#### Complex Gaussian multiplicative chaos

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#### Outline

1 Framework and Real Gaussian multiplicative Chaos

2 Complex Gaussian multiplicative chaos

3 Motivations from Liouville Quantum Gravity

## Gaussian multiplicative Chaos

We want to give a precise meaning to distributions  $M^{\gamma,\beta}$  defined formally by:

$$M^{\gamma,\beta}(A) = \int_A e^{\gamma X(x) + i\beta Y(x)} dx, \ A \subset D$$

where:

- $\gamma, \beta \geq 0$
- $D \subset \mathbb{C}$  a bounded domain
- X, Y two centered independent GFF with covariance given by:

$$E[X(x)X(y)] = G_D(x,y) \underset{|x-y| \to 0}{\sim} \ln \frac{1}{|x-y|}$$

where  $G_D$  is the Green kernel:

$$-\Delta_{V}G_{D}(x,y)=2\pi\delta_{X}$$



#### Framework

We consider a family of centered Gaussian processes  $(X_{\varepsilon}(x))_{x\in D}$   $(\varepsilon \leq 1)$ :

- Covariance:  $E[X_{\varepsilon}(x)X_{\varepsilon}(y)] \sim \ln \frac{1}{|x-y|+\varepsilon} \underset{\varepsilon \to 0}{\to} G_D(x,y)$
- Variance:  $E[X_{\varepsilon}(x)^2] = \ln \frac{1}{\varepsilon} + \ln C(x,D) + o(1)$  where C(x,D) conformal radius.
- $\varepsilon \mapsto X_{\varepsilon}$  independent increments

Same for  $(Y_{\varepsilon}(x))_{x \in D}$  ( $\varepsilon \leq 1$ ) independent from X.

## Gaussian multiplicative Chaos: notations

We define:

$$M_{\varepsilon}^{\gamma,\beta}(A) = \int_A e^{\gamma X_{\varepsilon}(x) + i\beta Y_{\varepsilon}(x)} dx, \ A \subset D$$

Observe that:

$$M_{\varepsilon}^{\gamma,0}(A) = \int_A e^{\gamma X_{\varepsilon}(x)} dx, \ A \subset D$$

#### Other Frameworks

One can also work with other "smooth" cut-offs

- 1985: Kahane, H¹-basis decomposition
- 2006-2008: Robert, V., general convolutions
- 2008: Duplantier, Sheffield, circle averages

In fact, one can work with any log-correlated field in any dimension (Kahane, 1985, Robert, V., 2006, 2008): see our review with Rhodes.

## Gaussian multiplicative chaos: $\beta = 0$

#### Theorem (Kahane, 1985)

There exists a random measure  $M^{\gamma,0}$  such that following limit exists almost surely in the space of Radon measures:

$$\varepsilon^{\frac{\gamma^2}{2}}M_{\varepsilon}^{\gamma,0}(dx)\underset{\varepsilon\to 0}{\to}M^{\gamma,0}(dx).$$

 $M^{\gamma,0}$  is called Gaussian multiplicative chaos associated to the Green kernel.

#### Remark

When J.P. Kahane meets Paul Levy...



## Gaussian multiplicative chaos: $\beta = 0$

#### Theorem (Kahane, 1985)

The measure  $M^{\gamma,0}$  is different from 0 if and only if  $\gamma < 2$ .

#### Theorem (Kahane, 1985)

For  $\gamma <$  2, the measure  $M^{\gamma,0}$  "lives" almost surely on a set of Hausdorff dimension  $2-\frac{\gamma^2}{2}$  (the set of  $\gamma$ -thick points).

## Density of Gaussian multiplicative chaos with respect to $\boldsymbol{\gamma}$

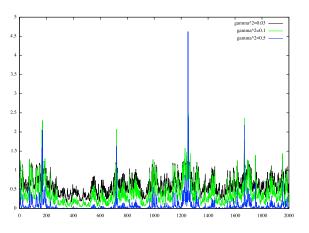
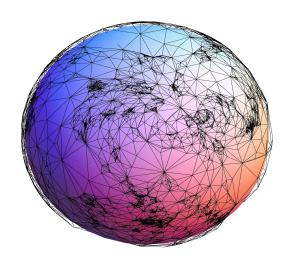


Figure: Density of Gaussian multiplicative chaos

# Uniformisation of a uniform triangulation: courtesy of N. Curien



## Liouville quantum gravity

It is conjectured to be the limit of random planar maps weighted by a critical statistical physics system (CFT with central charge  $c \leq 1$ ) and conformally mapped to a domain D:

- Ambjorn-Durhuus-Jonsson (2005): Quantum geometry: A Statistical Field Theory Approach.
- Sheffield (2010): Conformal weldings of random surfaces: SLE and the quantum gravity zipper. Precise math conjectures.
- Curien (2013): A glimpse of the conformal structure of random planar maps (c=0). First step in a mathematical proof.
- Miller, Sheffield (2014): Quantum Loewner evolution.

## Critical Gaussian multiplicative chaos: $\gamma = 2, \beta = 0$

#### Theorem (Duplantier, Rhodes, Sheffield, V, 2012)

There exists a random measure M such that following limit exists almost surely in the space of Radon measures:

$$\varepsilon^2(2\ln\frac{1}{\varepsilon}-X_{\varepsilon}(x))M_{\varepsilon}^{2,0}(dx)\underset{\varepsilon\to 0}{\to}M'(dx).$$

 $M^{'}$  is called critical Gaussian multiplicative chaos associated to the Green kernel.

## Critical Gaussian multiplicative chaos: $\gamma = 2, \beta = 0$

#### Theorem (Duplantier, Rhodes, Sheffield, V, 2012)

The following limit exists almost surely (along suitable subsequences) in the space of Radon measures:

$$\sqrt{\ln\frac{1}{\varepsilon}}\varepsilon^2 M_\varepsilon^{2,0}(dx)\underset{\varepsilon\to 0}{\to} \sqrt{\frac{2}{\pi}}M'(dx).$$

#### Theorem (Barral, Kupiainen, Nikula, Saksman, Webb, 2013)

The measure M' lives on a set of Hausdorff dimension 0.

## Complex Gaussian multiplicative chaos: Phase diagram

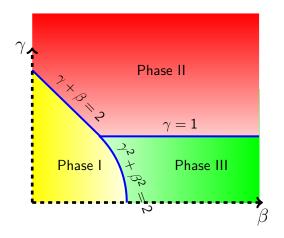


Figure: Phase diagram

## Previous works on the topic

#### Previous work on the complex case:

 Computation of the Free Energy of complex multiplicative cascades: Derrida, Evans, Speer, 1993. In our context:

$$\lim_{\varepsilon \to 0} \frac{\ln |M_\varepsilon^{\gamma,\beta}([0,1]^2)|}{\ln \frac{1}{\varepsilon}}$$

 Complex multiplicative cascades: series of works in dimension 1 by Barral, Jin, Mandelbrot, 2010. Essentially investigated phase I. Partial results in phase III.

# Convergence in phase I and it's frontier I/II (excluding extremal points)

#### Theorem (Lacoin, Rhodes, V., 2013)

On phase I and it's frontier I/II (excluding the extremal points), the  $\mathcal{D}'(D)$ -valued distribution:

$$M_{\varepsilon}^{\gamma,\beta}: \varphi o arepsilon^{rac{\gamma^2}{2}-rac{eta^2}{2}} \int_D arphi(x) M_{arepsilon}^{\gamma,eta}(dx)$$

converges almost surely in the space  $\mathcal{D}_2'(D)$  of distributions of order 2 towards a non trivial limit  $M^{\gamma,\beta}$ .

#### The Sine Gordon model

The probability measure is  $e^{-S(Y)}dY$  where S(Y) is the action:

$$S(Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^{2} dx + \mu \int_{D} \cos(\beta Y(x)) dx$$

Different regimes ( $\varepsilon \to 0$ ):

- $\beta^2 < 2$ : non trivial convergence of  $\mathbb{E}[e^{-\mu \varepsilon^{-\beta^2/2} \int_D \cos(\beta Y_\varepsilon(x)) dx}]$
- $2 \le \beta^2 < \beta_c^2$ :  $\varepsilon^{-\beta^2/2} \int_D \cos(\beta Y_\varepsilon(x)) dx \approx \sigma_\varepsilon N + O(1)$  where  $\sigma_\varepsilon \to \infty$  and N Gaussian variable.
- $\beta^2 > \beta_c^2$ : more and more cumulants to take out...

## Convergence in the inner phase III and it's frontier I/III

#### Theorem (Lacoin, Rhodes, V., 2013)

• When  $\gamma \in [0,1[$  and  $\beta^2 + \gamma^2 > 2$ , we have

$$\left(\varepsilon^{\gamma^2-1}M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D}\Rightarrow (W_{\sigma^2M^{2\gamma,0}}(A))_{A\subset D}.$$
 (1)

where  $\sigma^2 := \sigma^2(\beta^2 + \gamma^2) > 0$  and W is a complex Gaussian measure on D with intensity  $\sigma^2 M^{2\gamma,0}$ .

• When  $\gamma \in [0,1[$  and  $\beta^2 + \gamma^2 = 2$ , we have

$$\left(\varepsilon^{\gamma^2-1}|\log\varepsilon|^{-1/2}M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D}\Rightarrow (W_{\sigma^2M^{2\gamma,0}}(A))_{A\subset D}. \quad (2)$$

where  $\sigma^2 > 0$  and W is a complex Gaussian measure on D with intensity  $\sigma^2 M^{2\gamma,0}$ .



## Convergence in the frontier phase II/III

#### Theorem (Lacoin, Rhodes, V., 2013)

When  $\gamma = 1$  and  $\beta^2 + \gamma^2 > 2$ , we have

$$\left(|\ln\varepsilon|^{1/4}M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D}\Rightarrow (W_{\sigma^2M'}(A))_{A\subset D}.$$

with  $\sigma^2 := \sigma^2(\beta) > 0$  and  $W_{\sigma^2M'}(\cdot)$  is a complex Gaussian random measure with intensity  $\sigma^2M'$ .

## Conformal Field theory c=1 coupled to Liouville Quantum Gravity

Recall that, on phase I and it's frontier I/II (excluding the extremal points), the  $\mathcal{D}'(D)$ -valued distribution:

$$M_{arepsilon}^{\gamma,eta}:arphi
ightarrowarepsilon^{rac{\gamma^2}{2}-rac{eta^2}{2}}\int_Darphi(x)M_{arepsilon}^{\gamma,eta}(dx)$$

converges almost surely in the space  $\mathcal{D}_2'(D)$  of distributions of order 2 towards a non trivial limit  $M_{X,Y}^{\gamma,\beta}$ .

## Setup

In fact, we must denote

$$M_{X,Y}^{\gamma,\beta}(dx) = e^{\gamma X(x) + i\beta Y(x) - \frac{\gamma^2}{2} \mathbb{E}[X(x)^2] + \frac{\beta^2}{2} \mathbb{E}[Y(x)^2]} C(x,D)^{\frac{\gamma^2}{2} - \frac{\beta^2}{2}} dx,$$

where C(x, D) is the conformal radius. This is because we do not renormalize by the mean!

## CFT with central charge c=1 coupled to Gravity

Polyakov action on a domain D

$$S(X,Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^{2} dx + \frac{1}{4\pi} \int_{D} |\nabla X(x)|^{2} + QR(x)X(x) dx,$$

R is the curvature and Q=2

• Equivalence class of random surfaces:

$$(X, Y) \rightarrow (X \circ \psi + 2 \ln |\psi'|, Y \circ \psi),$$

where  $\psi: \tilde{D} \to D$  is a conformal map. See Ginsparg, Moore (1993), Lectures on 2D gravity and 2D string theory.



## The Tachyon fields

Under the above equivalence class  $(\psi: ilde{D} o D)$ 

$$M_{X\circ\psi+2\ln|\psi'|,Y\circ\psi}^{\gamma,\beta}(\varphi)=|\psi'\circ\psi^{-1}|^{2\gamma-\frac{\gamma^2}{2}+\frac{\beta^2}{2}-2}M_{X,Y}^{\gamma,\beta}(\varphi\circ\psi^{-1}),$$

for every function  $arphi \in \mathcal{C}^2_c( ilde{D})$ 

Tachyon Fields are conformally invariant. One must solve

$$2\gamma - \frac{\gamma^2}{2} + \frac{\beta^2}{2} - 2 = 0 \leftrightarrow \gamma \pm \beta = 2, \ \gamma \in ]1, 2[.$$

## The Tachyon field for $(\gamma = 2, \beta = 0)$

At the special point ( $\gamma=2,\beta=0$ ), we recover the special tachyon field, i.e. the background measure

$$M_{X,Y}^{\gamma,\beta}(A) = M'(A)$$

where M' is critical Gaussian multiplicative chaos.

#### Sine-Gordon model

The probability measure is  $e^{-S(Y)}dY$  where S(Y) is the action:

$$S(Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^{2} dx + \mu \int_{D} \cos(\beta Y(x)) dx$$

Representation of the density of charge  $\rho$  of the Coulomb gas:

$$<\rho(x)\rho(y)>=\mathbb{E}[\sin(\beta Y(x))\sin(\beta Y(y))e^{-\mu\int_{D}\cos(\beta Y(z))dz}]$$

## Sine-Gordon model coupled to gravity?

The probability measure is  $e^{-S(X,Y)}dXdY$  where S(X,Y) is the action:

$$S(X,Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^{2} dx + \frac{1}{4\pi} \int_{D} |\nabla X(x)|^{2} dx + \mu_{1} \int_{D} \cos(\beta Y(x)) e^{\gamma X(x)} dx + \mu_{2} \int_{D} e^{2X(x)} dx$$

where  $\gamma + \beta = 2$  (see G. Moore, Gravitational Phase transitions and the Sine-Gordon model).

The problem is linked to defining the Coulomb gas on a random lattice.