

Taming infinities

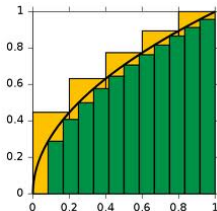
M. Hairer

University of Warwick

29.05.2014

Infinites and infinitesimals

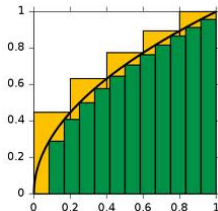
Bonaventura Cavalieri (1635): Computes areas by adding infinitesimals. (Even much earlier: Archimedes!)



Paul Guldin: “Things that do not exist, nor could they exist, cannot be compared.”

Infinites and infinitesimals

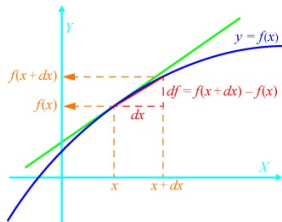
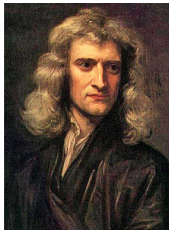
Bonaventura Cavalieri (1635): Computes areas by adding infinitesimals. (Even much earlier: Archimedes!)



Paul Guldin: “Things that do not exist, nor could they exist, cannot be compared.”

Calculus

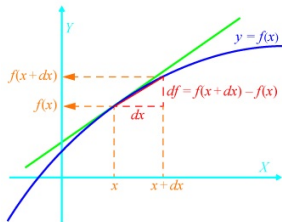
Leibniz and Newton (1680's): Compute derivatives by dividing infinitesimals.



Lord Berkeley: “They are neither finite quantities, nor quantities infinitely small, nor yet nothing. May we not call them the ghosts of departed quantities?”

Calculus

Leibniz and Newton (1680's): Compute derivatives by dividing infinitesimals.



Lord Berkeley: “They are neither finite quantities, nor quantities infinitely small, nor yet nothing. May we not call them the ghosts of departed quantities?”

Resolution

Cauchy, Bolzano, Weierstrass, etc (1820's onward): Rigorous formulation of limits without need for infinitesimals.

Dramatic consequences: Development of a consistent calculus allowing to model many physical processes and underpinning much of modern technology.

Moral of the story: Expressions like $\frac{df}{dx}$ can be given a meaning (as a real number) without having to give meaning to df and dx separately.

Resolution

Cauchy, Bolzano, Weierstrass, etc (1820's onward): Rigorous formulation of limits without need for infinitesimals.

Dramatic consequences: Development of a consistent calculus allowing to model many physical processes and underpinning much of modern technology.

Moral of the story: Expressions like $\frac{df}{dx}$ can be given a meaning (as a real number) without having to give meaning to df and dx separately.

Resolution

Cauchy, Bolzano, Weierstrass, etc (1820's onward): Rigorous formulation of limits without need for infinitesimals.

Dramatic consequences: Development of a consistent calculus allowing to model many physical processes and underpinning much of modern technology.

Moral of the story: Expressions like $\frac{df}{dx}$ can be given a meaning (as a real number) without having to give meaning to df and dx separately.

Quantum field theory

General methodology: “Guess” form of Lagrangian, predict outcomes of experiments as function of free parameters, perform experiments to determine them.

Problem: In QFT, when matching theory and experiments, infinities appear!

Cure: Discard infinities in a systematic way to extract finite parts.

Quantum field theory

General methodology: “Guess” form of Lagrangian, predict outcomes of experiments as function of free parameters, perform experiments to determine them.

Problem: In QFT, when matching theory and experiments, infinities appear!

Cure: Discard infinities in a systematic way to extract finite parts.

Quantum field theory

General methodology: “Guess” form of Lagrangian, predict outcomes of experiments as function of free parameters, perform experiments to determine them.

Problem: In QFT, when matching theory and experiments, infinities appear!

Cure: Discard infinities in a systematic way to extract finite parts.

Some reactions

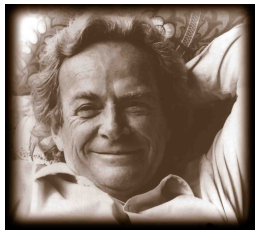
Not everybody liked these techniques...



“This is just not sensible mathematics. Sensible mathematics involves neglecting a quantity when it is small - not neglecting it just because it is infinitely great and you do not want it.” – Paul Dirac

More reactions

... not even those who developed them!

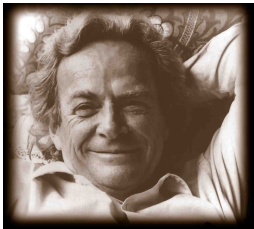


“The shell game that we play [...] is technically called ‘renormalization’. But no matter how clever the word, it is still what I would call a dippy process!” – Richard Feynman

However: Experimental verification of quantum electrodynamics predictions to within 9 digits of accuracy!

More reactions

... not even those who developed them!



“The shell game that we play [...] is technically called ‘renormalization’. But no matter how clever the word, it is still what I would call a dippy process!” – Richard Feynman

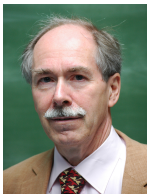
However: Experimental verification of quantum electrodynamics predictions to within 9 digits of accuracy!

Renormalisability

Some models are perturbatively renormalisable: at every order, parameters can be adjusted (in a diverging way!) to provide consistent answers.

Outcome: Theory with as many parameters as the naïve model.

Moral: “Form” of a model matters, not finiteness of constants.



't Hooft shows that the “standard model” is perturbatively renormalisable.

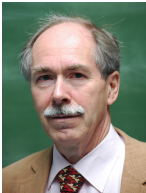
Despite investing billions (LHC), that model has not been faulted yet.

Renormalisability

Some models are perturbatively renormalisable: at every order, parameters can be adjusted (in a diverging way!) to provide consistent answers.

Outcome: Theory with as many parameters as the naïve model.

Moral: “Form” of a model matters, not finiteness of constants.



't Hooft shows that the “standard model” is perturbatively renormalisable.

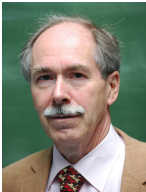
Despite investing billions (LHC), that model has not been faulted yet.

Renormalisability

Some models are perturbatively renormalisable: at every order, parameters can be adjusted (in a diverging way!) to provide consistent answers.

Outcome: Theory with as many parameters as the naïve model.

Moral: “Form” of a model matters, not finiteness of constants.



't Hooft shows that the “standard model” is perturbatively renormalisable.

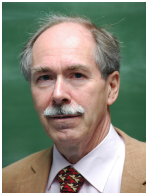
Despite investing billions (LHC), that model has not been faulted yet.

Renormalisability

Some models are perturbatively renormalisable: at every order, parameters can be adjusted (in a diverging way!) to provide consistent answers.

Outcome: Theory with as many parameters as the naïve model.

Moral: “Form” of a model matters, not finiteness of constants.



't Hooft shows that the “standard model” is perturbatively renormalisable.

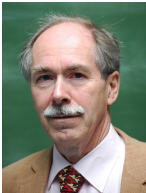
Despite investing billions (LHC), that model has not been faulted yet.

Renormalisability

Some models are perturbatively renormalisable: at every order, parameters can be adjusted (in a diverging way!) to provide consistent answers.

Outcome: Theory with as many parameters as the naïve model.

Moral: “Form” of a model matters, not finiteness of constants.



't Hooft shows that the “standard model” is perturbatively renormalisable.

Despite investing billions (LHC), that model has not been faulted yet.

Example

Try to define distribution “ $\eta(x) = \frac{1}{|x|}$ ”.

Problem: Integral of $1/|x|$ diverges, so we need to set “ $C = \infty$ ” to compensate!

Formal definition:

$$\eta_\chi(\varphi) = \int_{\mathbf{R}} \frac{\varphi(x) - \chi(x)\varphi(0)}{|x|} dx ,$$

for some smooth compactly supported cut-off χ with $\chi(0) = 1$.
Yields **one-parameter** family $c \mapsto \eta_c$ of models, but no canonical “choice of origin” for c .

Approximation: $1/(\varepsilon + |x|) - 2|\log \varepsilon| \delta(x)$ converges to η_c for some c .

Example

Try to define distribution “ $\eta(x) = \frac{1}{|x|} - C\delta(x)$ ”.

Problem: Integral of $1/|x|$ diverges, so we need to set “ $C = \infty$ ” to compensate!

Formal definition:

$$\eta_\chi(\varphi) = \int_{\mathbf{R}} \frac{\varphi(x) - \chi(x)\varphi(0)}{|x|} dx ,$$

for some smooth compactly supported cut-off χ with $\chi(0) = 1$.
Yields **one-parameter** family $c \mapsto \eta_c$ of models, but no canonical “choice of origin” for c .

Approximation: $1/(\varepsilon + |x|) - 2|\log \varepsilon| \delta(x)$ converges to η_c for some c .

Example

Try to define distribution “ $\eta(x) = \frac{1}{|x|} - C\delta(x)$ ”.

Problem: Integral of $1/|x|$ diverges, so we need to set “ $C = \infty$ ” to compensate!

Formal definition:

$$\eta_\chi(\varphi) = \int_{\mathbf{R}} \frac{\varphi(x) - \chi(x)\varphi(0)}{|x|} dx ,$$

for some smooth compactly supported cut-off χ with $\chi(0) = 1$.
Yields **one-parameter** family $c \mapsto \eta_c$ of models, but no canonical “choice of origin” for c .

Approximation: $1/(\varepsilon + |x|) - 2|\log \varepsilon| \delta(x)$ converges to η_c for some c .

Example

Try to define distribution “ $\eta(x) = \frac{1}{|x|} - C\delta(x)$ ”.

Problem: Integral of $1/|x|$ diverges, so we need to set “ $C = \infty$ ” to compensate!

Formal definition:

$$\eta_\chi(\varphi) = \int_{\mathbf{R}} \frac{\varphi(x) - \chi(x)\varphi(0)}{|x|} dx ,$$

for some smooth compactly supported cut-off χ with $\chi(0) = 1$.
Yields **one-parameter** family $c \mapsto \eta_c$ of models, but no canonical “choice of origin” for c .

Approximation: $1/(\varepsilon + |x|) - 2|\log \varepsilon| \delta(x)$ converges to η_c for some c .

Example

Try to define distribution “ $\eta(x) = \frac{1}{|x|} - C\delta(x)$ ”.

Problem: Integral of $1/|x|$ diverges, so we need to set “ $C = \infty$ ” to compensate!

Formal definition:

$$\eta_\chi(\varphi) = \int_{\mathbf{R}} \frac{\varphi(x) - \chi(x)\varphi(0)}{|x|} dx ,$$

for some smooth compactly supported cut-off χ with $\chi(0) = 1$.
Yields **one-parameter** family $c \mapsto \eta_c$ of models, but no canonical “choice of origin” for c .

Approximation: $1/(\varepsilon + |x|) - 2|\log \varepsilon| \delta(x)$ converges to η_c for some c .

Stochastics / Finance

Random walk: $W_{t+\varepsilon}^\varepsilon = W_t^\varepsilon + \sqrt{\varepsilon}\xi_t$ with $\{\xi_t\}$ independent identically distributed random variables, zero mean, unit variance.
Donsker's invariance principle: $W_t^\varepsilon \rightarrow W_t$ with W_t a Brownian motion.

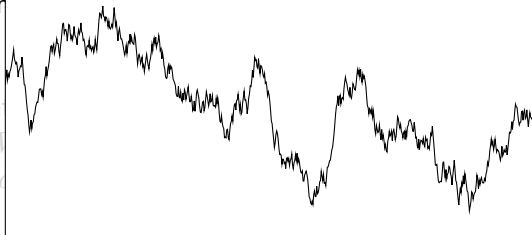
Simple asset price model: $S_{t+\varepsilon}^\varepsilon = S_t^\varepsilon(1 + \sqrt{\varepsilon}\xi_t)$. (So $\delta S_t^\varepsilon = S_t^\varepsilon \delta W_t^\varepsilon$.) **Formally**, one expects in the limit $\varepsilon \rightarrow 0$ to have $dS/dt = S dW/dt$, so that $S_t = S_0 \exp(W_t)$.

Wrong: The limit satisfies $S_t = S_0 \exp(W_t - t/2)$. (Can be "guessed" from $\mathbf{E}S_t = \mathbf{E}S_0$.)

Stochastics / Finance

Random walk: $W_t^\varepsilon = W_t^\varepsilon + \sqrt{\varepsilon} \xi_t$ with $\{\xi_t\}$ independent
identically distributed with mean 0 and variance ε .
Donsker's invariance principle: Brownian motion.

Sample path of Brownian motion



Simple asset price model: $\delta S_t^\varepsilon = S_t^\varepsilon \delta W_t^\varepsilon$
 $dS/dt = S \sigma^2/2$ to have

Wrong: The limit satisfies $S_t = S_0 \exp(W_t - t/2)$. (Can be
“guessed” from $\mathbf{E}S_t = \mathbf{E}S_0$.)

Stochastics / Finance

Random walk: $W_{t+\varepsilon}^\varepsilon = W_t^\varepsilon + \sqrt{\varepsilon}\xi_t$ with $\{\xi_t\}$ independent identically distributed random variables, zero mean, unit variance.
Donsker's invariance principle: $W_t^\varepsilon \rightarrow W_t$ with W_t a Brownian motion.

Simple asset price model: $S_{t+\varepsilon}^\varepsilon = S_t^\varepsilon(1 + \sqrt{\varepsilon}\xi_t)$. (So $\delta S_t^\varepsilon = S_t^\varepsilon \delta W_t^\varepsilon$.) **Formally**, one expects in the limit $\varepsilon \rightarrow 0$ to have $dS/dt = S dW/dt$, so that $S_t = S_0 \exp(W_t)$.

Wrong: The limit satisfies $S_t = S_0 \exp(W_t - t/2)$. (Can be "guessed" from $\mathbf{E}S_t = \mathbf{E}S_0$.)

Stochastics / Finance

Random walk: $W_{t+\varepsilon}^\varepsilon = W_t^\varepsilon + \sqrt{\varepsilon}\xi_t$ with $\{\xi_t\}$ independent identically distributed random variables, zero mean, unit variance.
Donsker's invariance principle: $W_t^\varepsilon \rightarrow W_t$ with W_t a Brownian motion.

Simple asset price model: $S_{t+\varepsilon}^\varepsilon = S_t^\varepsilon(1 + \sqrt{\varepsilon}\xi_t)$. (So $\delta S_t^\varepsilon = S_t^\varepsilon \delta W_t^\varepsilon$.) **Formally**, one expects in the limit $\varepsilon \rightarrow 0$ to have $dS/dt = S dW/dt$, so that $S_t = S_0 \exp(W_t)$.

Wrong: The limit satisfies $S_t = S_0 \exp(W_t - t/2)$. (Can be "guessed" from $\mathbf{E}S_t = \mathbf{E}S_0$.)

What went wrong??

Problem: While S^ε converges to a limit and dW^ε/dt converges to a limit, these are **too rough** for their product to be well-posed.

In general $(f, \xi) \mapsto f \cdot \xi$ well-posed on $C^\alpha \times C^\beta$ if and only if $\alpha + \beta > 0$. Here: **just below borderline**.

Consequence: Limit depends on details of discretisation. For example, if we set instead $S_{t+\varepsilon} = \frac{S_t + S_{t+\varepsilon}}{2} (1 + \sqrt{\varepsilon} \xi_t)$, then $S_t^\varepsilon \rightarrow S_0 \exp(W_t)$ as expected. In general: one-parameter family $S_0 \exp(W_t - ct)$ with $c \in \mathbf{R}$.

Moral: For singular objects, details of the approximation may matter.

What went wrong??

Problem: While S^ε converges to a limit and dW^ε/dt converges to a limit, these are **too rough** for their product to be well-posed.

In general $(f, \xi) \mapsto f \cdot \xi$ **well-posed** on $C^\alpha \times C^\beta$ if and **only if** $\alpha + \beta > 0$. Here: **just below borderline**.

Consequence: Limit depends on details of discretisation. For example, if we set instead $S_{t+\varepsilon} = \frac{S_t + S_{t+\varepsilon}}{2} (1 + \sqrt{\varepsilon} \xi_t)$, then $S_t^\varepsilon \rightarrow S_0 \exp(W_t)$ as expected. In general: one-parameter family $S_0 \exp(W_t - ct)$ with $c \in \mathbf{R}$.

Moral: For singular objects, details of the approximation may matter.

What went wrong??

Problem: While S^ε converges to a limit and dW^ε/dt converges to a limit, these are **too rough** for their product to be well-posed.

In general $(f, \xi) \mapsto f \cdot \xi$ **well-posed** on $C^\alpha \times C^\beta$ if and **only if** $\alpha + \beta > 0$. Here: **just below borderline**.

Consequence: Limit depends on details of discretisation. For example, if we set instead $S_{t+\varepsilon} = \frac{S_t + S_{t+\varepsilon}}{2} (1 + \sqrt{\varepsilon} \xi_t)$, then $S_t^\varepsilon \rightarrow S_0 \exp(W_t)$ as expected. In general: one-parameter family $S_0 \exp(W_t - ct)$ with $c \in \mathbf{R}$.

Moral: For singular objects, details of the approximation may matter.

What went wrong??

Problem: While S^ε converges to a limit and dW^ε/dt converges to a limit, these are **too rough** for their product to be well-posed.

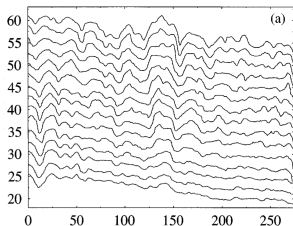
In general $(f, \xi) \mapsto f \cdot \xi$ **well-posed** on $C^\alpha \times C^\beta$ if and **only if** $\alpha + \beta > 0$. Here: **just below borderline**.

Consequence: Limit depends on details of discretisation. For example, if we set instead $S_{t+\varepsilon} = \frac{S_t + S_{t+\varepsilon}}{2} (1 + \sqrt{\varepsilon} \xi_t)$, then $S_t^\varepsilon \rightarrow S_0 \exp(W_t)$ as expected. In general: one-parameter family $S_0 \exp(W_t - ct)$ with $c \in \mathbf{R}$.

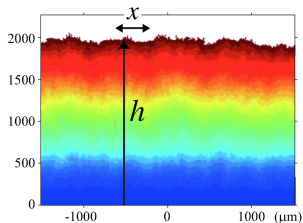
Moral: For singular objects, details of the approximation may matter.

Universality

Universality: Many random systems “look the same” and are “scale invariant” when viewed at scales much larger than that of the mechanism producing them, provided that they share some basic features:



Maunuksela & AI, PRL

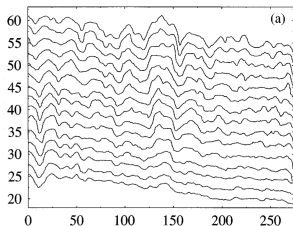


Takeuchi & AI, Sci. Rep.

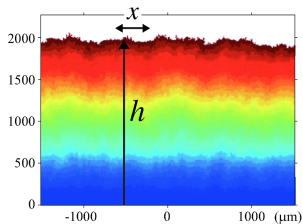
Very poorly understood in many cases!

Universality

Universality: Many random systems “look the same” and are “scale invariant” when viewed at scales much larger than that of the mechanism producing them, provided that they share some basic features:



Maunuksela & AI, PRL



Takeuchi & AI, Sci. Rep.

Very poorly understood in many cases!

Crossover regimes

Described by simple “normal form” equations:

$$\begin{aligned}\partial_t h &= \partial_x^2 h + (\partial_x h)^2 + \xi - C, & (\text{KPZ}; d = 1) \\ \partial_t \Phi &= -\Delta(\Delta\Phi + C\Phi - \Phi^3) + \nabla\xi. & (\Phi^4; d = 2, 3)\end{aligned}$$

Here ξ is **space-time white noise** (think of independent random variables at every space-time point).

KPZ: universal model for weakly asymmetric interface growth.

Φ^4 : universal model for phase coexistence near mean-field.

Problem: **red** terms ill-posed, requires $C = \infty!!$

Crossover regimes

Described by simple “normal form” equations:

$$\begin{aligned}\partial_t h &= \partial_x^2 h + (\partial_x h)^2 + \xi - C, & (\text{KPZ}; d = 1) \\ \partial_t \Phi &= -\Delta(\Delta\Phi + C\Phi - \Phi^3) + \nabla\xi. & (\Phi^4; d = 2, 3)\end{aligned}$$

Here ξ is **space-time white noise** (think of independent random variables at every space-time point).

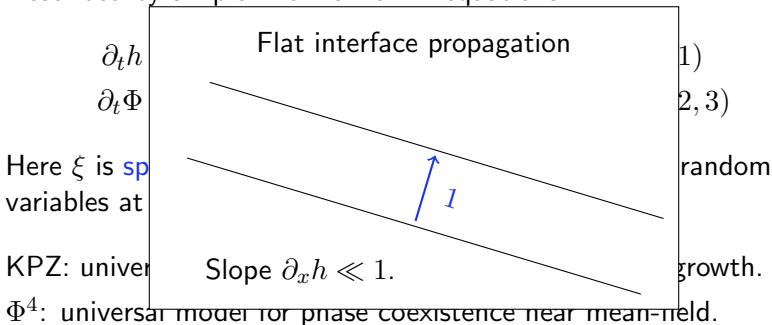
KPZ: universal model for weakly asymmetric interface growth.

Φ^4 : universal model for phase coexistence near mean-field.

Problem: **red** terms ill-posed, requires $C = \infty!!$

Crossover regimes

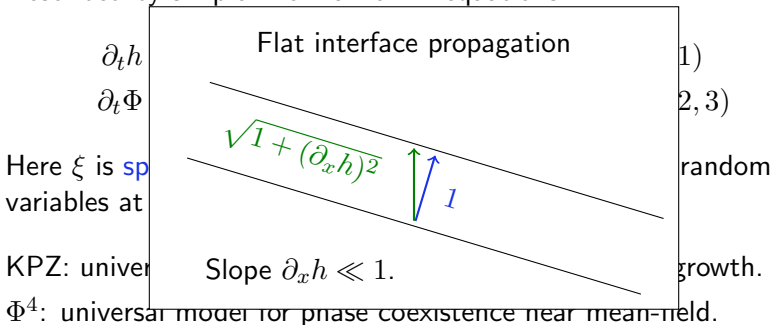
Described by simple “normal form” equations:



Problem: red terms ill-posed, requires $C = \infty!!$

Crossover regimes

Described by simple “normal form” equations:



Here ξ is sp variables at

KPZ: univer

Φ^4 : universal model for phase coexistence near mean-field.

Problem: red terms ill-posed, requires $C = \infty!!$

Crossover regimes

Described by simple “normal form” equations:

$$\begin{aligned}\partial_t h &= \partial_x^2 h + (\partial_x h)^2 + \xi - C, & (\text{KPZ}; d = 1) \\ \partial_t \Phi &= -\Delta(\Delta\Phi + C\Phi - \Phi^3) + \nabla\xi. & (\Phi^4; d = 2, 3)\end{aligned}$$

Here ξ is **space-time white noise** (think of independent random variables at every space-time point).

KPZ: universal model for weakly asymmetric interface growth.

Φ^4 : universal model for phase coexistence near mean-field.

Problem: **red** terms ill-posed, requires $C = \infty!!$

Well-posedness results

Write ξ_ε for mollified version of space-time white noise. Consider

$$\partial_t h = \partial_x^2 h + (\partial_x h)^2 - C_\varepsilon + \xi_\varepsilon, \quad (d = 1)$$

$$\partial_t \Phi = -\Delta(\Delta\Phi + C_\varepsilon\Phi - \Phi^3) + \nabla\xi_\varepsilon, \quad (d = 2, 3)$$

(Periodic boundary conditions on torus / circle.)

Theorem (H. 2013): There are $C_\varepsilon \rightarrow \infty$ so that solutions converge to a limit independent of the regularisation. (The constants themselves do depend on that choice.) For KPZ, limit coincides with Cole-Hopf solution (if C_ε is chosen appropriately).

Corollary of proof: Rates of convergence, precise local description of limit, suitable continuity, etc.

Well-posedness results

Write ξ_ε for mollified version of space-time white noise. Consider

$$\partial_t h = \partial_x^2 h + (\partial_x h)^2 - C_\varepsilon + \xi_\varepsilon, \quad (d = 1)$$

$$\partial_t \Phi = -\Delta(\Delta\Phi + C_\varepsilon\Phi - \Phi^3) + \nabla\xi_\varepsilon, \quad (d = 2, 3)$$

(Periodic boundary conditions on torus / circle.)

Theorem (H. 2013): There are $C_\varepsilon \rightarrow \infty$ so that solutions converge to a limit **independent** of the regularisation. (The constants themselves **do** depend on that choice.) For KPZ, limit coincides with Cole-Hopf solution (if C_ε is chosen appropriately).

Corollary of proof: Rates of convergence, precise local description of limit, suitable continuity, etc.

Well-posedness results

Write ξ_ε for mollified version of space-time white noise. Consider

$$\partial_t h = \partial_x^2 h + (\partial_x h)^2 - C_\varepsilon + \xi_\varepsilon, \quad (d = 1)$$

$$\partial_t \Phi = -\Delta(\Delta\Phi + C_\varepsilon\Phi - \Phi^3) + \nabla\xi_\varepsilon, \quad (d = 2, 3)$$

(Periodic boundary conditions on torus / circle.)

Theorem (H. 2013): There are $C_\varepsilon \rightarrow \infty$ so that solutions converge to a limit **independent** of the regularisation. (The constants themselves **do** depend on that choice.) For KPZ, limit coincides with Cole-Hopf solution (if C_ε is chosen appropriately).

Corollary of proof: Rates of convergence, precise local description of limit, suitable continuity, etc.

Universality result for KPZ

(In progress; joint with J. Quastel.) Consider

$$\partial_t h = \partial_x^2 h + \sqrt{\varepsilon} P(\partial_x h) + \xi ,$$

with P **even polynomial**, ξ **smooth** space-time Gaussian field with compactly supported correlations.

Theorem: As $\varepsilon \rightarrow 0$, there is a choice of $C_\varepsilon \sim \varepsilon^{-1}$ such that $\varepsilon^{1/2} h(\varepsilon^{-1} x, \varepsilon^{-2} t) - C_\varepsilon t$ converges to solutions to $(\text{KPZ})_\lambda$ with λ depending in a non-trivial way on all coefficients of P .

Remark: Convergence to KPZ with $\lambda \neq 0$ even if $P(u) = u^4!!$

Universality result for KPZ

(In progress; joint with J. Quastel.) Consider

$$\partial_t h = \partial_x^2 h + \sqrt{\varepsilon} P(\partial_x h) + \xi ,$$

with P **even polynomial**, ξ **smooth** space-time Gaussian field with compactly supported correlations.

Theorem: As $\varepsilon \rightarrow 0$, there is a choice of $C_\varepsilon \sim \varepsilon^{-1}$ such that $\varepsilon^{1/2} h(\varepsilon^{-1} x, \varepsilon^{-2} t) - C_\varepsilon t$ converges to solutions to $(\text{KPZ})_\lambda$ with λ depending in a non-trivial way on **all** coefficients of P .

Remark: Convergence to KPZ with $\lambda \neq 0$ even if $P(u) = u^4!!$

Universality result for KPZ

(In progress; joint with J. Quastel.) Consider

$$\partial_t h = \partial_x^2 h + \sqrt{\varepsilon} P(\partial_x h) + \xi,$$

with P even polynomial, ξ smooth space-time Gaussian field with compactly supported Nonlinearity $\lambda(\partial_x h)^2$

Theorem: As $\varepsilon \rightarrow 0$, there is a choice of $C_\varepsilon \sim \varepsilon^{-1}$ such that $\varepsilon^{1/2} h(\varepsilon^{-1} x, \varepsilon^{-2} t) - C_\varepsilon t$ converges to solutions to $(\text{KPZ})_\lambda$ with λ depending in a non-trivial way on all coefficients of P .

Remark: Convergence to KPZ with $\lambda \neq 0$ even if $P(u) = u^4!!$

Universality result for KPZ

(In progress; joint with J. Quastel.) Consider

$$\partial_t h = \partial_x^2 h + \sqrt{\varepsilon} P(\partial_x h) + \xi ,$$

with P **even polynomial**, ξ **smooth** space-time Gaussian field with compactly supported correlations.

Theorem: As $\varepsilon \rightarrow 0$, there is a choice of $C_\varepsilon \sim \varepsilon^{-1}$ such that $\varepsilon^{1/2} h(\varepsilon^{-1} x, \varepsilon^{-2} t) - C_\varepsilon t$ converges to solutions to $(\text{KPZ})_\lambda$ with λ depending in a non-trivial way on **all** coefficients of P .

Remark: Convergence to KPZ with $\lambda \neq 0$ even if $P(u) = u^4!!$

Universality result for KPZ

(In progress; joint with J. Quastel.) Consider

$$\partial_t h = \partial_x^2 h + \sqrt{\varepsilon} P(\partial_x h) + \xi ,$$

with P **even polynomial**, ξ **smooth** space-time Gaussian field with compactly supported correlations.

Theorem: As $\varepsilon \rightarrow 0$, there is a choice of $C_\varepsilon \sim \varepsilon^{-1}$ such that $\varepsilon^{1/2} h(\varepsilon^{-1} x, \varepsilon^{-2} t) - C_\varepsilon t$ converges to solutions to $(\text{KPZ})_\lambda$ with λ depending in a non-trivial way on **all** coefficients of P .

Remark: Convergence to KPZ with $\lambda \neq 0$ even if $P(u) = u^4!!$

Main Idea

Problem: Solutions are not smooth.

Insight: What is “smoothness”? Proximity to polynomials; we know how to multiply these...

Idea: Replace polynomials by a (finite / countable) collection of **taylor-made** space-time functions / distributions with similar algebraic / analytic properties. Depends on the realisation of the noise, but not on “details” of the equation. (Values of constants, initial condition, boundary conditions, etc.)

Amazing fact: If we chose the objects replacing polynomials in a smart way, these very singular solutions are “smooth”!

Main Idea

Problem: Solutions are not smooth.

Insight: What is “smoothness”? Proximity to polynomials; we know how to multiply these...

Idea: Replace polynomials by a (finite / countable) collection of *taylor-made* space-time functions / distributions with similar algebraic / analytic properties. Depends on the realisation of the noise, but not on “details” of the equation. (Values of constants, initial condition, boundary conditions, etc.)

Amazing fact: If we chose the objects replacing polynomials in a smart way, these very singular solutions are “smooth”!

Main Idea

Problem: Solutions are not smooth.

Insight: What is “smoothness”? Proximity to polynomials; we know how to multiply these...

Idea: Replace polynomials by a (finite / countable) collection of **taylor-made** space-time functions / distributions with similar algebraic / analytic properties. Depends on the realisation of the noise, but not on “details” of the equation. (Values of constants, initial condition, boundary conditions, etc.)

Amazing fact: If we chose the objects replacing polynomials in a smart way, these very singular solutions are “smooth”!

Main Idea

Problem: Solutions are not smooth.

Insight: What is “smoothness”? Proximity to polynomials; we know how to multiply these...

Idea: Replace polynomials by a (finite / countable) collection of **taylor-made** space-time functions / distributions with similar algebraic / analytic properties. Depends on the realisation of the noise, but not on “details” of the equation. (Values of constants, initial condition, boundary conditions, etc.)

Amazing fact: If we chose the objects replacing polynomials in a smart way, these very singular solutions are “smooth”!

General picture

Method of proof: Build objects for the following diagram:

$$\begin{array}{ccccc}
 \mathcal{F} \times \mathcal{M} \times C^\alpha(\mathbf{R}^d) & \xrightarrow{\mathcal{S}_A} & \mathcal{D}^\gamma \\
 \updownarrow \mathcal{F} & & \downarrow \mathcal{R} \\
 \mathcal{F} \times \text{?} \times C^\alpha(\mathbf{R}^d) & \xrightarrow{\mathcal{S}_C} & \mathcal{S}'(\mathbf{R}^{d+1}) \\
 \updownarrow \mathcal{R} & & \updownarrow \mathcal{U} \\
 \mathcal{F} \times \xi_\varepsilon \times u_0 & &
 \end{array}$$

The diagram shows a commutative-like structure. The top row is $\mathcal{F} \times \mathcal{M} \times C^\alpha(\mathbf{R}^d) \xrightarrow{\mathcal{S}_A} \mathcal{D}^\gamma$. The bottom row is $\mathcal{F} \times \text{?} \times C^\alpha(\mathbf{R}^d) \xrightarrow{\mathcal{S}_C} \mathcal{S}'(\mathbf{R}^{d+1})$. A vertical arrow labeled Ψ points from the ? to \mathcal{M} . Vertical arrows labeled \mathcal{U} point from \mathcal{D}^γ to $\mathcal{S}'(\mathbf{R}^{d+1})$ and from $C^\alpha(\mathbf{R}^d)$ to u_0 . Vertical arrows labeled \mathcal{F} and \mathcal{R} connect the \mathcal{F} terms. Self-loops labeled \mathcal{R} are on \mathcal{M} and \mathcal{F} .

\mathcal{F} : Formal right-hand side of the equation.

\mathcal{S}_C : Classical solution to the PDE with smooth input.

\mathcal{S}_A : Abstract fixed point: locally jointly continuous!

General picture

Method of proof: Build objects for the following diagram:

$$\begin{array}{ccccc}
 \mathcal{F} \times \mathcal{M} \times C^\alpha(\mathbf{R}^d) & \xrightarrow{S_A} & \mathcal{D}^\gamma & & \\
 \updownarrow & & \cdot & & \downarrow \mathcal{R} \\
 \mathcal{F} \times \text{?} \times C^\alpha(\mathbf{R}^d) & \xrightarrow{S_C} & \mathcal{S}'(\mathbf{R}^{d+1}) & & \\
 \updownarrow & & \Psi & & \\
 \mathcal{R} & & \xi_\varepsilon & & u_0
 \end{array}$$

\mathcal{R} (curved arrow from \mathcal{M} to \mathcal{F})
 Ψ (red arrow from ? to \mathcal{M})
 \mathcal{R} (curved arrow from \mathcal{F} to \mathcal{F})

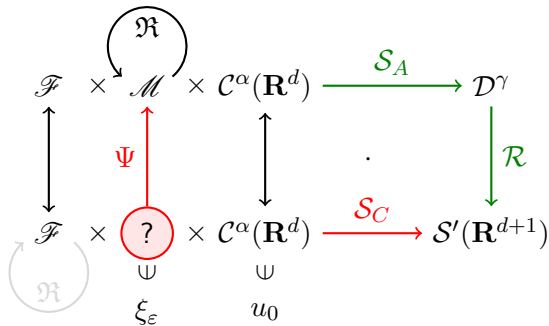
\mathcal{F} : Formal right-hand side of the equation.

S_C : Classical solution to the PDE with smooth input.

S_A : Abstract fixed point: locally **jointly continuous!**

General picture

Method of proof: Build objects for the following diagram:



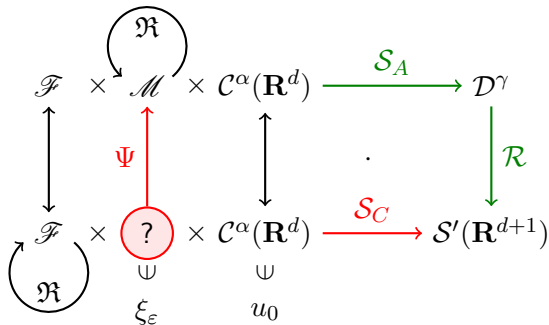
\mathcal{F} : Formal right-hand side of the equation.

S_C : Classical solution to the PDE with smooth input.

S_A : Abstract fixed point: locally **jointly continuous!**

General picture

Method of proof: Build objects for the following diagram:



\mathcal{F} : Formal right-hand side of the equation.

\mathcal{S}_C : Classical solution to the PDE with smooth input.

\mathcal{S}_A : Abstract fixed point: locally **jointly continuous!**

General picture

Method of proof: Build objects for the following diagram:

$$\begin{array}{ccccc} \mathcal{F} \times \mathcal{M} \times \mathcal{C}^\alpha(\mathbf{R}^d) & \xrightarrow{\mathcal{S}_A} & \mathcal{D}^\gamma & & \\ \updownarrow & & \cdot & & \downarrow \mathcal{R} \\ \mathcal{F} \times \text{?} \times \mathcal{C}^\alpha(\mathbf{R}^d) & \xrightarrow{\mathcal{S}_C} & \mathcal{S}'(\mathbf{R}^{d+1}) & & \\ \updownarrow & & \downarrow \Psi & & \\ \mathcal{F} \times \xi_\epsilon & & u_0 & & \end{array}$$

The diagram consists of two rows of objects. The top row is $\mathcal{F} \times \mathcal{M} \times \mathcal{C}^\alpha(\mathbf{R}^d)$ and the bottom row is $\mathcal{F} \times \text{?} \times \mathcal{C}^\alpha(\mathbf{R}^d)$. A red arrow labeled Ψ points from the question mark to \mathcal{M} . A green arrow labeled \mathcal{S}_A points from the top row to \mathcal{D}^γ . A red arrow labeled \mathcal{S}_C points from the bottom row to $\mathcal{S}'(\mathbf{R}^{d+1})$. A green arrow labeled \mathcal{R} points from \mathcal{D}^γ to $\mathcal{S}'(\mathbf{R}^{d+1})$. Vertical double-headed arrows connect the first and third columns. A black circle with \mathfrak{R} is around \mathcal{M} and another around \mathcal{F} . Below ξ_ϵ and u_0 are the labels Ψ and Ψ respectively.

Strategy: find $M_\epsilon \in \mathfrak{R}$ such that $M_\epsilon \Psi(\xi_\epsilon)$ converges.

Some concluding remarks

1. Diverging terms can (sometimes) be **cured** by suitable **counterterms**, in probabilistic models, not just in QFT.
2. Forces one to deal with families of models parametrised by constants defined “up to an infinite part”. Not a problem as soon as “observables” are finite and answers are consistent...
3. Goal: obtain “universality” results for models from statistical mechanics with tuneable parameters. (Current theory works well for continuous rather than discrete models.)

Thank you for your attention!

Some concluding remarks

1. Diverging terms can (sometimes) be **cured** by suitable **counterterms**, in probabilistic models, not just in QFT.
2. Forces one to deal with families of models parametrised by constants defined “up to an infinite part”. Not a problem as soon as “observables” are finite and answers are consistent...
3. Goal: obtain “universality” results for models from statistical mechanics with tuneable parameters. (Current theory works well for continuous rather than discrete models.)

Thank you for your attention!

Some concluding remarks

1. Diverging terms can (sometimes) be **cured** by suitable **counterterms**, in probabilistic models, not just in QFT.
2. Forces one to deal with families of models parametrised by constants defined “up to an infinite part”. Not a problem as soon as “observables” are finite and answers are consistent...
3. Goal: obtain “universality” results for models from statistical mechanics with tuneable parameters. (Current theory works well for continuous rather than discrete models.)

Thank you for your attention!

Some concluding remarks

1. Diverging terms can (sometimes) be **cured** by suitable **counterterms**, in probabilistic models, not just in QFT.
2. Forces one to deal with families of models parametrised by constants defined “up to an infinite part”. Not a problem as soon as “observables” are finite and answers are consistent...
3. Goal: obtain “universality” results for models from statistical mechanics with tuneable parameters. (Current theory works well for continuous rather than discrete models.)

Thank you for your attention!