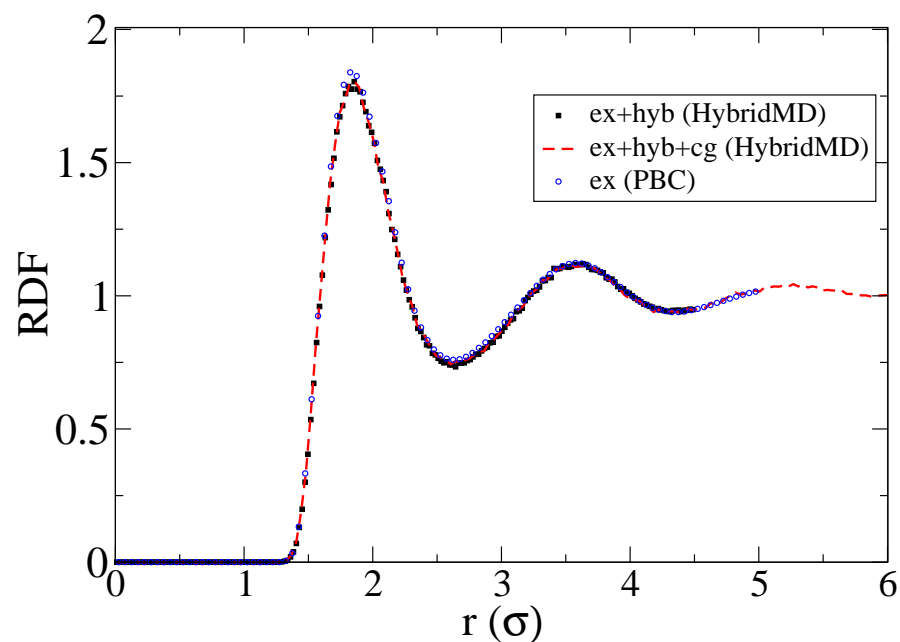
# Molecular Density Profile



(a)  $\rho_m = 0.1\sigma^{-3}$ . (b)  $\rho_m = 0.175\sigma^{-3}$ .

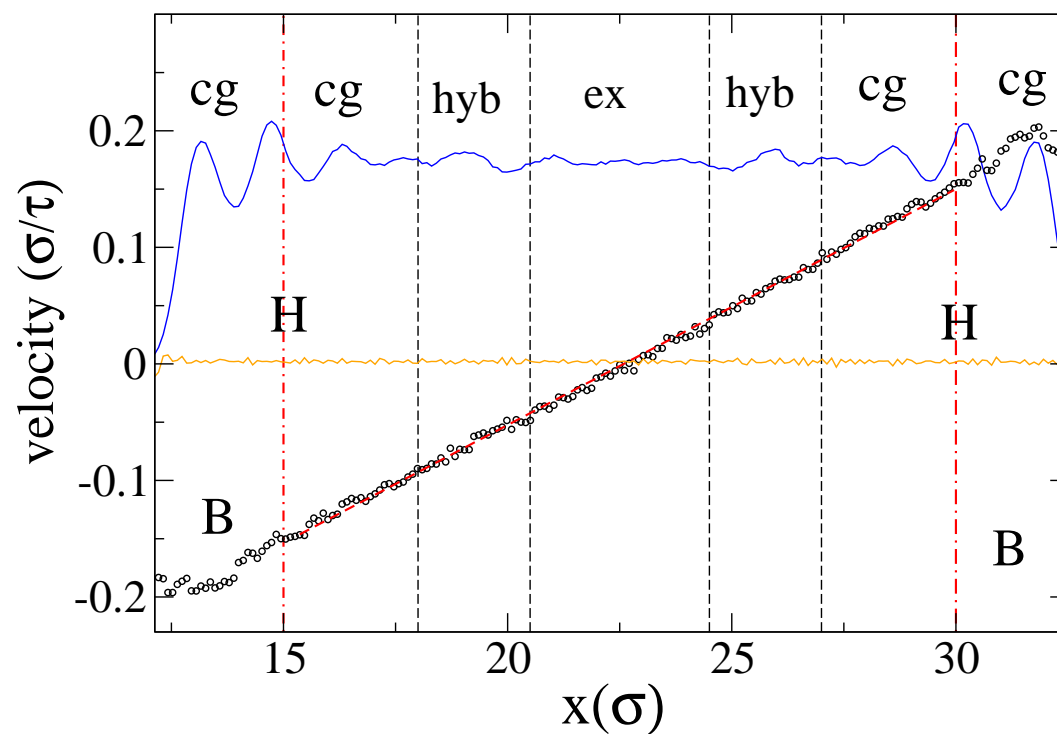


# Radial Distribution Functions: Equilibrium

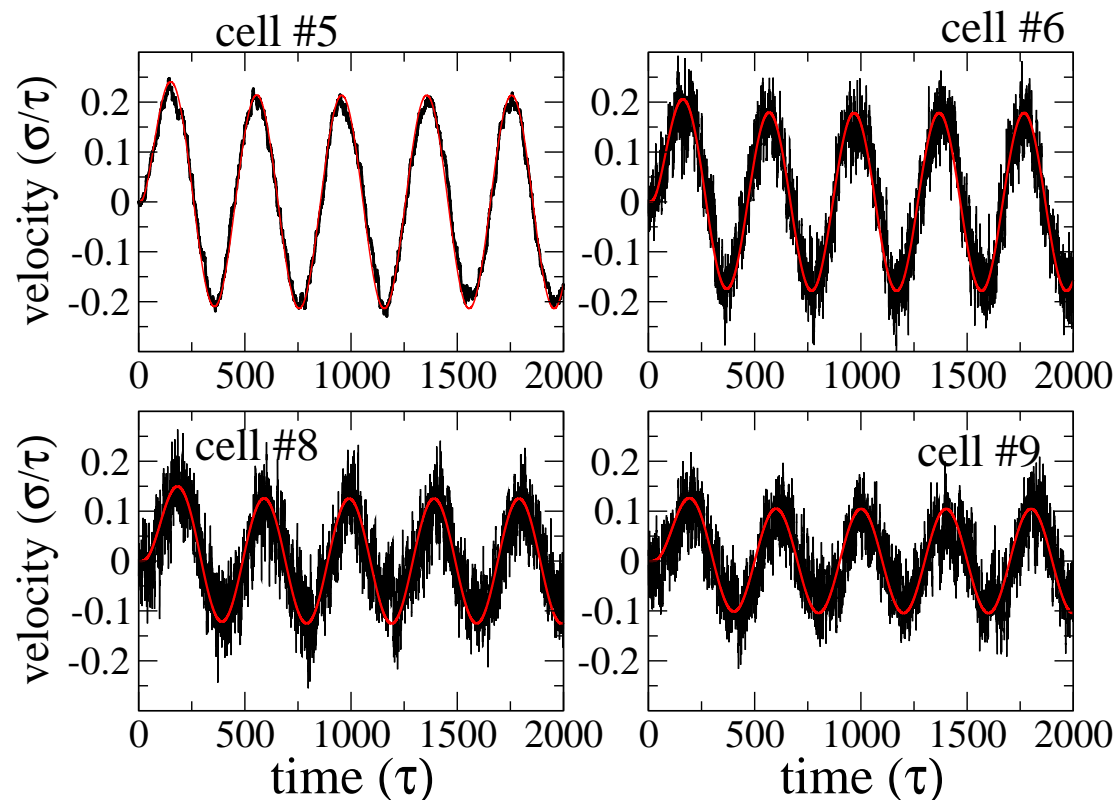


- ⑥ RDF<sub>cm</sub>s of the liquid in the atomistic and transition domains ( $ex + hyb$ ) and in the total molecular region ( $ex + hyb + cg$ ) of the triple-scale model together with the reference RDF<sub>cm</sub> of the all-atom system ( $ex(PBC)$ ) at  $\rho = 0.175/\sigma^3$ .

# Couette Flow



- 6 Velocity profile at the particle region of an hybrid simulation of a Couette flow.



- Velocity in the y-direction at some selected cells in a hybrid simulation of a Stokes flow.

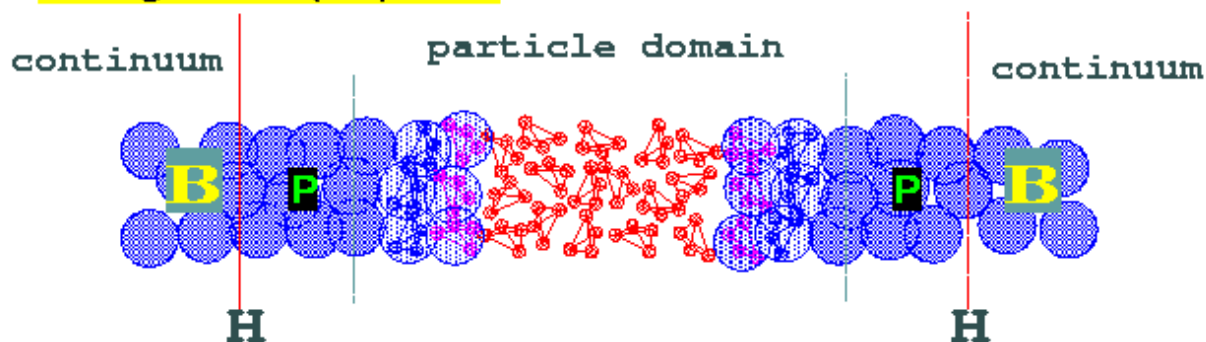




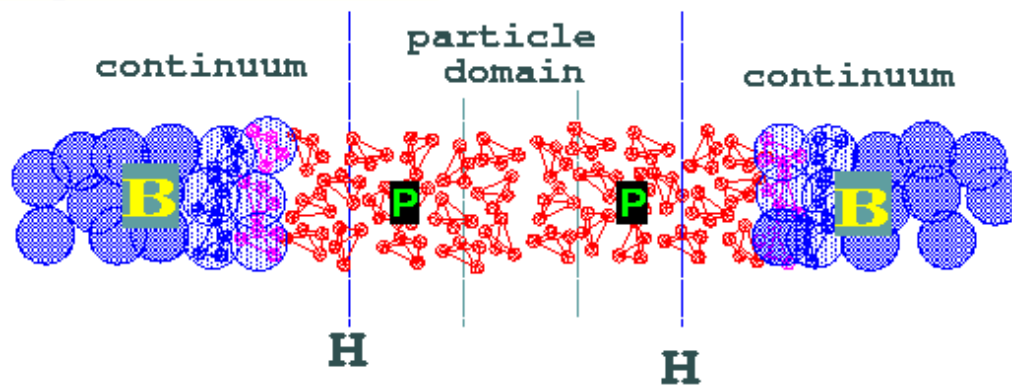
# Triple-Scale Simulation: Two Possible Setups



## Homogeneous (CG) buffer



## Heterogeneous model buffer



R. Delgado Buscalioni, K. Kremer, MP, J. Chem. Phys. **128**, 114110 (2008).

# Buffer: Two Possible Setups



## Homogeneous buffer:

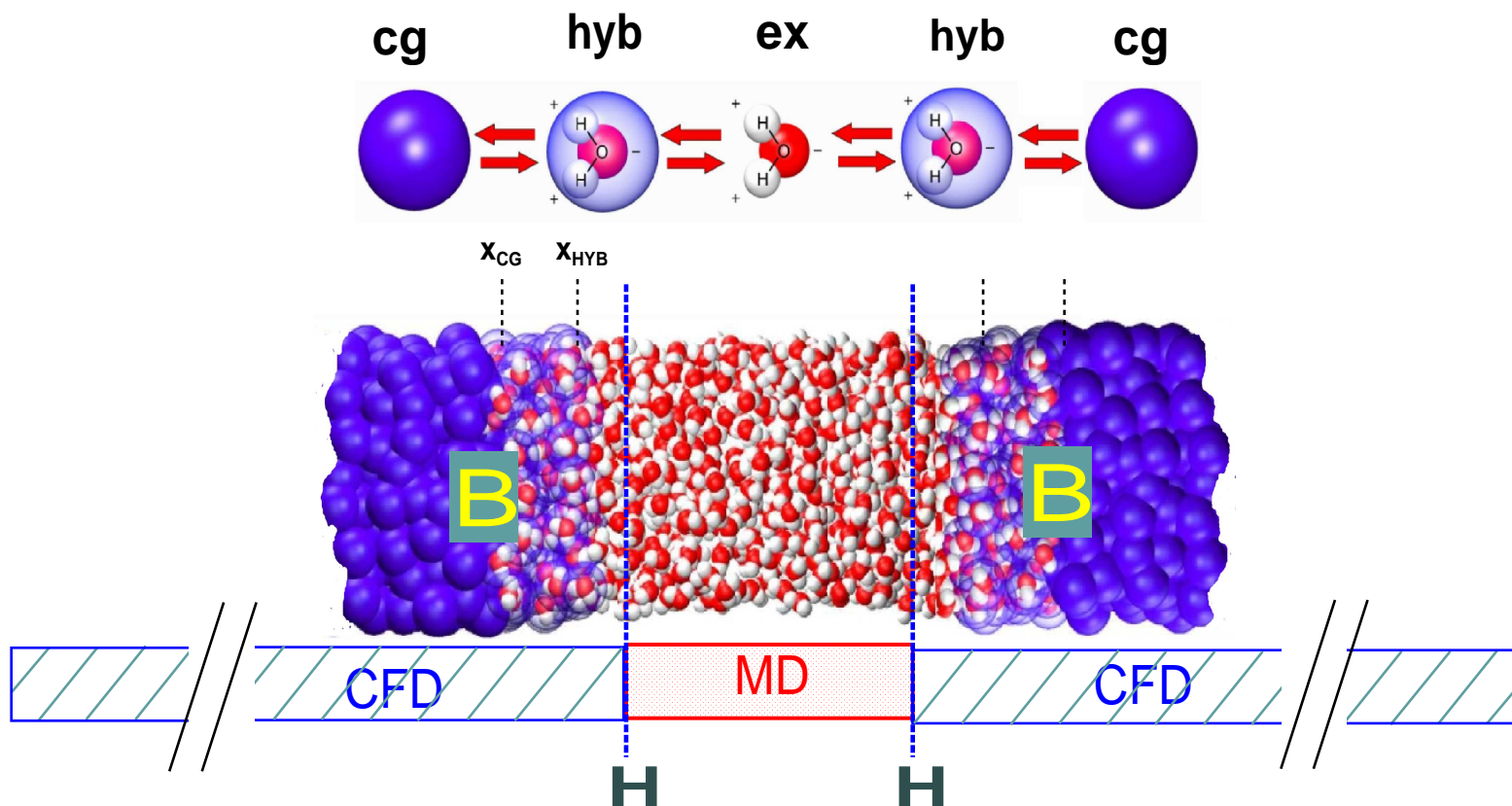
- △ *pro*: Requires small buffer size.
- △ *pro*: Allows us to introduce CG molecular information into the explicit MD region (structure, diffusion rates, etc.)
- △ *con*: Requires fine tuning of CG model ,i.e., EOS, viscosity or diffusion coefficient: *Transversal DPD*; C. Junghans, MP, K. Kremer, *Soft Matter* 4, 156 (2008).



## Heterogeneous buffer:

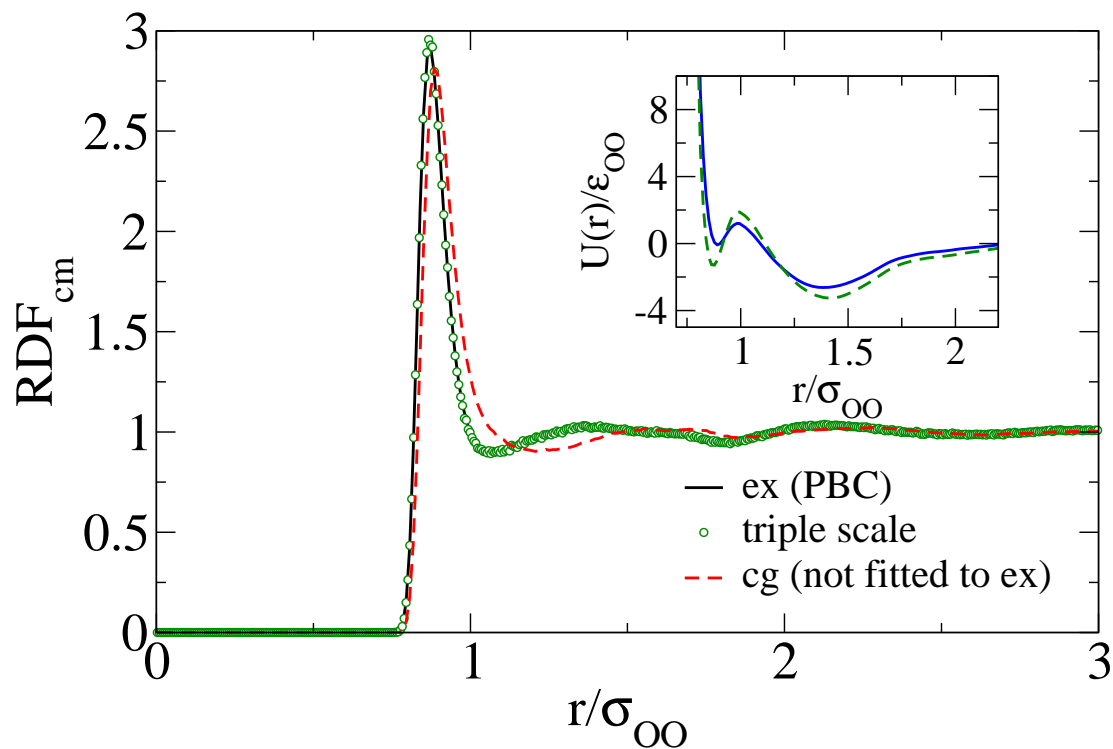
- △ *pro*: All-atom MD region: correct viscosity, EOS, fluctuations.
- △ *pro*: Does not requires fine tuning of CG and HYB models.
- △ *pro*: Enables energy exchange, as the MD region is fully explicit.
- △ *con*: Larger buffer size.

# Triple-Scale Simulation: Liquid Water



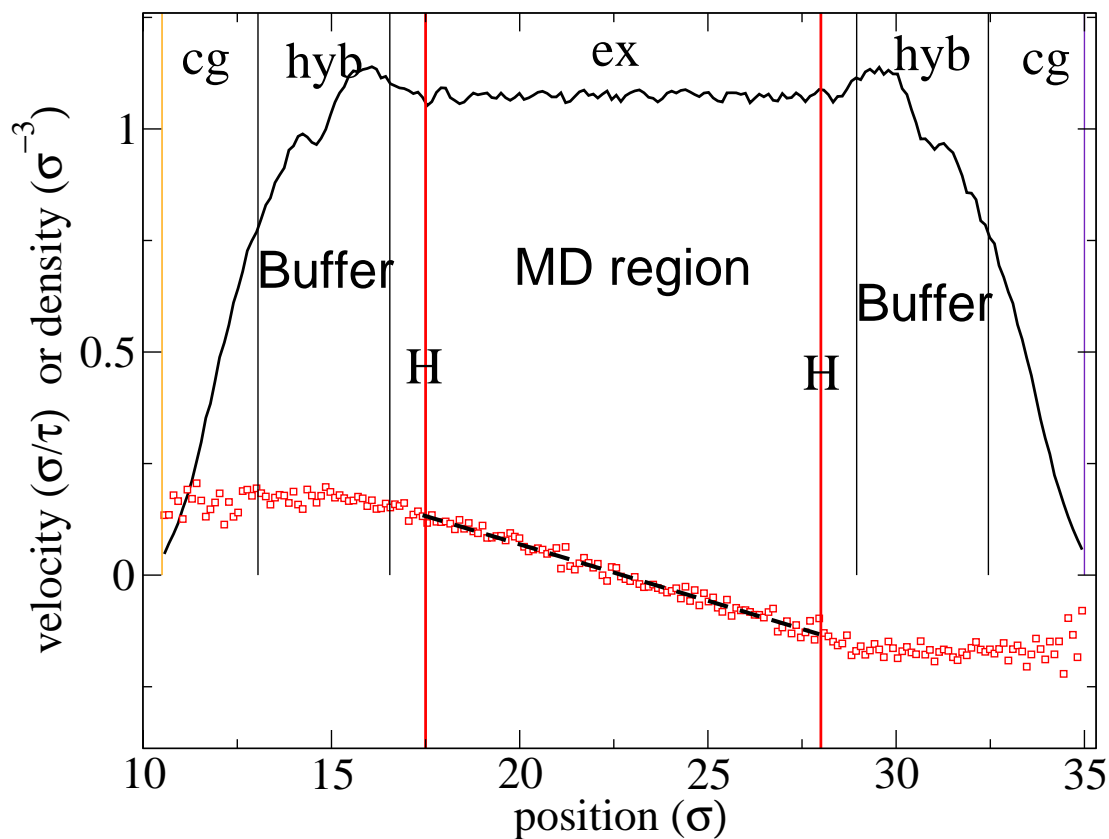
R. Delgado Buscalioni, K. Kremer, MP, submitted.

# Coarse-Grained Model



- ⑥ Center-of-mass RDF of the flexible TIP3P water model and the effective potential.

# Couette Flow: Liquid Water



- ⑥ Density profile and velocity distribution across the particle domain.

# Grand Canonical Ensemble



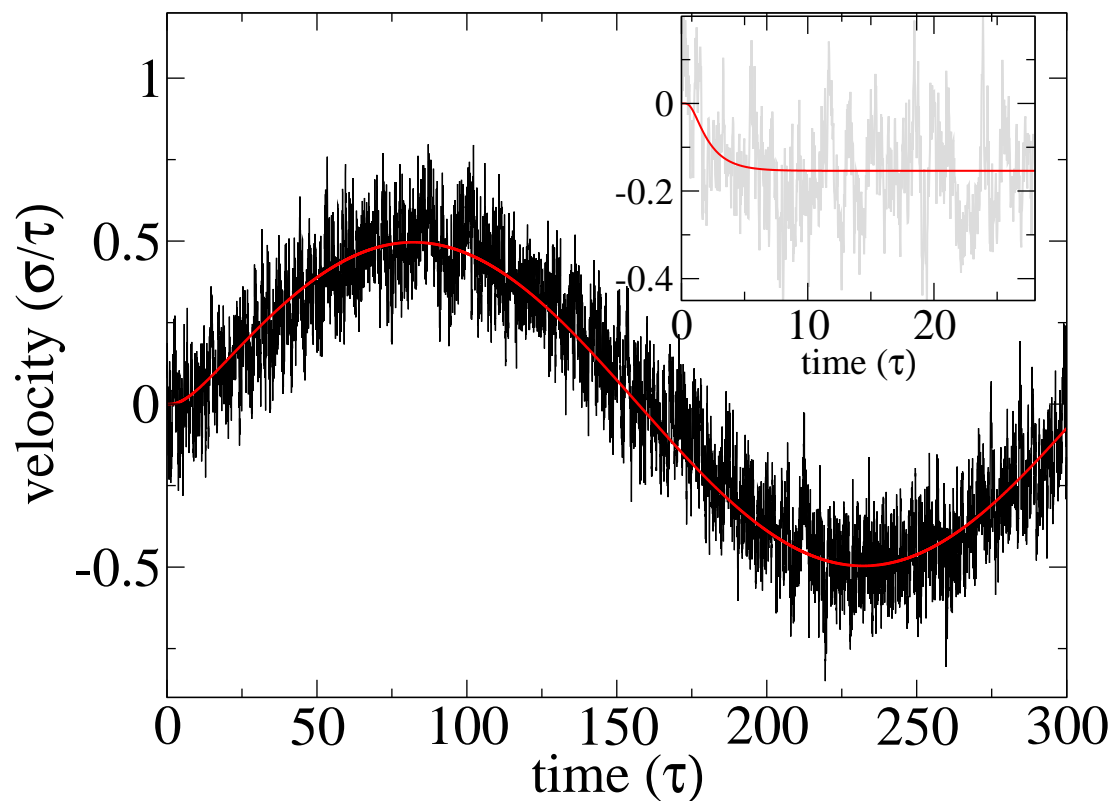
## Mass fluctuations

Standard deviation of mass  $Var[\rho] = \rho k_B T / (V c_T^2)$ ,  $c_T^2 = \left( \frac{\partial p}{\partial \rho} \right)_T$ ,  $\beta_T = (c_T^2 \rho)^{-1}$

Flexible TIP3P:  $c_T = 7.38(\epsilon_{OO}/m_O)^{1/2}$ ,  $\rho = 1.20m_O/\sigma_{OO}^3$

- ⑥  $V = 3.50 \times 6.18 \times 11.12 \sigma_{OO}^3$ :  
grand canonical:  $Var[\rho] = 0.0187$ , simulation:  $Var[\rho] = 0.020 \pm 0.002$
- ⑥  $V = 10.50 \times 6.18 \times 11.12 \sigma_{OO}^3$ :  
grand canonical:  $Var[\rho] = 0.0108$ , simulation:  $Var[\rho] = 0.011 \pm 0.005$

# Oscillatory Shear Flow



- ⑥ Time evolution of the velocity. For comparison, the deterministic Navier-Stokes solution is shown.





## ⑥ AdResS-HybridMD Scheme:

- △ We performed a **triple-scale simulation** of a molecular liquid.
- △ Length scales from the **micro-** to **macro-scale** are concurrently coupled.
- △ The triple-scale scheme is **robust** against the details of the mesoscopic model.
- △ The method allows us to perform **efficient molecular dynamics simulations** of molecular liquids in the **grand canonical ensemble** or under **non-equilibrium flows**.

## ⑥ Future work:

- △ Applications to study phenomena involving flow-matter interactions at multiple length scales.
- △ Grand canonical molecular simulations involving complex molecules.