α -Farey and α -Lüroth maps - new types of phase transitions

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Farey map

• The *Farey map* $F : [0,1] \rightarrow [0,1]$ is given by







Jump Transformation

- *G* is invariant with respect to the (finite) Gauss measure $d\mu(x) := ((1+x)\log 2)^{-1}d\lambda(x)$.
- *F* is invariant with respect to the (infinite) measure $dm(x) := 1/x \cdot d\lambda(x)$.
- Fix $A_1:=(1/2,1]$. For $x\in [0,1]\setminus \mathbb{Q}$ define the jump time

$$\varphi_{A_1}(x) := \inf \left\{ n \in \mathbb{N}_0 : F^n(x) \in A_1 \right\}$$

and let the jump transformation of the Farey map F with respect to A_1 for $x \in [0,1] \setminus \mathbb{Q}$ be given by

$$F_{A_1}(x) := F^{\varphi_{A_1}(x)+1}(x)$$



Continued Fractions: Sum-level Result

• n-th Sum-Level-Set:

$$\mathscr{C}_n := \left\{ x \in [a_1, \dots, a_k] : \sum_{i=1}^k a_i = n, \text{ for some } k \in \mathbb{N} \right\},$$

Theorem (K/Stratmann '10)

$$\lambda\left(\mathscr{C}_{n}\right)\sim \frac{\log 2}{\log n} \text{ and } \sum_{k=1}^{n}\lambda\left(\mathscr{C}_{k}\right)\sim \frac{n\log 2}{\log n}.$$

Proof.

Observe $F^{-n+1}([1/2,1]) = \mathscr{C}_n$ and use Infinite Ergodic Theory for the transfer operator \widehat{F} of F on $([0,1], \mathscr{B}, x^{-1}d\lambda(x))$.

Farey Spectrum

- $S: [0,1] \rightarrow [0,1]$ diff'able, $x \in [0,1]$, $\Lambda(S,x) := \lim_{n \to \infty} \frac{1}{n} \sum_{k=0}^{n-1} \log \left| S'(S^k(x)) \right|.$
- Lyapunov spectra (K./Stratmann '07)

$$\mathscr{L}_1(\alpha) := \{x \in [0,1] : \Lambda(F,x) = \alpha\}.$$



Gauss Spectrum

 Lyapunov spectrum for the Gauss map (Pollicott/Weiss '99, K./Stratmann '07, Fan/Liao/Wang/Wu '09)

$$\mathscr{L}_{3}(\alpha) := \{x \in [0,1] : \Lambda(G,x) = s\}.$$



Linearised versions: α -Lüroth and α -Farey maps



α-Lüroth map

- α-Farey map
- J. Lüroth. Über eine eindeutige Entwicklung von Zahlen in eine unendliche Reihe. *Math. Ann.* **21**:411–423, 1883.

Generating partition α

 countable partition α := {A_n : n ∈ ℕ} of [0,1] consisting of left open, right closed intervals; ordered from right to left, starting with A₁.

•
$$a_n := \lambda(A_n); t_n := \sum_{k=n}^{\infty} a_k.$$

• α -Lüroth map $L_{\alpha}(x) := \begin{cases} (t_n - x)/a_n & \text{for } x \in A_n, n \in \mathbb{N}, \\ 0 & \text{for } x = 0. \end{cases}$

• α-Farey map

$$F_{\alpha}(x) := \begin{cases} (1-x)/a_1 & \text{for } x \in A_1, \\ a_{n-1}(x-t_{n+1})/a_1 + t_n & \text{for } x \in A_n, n \ge 2, \\ 0 & \text{for } x = 0. \end{cases}$$

α -Lüroth and α -Farey

- λ is invariant with respect to L_{α} .
- L_{α} is the jump transformation of F_{α} with respect to A_1 .
- α is said to be of finite type if $\sum_{n=1}^{\infty} t_n < \infty$
- α is said to be of infinite type if $\sum_{n=1}^{\infty} t_n = \infty$
- α is called expansive of exponent $\theta \ge 0$ if $t_n = \psi(n)n^{-\theta}$, for all $n \in \mathbb{N}$ and some slowly varying function ψ . Then:

$$\lim_{n\to\infty}\frac{t_n}{t_{n+1}}=1 \text{ and } F'_{\alpha}\left(0+\right)=1$$

• α is said to be expanding if $\lim_{n\to\infty} t_n/t_{n+1} = \rho > 1$. Then:

$$F'_{\alpha}(0+) = \rho.$$

α -Lüroth and α -Farey

- $\exists v_{\alpha} \ll \lambda$ invariant with respect to F_{α} and density $\sum_{n=1}^{\infty} t_n / a_n \cdot \mathbf{1}_{A_n}$.
- $v_{\alpha}([0,1]) = +\infty \iff \alpha$ of infinite type.
- F_{α} and the tend map are topologically conjugate with conjugating homeomorphism given by (the α -Minkowski-? function)

$$heta_{lpha}\left(x
ight):=-2\sum\left(-1
ight)^{k}2^{-\sum_{i=1}^{k}\ell_{i}}$$

for $x = [\ell_1, \ell_2, ...]_{\alpha} = \sum_{n=1}^{\infty} (-1)^{n-1} (\prod_{i < n} a_{\ell_i}) t_{\ell_n}$ (α -Lüroth Expansion).

Examples for different expansive α



• *α*-sum-level sets

$$\mathscr{L}_n^{(\alpha)} := \left\{ x \in C_\alpha(\ell_1, \ell_2, \dots, \ell_k) : \sum_{i=1}^k \ell_i = n, \text{ for some } k \in \mathbb{N} \right\},$$

where

$$\mathcal{C}_{\alpha}(\ell_1,\ldots,\ell_k) := \{ x \in [0,1] : L^{i-1}_{\alpha}(x) \in \mathcal{A}_{l_i}, \forall i = 1,\ldots,k \}.$$

• Important fact: $\mathscr{L}_n^{(\alpha)} = F_\alpha^{-(n-1)}(A_1)$, for all $n \in \mathbb{N}$.

Theorem (K./Munday/Stratmann '11)

1 We have that $\sum_{n=1}^{\infty} \lambda(\mathscr{L}_n^{(\alpha)})$ diverges, and that

$$\lim_{n\to\infty}\lambda\left(\mathscr{L}_n^{(\alpha)}\right) = \begin{cases} 0, & \text{if } \alpha \text{ is of infinite type;} \\ \left(\sum_{k=1}^{\infty} t_k\right)^{-1}, & \text{if } \alpha \text{ is of finite type.} \end{cases}$$

2 Let α be either expansive of exponent $\theta \in [0,1]$ $(K_{\alpha} := \frac{1}{\Gamma(2-\theta)\Gamma(1+\theta)}, k_{\alpha} := \frac{1}{\Gamma(2-\theta)\Gamma(\theta)})$, or of finite type $K_{\alpha} := k_{\alpha} := 1$.

(a) Weak renewal law.
$$\sum_{k=1}^{n} \lambda \left(\mathscr{L}_{k}^{(\alpha)} \right) \sim K_{\alpha} \cdot n \cdot \left(\sum_{k=1}^{n} t_{k} \right)^{-1}$$
.
(b) Strong renewal law. $\lambda \left(\mathscr{L}_{n}^{(\alpha)} \right) \sim k_{\alpha} \cdot \left(\sum_{k=1}^{n} t_{k} \right)^{-1}$.

Proof of Part (1)

Fact (Renewal Equation)

For each $n \in \mathbb{N}$, we have that

$$\lambda\left(\mathscr{L}_{n}^{(\alpha)}\right) = \sum_{m=1}^{n} a_{m} \lambda\left(\mathscr{L}_{n-m}^{(\alpha)}\right).$$

Proof.

Proved by induction using linearity.

Proof of $\sum_{n=0}^{\infty} \lambda(\mathscr{L}_{n}^{(\alpha)})$ diverges. Define $a(s) := \sum_{n=1}^{\infty} a_{n}s^{n}$ and $\ell(s) := \sum_{m=0}^{\infty} \lambda\left(\mathscr{L}_{m}^{(\alpha)}\right)s^{m}$. Then for $s \in (0,1)$ we have that $\ell(s) - 1 = \ell(s)a(s)$, and hence, $\ell(s) = 1/(1 - a(s))$. Since a(1) = 1 we have $\lim_{s \nearrow 1} \ell(s) = \infty$

Proof of First Theorem

Proof Part (1).

Classical Renewal Theorem by Erdős, Pollard and Feller gives

$$\lim_{n\to\infty}\lambda(\mathscr{L}_n^{(\alpha)})=\frac{1}{\sum_{m=1}^{\infty}m\cdot a_m}=\frac{1}{\sum_{k=1}^{\infty}t_k}$$

(P. Erdős, H. Pollard, W. Feller. A property of power series with positive coefficients. *Bull. Amer. Math. Soc.* **55**:201-204, 1949)

Proof Part (2).

For the finite case consider part (1). For the expansive case apply a strong renewal theorems obtained in [K. B. Erickson. Strong renewal theorems with infinite mean. *Trans. Amer. Math. Soc.* **151**, 1970], [A. Garsia, J. Lamperti. A discrete renewal theorem with infinite mean. *Comment. Math. Helv.* **37**, 1963].

α -Farey Free Energy Function

• $S:[0,1] \rightarrow [0,1]$ diff'able, $x \in [0,1]$,

$$\Lambda(S,x) := \lim_{n \to \infty} \frac{1}{n} \sum_{k=0}^{n-1} \log \left| S'(S^k(x)) \right|.$$

• α -Farey Lyapunov spectrum, $s \in \mathbb{R}$,

$$\sigma_{\alpha}(s) := \dim_{H}(\{x \in [0,1] : \Lambda(F_{\alpha},x) = s\}).$$

• α -Farey free energy function $\nu : \mathbb{R} \to \mathbb{R} \cup \{\infty\}$

$$v(u) := \inf \left\{ r \in \mathbb{R} : \sum_{n=1}^{\infty} a_n^u \exp(-rn) \leq 1 \right\}.$$

• We say that F_{α} exhibits no phase transition if and only if v is diff'able everywhere.

Theorem (K./Munday/Stratmann '11)

Let α either expanding, or expansive and eventually decreasing. For $s_- := \inf\{-(\log a_n)/n : n \in \mathbb{N}\}$ and $s_+ := \sup\{-(\log a_n)/n : n \in \mathbb{N}\}$, we have that $\sigma_{\alpha}(s)$ vanishes outside the interval $[s_-, s_+]$ and for each $s \in (s_-, s_+)$, we have

$$\sigma_{\alpha}(s) = \inf_{u \in \mathbb{R}} \left(u + s^{-1} v(u) \right).$$

- **1** α expanding: F_{α} exhibits no phase transition. In particular, v is strictly decreasing and bijective.
- **2** α expansive of exponent θ and eventually decreasing: F_{α} exhibits no phase transition $\iff \alpha$ is of infinite type. In particular, $v \ge 0$ and $v|_{[1,\infty)} = 0$.

α -Lüroth Pressure

• α -Lüroth Lyapunov spectrum, $s \in \mathbb{R}$

$$\tau_{\alpha}(s) := \dim_{H}(\{x \in \mathscr{U} : \Lambda(L_{\alpha}, x) = s\}).$$

• α -Lüroth pressure function $p : \mathbb{R} \to \mathbb{R} \cup \{\infty\}$

$$p: u \mapsto \log \sum_{n=1}^{\infty} a_n^u.$$

• We say that L_{α} exhibits no phase transition if and only if the pressure function p is differentiable everywhere (that is, the right and left derivatives of p coincide everywhere, with the convention that $p'(u) = \infty$ if $p(u) = \infty$).

Theorem (K./Munday/Stratmann '11)

For $t_- := \min\{-\log a_n : n \in \mathbb{N}\}$ we have that τ_{α} vanishes on $(-\infty, t_-)$, and for each $s \in (t_-, \infty)$ we have

$$au_{lpha}(s) = \inf_{u \in \mathbb{R}} \left(u + s^{-1} p(u)
ight).$$

Moreover, $\lim_{s\to\infty} \tau_{\alpha}(s) = t_{\infty} := \inf\{r > 0 : \sum_{k=1}^{\infty} a_n^r < \infty\} \le 1$.

- **1** α expanding: L_{α} exhibits no phase transition and $t_{\infty} = 0$.
- **2** α expansive of exponent $\theta > 0$ and eventually decreasing: $t_{\infty} = 1/(1+\theta)$.

 L_{α} exhibits no phase trans. $\iff \sum_{n=1}^{\infty} \psi(n)^{1/(1+\theta)} \frac{\log n}{n} = \infty.$

a expansive of exponent θ = 0 and eventually decreasing: t_∞ = 1.
 L_α exhibits no phase trans. ⇔ Σ_{n=1}[∞] a_n log(a_n) = ∞.

Good set

Theorem (Munday '10)

The critical value t_{∞} is also equal to the Hausdorff dimension of the Good-type set $G_{\infty}^{(\alpha)}$ associated to L_{α} , given by

$$G_{\infty}^{(\alpha)} := \{ [\ell_1, \ell_2, \ldots]_{\alpha} : \lim_{n \to \infty} \ell_n = \infty \}.$$

If L_α exhibits a phase transition, that is ∑ a_n^{t_∞} < +∞ with finite right derivative t₀ in t_∞, then for t ∈ [t₀, +∞),

$$\tau_{\alpha}(t) = \frac{\log \sum_{n=1}^{\infty} a_n^{t_{\infty}}}{t} + t_{\infty}.$$

Expansive Example : The classical alternating Lüroth system



• For $\alpha_H := \{(1/(n+1), 1/n], n \in \mathbb{N}\}$ The figure shows the α_H -Farey free energy v (solid line), the α_H -Lüroth pressure function p (dashed line), and the associated dimension graphs σ_{α_H} and τ_{α_H} . Here, $t_- = \log 2, t_{\infty} = 1/2$ and $s_+ = (\log 6)/2$. We have $p(t_{\infty}) = \infty$, no phase transition for the α_H -Farey free energy function and the α_H -Lüroth pressure function.

Expansive Example: $a_n := \zeta (5/4)^{-1} n^{-5/4}$



 The Farey spectrum and the Lüroth spectrum intersect in a single point, for α expansive. The α-Farey free energy v (solid line), the α-Lüroth pressure function p (dashed line), and the associated dimension graphs for a_n := ζ (5/4)⁻¹ n^{-5/4}. Here, F_α exhibits no phase transition.

Expanding Example: $a_n := 2 \cdot 3^{-n}$



• The Farey spectrum is completely contained in the Lüroth spectrum, for α expanding. The α -Farey free energy v (solid line), the α -Lüroth pressure function p (dashed line), and the associated dimension graphs. The α -Farey system is given in this situation by the tent map with slopes 3 and -3/2.

Example for Lüroth Phase Transition $a_n := \frac{C}{n^2 \cdot (\log(n+5))^{12}}$



 Finite critical value p(t_∞) < ∞ with phase transition for the α-Lüroth pressure function and α expansive. The α-Lüroth pressure function p, and the associated dimension graphs. In this case t_∞ = 1/2 and p(1/2) < ∞ and L_α has a phase transition.

Examples: No Lüroth Phase Transition $a_n := \frac{C}{n^2 \cdot (\log(n+5))^4}$



 Finite critical value p(t_∞) < ∞ and no phase transition for the α-Lüroth pressure function and α expansive. The α-Lüroth pressure function p, and the associated dimension graphs for the α-Lüroth system. In this case t_∞ = 1/2 and p(1/2) < ∞, but L_α exhibits no phase transition.

Technical Lemma

Lemma

Let α be a partition such that $\lim_{n\to\infty} t_n/t_{n+1} = \rho \ge 1$ and such that α is either expanding, or expansive of exponent θ and eventually decreasing. Then:

 $\lim_{n \to \infty} \frac{\log a_n}{n} = \lim_{n \to \infty} \frac{\log t_n}{n} = -\log \rho. \ \alpha \text{ expansive with } \theta > 0,$ then $a_n \sim \theta n^{-1} t_n.$

2 If α expansive with $\theta = 0$, then we have $t_{\infty} = 1$.

- **3** If α is expanding, or expansive with $\theta > 0$, then $\lim_{n \to \infty} \frac{a_n}{a_{n+1}} = \rho$.
- **④** There exists a sequence $(\varepsilon_k)_{k \in \mathbb{N}}$, with $\lim_{k \to \infty} \varepsilon_k = 0$, such that for all *n* ∈ \mathbb{N} and *x* ∈ $\bigcup_{k > n} A_k$ we have that

$$\left|\frac{1}{n}\sum_{k=0}^{n-1}\log\left|F'_{\alpha}(F^{k}_{\alpha}(x))\right|-\log\rho\right|<\varepsilon_{n}.$$

Theorem (K./Jaerisch)

Consider the two potential functions $\varphi, \psi : \mathscr{U} \to \mathbb{R}$ given for $x \in A_n$, $n \in \mathbb{N}$, by $\varphi(x) := \log a_n$ and $\psi(x) := z_n$, for some fixed sequence $(z_n)_{n \in \mathbb{N}}$ of negative real numbers. For all $s \in \mathbb{R}$ we then have that

$$\dim_{H}\left\{x\in\mathscr{U}:\lim_{n\to\infty}\frac{\sum_{k=0}^{n-1}\psi(L_{\alpha}^{k}(x))}{\sum_{k=0}^{n-1}\varphi(L_{\alpha}^{k}(x))}=s\right\}\leq\max\{0,-t^{*}(-s)\}.$$

The function $t:\mathbb{R}\to\mathbb{R}\cup\{\infty\}$ is given by

$$t(v) := \inf \left\{ u \in \mathbb{R} : \sum_{n=1}^{\infty} a_n^u \exp(vz_n) \le 1 \right\}$$

and t^* is the Legendre transform of t.

Multifractal Formalism for Countable State Spaces II

Theorem (K./Jaerisch)

With

$$r_{-} := \inf \{-t^{+}(v) : v \in Int(dom(t))\},\\ r_{+} := \sup \{-t^{+}(v) : v \in Int(dom(t))\},\$$

we have for each $s \in (r_-, r_+)$,

$$\dim_{H}\left\{x\in\mathscr{U}:\lim_{n\to\infty}\frac{\sum_{k=0}^{n-1}\psi(L_{\alpha}^{k}(x))}{\sum_{k=0}^{n-1}\varphi(L_{\alpha}^{k}(x))}=s\right\}=-t^{*}(-s).$$

where t^+ denotes the right derivative of t, Int(A) denotes the interior of the set A, and $dom(t) := \{v \in \mathbb{R} : t(v) < +\infty\}.$

Proof of Theorem for α -Farey

• Set $z_n := -n$, then $v : u \mapsto \inf \{r \in \mathbb{R} : \sum_{n=1}^{\infty} a_n^u \exp(-rn) \le 1\}$ is the inverse of t.

•
$$s_{-} = 1/r_{+}$$
 and $s_{+} = 1/r_{-}$.

• For $s \in (s_-, s_+)$, it follows that

$$\begin{aligned} \sigma_{\alpha}(s) &= -t^* \left(-1/s\right) = \inf_{v \in \mathbb{R}} \left(t\left(v\right) + s^{-1}v\right) \\ &= \inf_{u \in \mathbb{R}} \left(u + s^{-1}\log\sum_{n=1}^{\infty} a_n^u\right) \end{aligned}$$

and $\sigma(s)$ vanishes outside of (s_-, s_+) .

Phase Transition for the α -Farey Free Energy

• Consider
$$Z(u,v) := \sum_{n=1}^{\infty} \exp\left(n\left(\frac{u\log a_n}{n} - v\right)\right)$$
.

- α expanding $\implies \forall u_0 \in \mathbb{R}\{Z(u_0, v) : v \in \mathbb{R}\} = (0, \infty) \implies \exists f(u_0) \text{ is unique solution of } Z(u_0, f(u_0)) = 1$. By the implicit function theorem there is no phase transition.
- α expansive
 - For *u* < 1 argue as above

• For
$$u \ge 1$$
 we have $\sum_{n=1}^{\infty} a_n^u e^{-wn} \begin{cases} <1 & \text{for } w \ge 0 \\ = \infty & \text{for } w < 0 \end{cases}$
 $v(u) = 0$

• Consider
$$f'(u) = \frac{\sum_{n=1}^{\infty} a_n^u e^{-f(u)n} \log a_n}{\sum_{n=1}^{\infty} n a_n^u e^{-f(u)n}}$$
 for $u \nearrow 1$.

- Infinite type: Denominator tends to ∞.
- $\lim_{u \nearrow 1} f'(u) = \lim_{u \nearrow 1} \sum_{n=1}^{\infty} \frac{\log a_n}{n} \frac{na_n e^{-f(u)n}}{\sum_{k=1}^{\infty} ka_k e^{-f(u)k}} = 0.$

Lemma

Let α be a partition which is either expanding, or expansive of exponent θ and eventually decreasing. With

$$\Pi(L_{\alpha},x) := \lim_{n \to \infty} \left(\sum_{k=0}^{n-1} \log \left| L'_{\alpha}(L^{k}_{\alpha}(x)) \right| \right) / \left(\sum_{k=0}^{n-1} N(L^{k}_{\alpha}(x)) \right),$$

we then have for each $s \ge 0$ that the sets

$$\{x \in \mathscr{U} : \Pi(L_{\alpha}, x) = s\}$$
 and $\{x \in \mathscr{U} : \Lambda(F_{\alpha}, x) = s\}$

coincide up to a countable set of points.

Proof of Theorem for α -Lüroth

- Set $z_n := -1$, $p : u \mapsto \log \sum_{n=1}^{\infty} a_n^u$ is the inverse of t.
- $t_-:=1/r_+=\inf\{-\log a_n:n\in\mathbb{N}\}$ and $t_+:=+\infty$
- For $s \in (t_-, +\infty)$, it follows that

$$egin{array}{rll} au_lpha(s) &=& -t^*(-1/s) = \inf_{v\in\mathbb{R}}ig(t\,(v)+s^{-1}vig) \ &=& \inf_{u\in\mathbb{R}}ig(u+s^{-1}\log\sum_{n=1}^\infty a_n^uig) \end{array}$$

and $\tau_{\alpha}(s)$ vanishes for $s < t_{-}$.

• For the right derivative of the pressure function p of L_{α} , we have that

$$p'(u) = \frac{\sum_{n=1}^{\infty} a_n^u \log a_n}{\sum_{n=1}^{\infty} a_n^u}$$

- Clearly, p is real-analytic on (t_{∞},∞) .
- Hence, we have that L_{α} exhibits no phase transition if and only if $\lim_{u \searrow t_{\infty}} -p'(u) = +\infty$.
- If α is expanding, then there is no phase transition. This follows, since, by the technical Lemma, we have that $p(u) < \infty$, for all u > 0. In particular, $t_{\infty} = 0$. If α is expansive with $\theta = 0$, we have by the Technical Lemma $t_{\infty} = 1$. Hence, $\lim_{u \searrow t_{\infty}} p'(u) = \infty$ if and only if $-\sum_{n=1}^{\infty} a_n \log(a_n) = \infty$.

Phase Transition for α -Lüroth

• If α is expansive such that $t_n = \psi(n)n^{-\theta}$, then the Technical Lemma implies that there exists ψ_0 such that $\psi_0(n) \sim \theta \psi(n)$ and $a_n = \psi_0(n)n^{-(1+\theta)}$. Consequently, we have that $t_{\infty} = 1/(1+\theta)$. Hence, we now observe that

$$\lim_{u \searrow t_{\infty}} -p'(u) = t_{\infty}^{-1} \lim_{u \searrow t_{\infty}} \frac{\sum_{n=1}^{\infty} \left(n^{-1-\theta} \psi_{0}(n) \right)^{u} \log \left(n(\psi_{0}(n))^{-\frac{1}{1+\theta}} \right)}{\sum_{n=1}^{\infty} \left(n^{-1-\theta} \psi_{0}(n) \right)^{u}}$$

- $\sum_{n=1}^{\infty} \psi(n)^{1/(1+\theta)} (\log n)/n < \infty \implies$ numerator and denominator both converge $\implies \lim_{u \searrow t_{\infty}} -p'(u) < \infty \implies$ phase transition.
- $\sum_{n=1}^{\infty} \psi(n)^{1/(1+\theta)} (\log n)/n = \infty$:
 - $\sum_{n=1}^{\infty} n^{-1} \psi_0(n)^{1/(1+\theta)} < \infty \implies \lim_{u \searrow t_{\infty}} -p'(u) = \infty.$
 - $\sum_{n=1}^{\infty} n^{-1} \psi_0(n)^{1/(1+\theta)} = \infty \implies$ $\forall k \in \mathbb{N} : \lim_{u \searrow t_{\infty}} (k^{-(1+\theta)} \psi_0(k))^u / \sum_{n=1}^{\infty} (n^{-(1+\theta)} \psi_0(n))^u = 0$ $\implies \lim_{u \searrow t_{\infty}} -p'(u) = \infty.$
- \implies no phase transition.

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