# Metastability in the reversible inclusion process II

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joint work with Sander Dommers & Cristian Giardinà



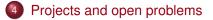
Warwick, 4-8 January 2016

# Outline









Interacting particles system with *N* particles moving on a (finite) set *S* following a given Markovian dynamics.

[Giardinà, Kurchan, Redig, Vafayi (2009); Giardinà, Redig, Vafayi (2010)]

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[Giardinà, Kurchan, Redig, Vafayi (2009); Giardinà, Redig, Vafayi (2010)]

• Configurations:  $\eta \in \{0, 1, 2, ...\}^{S}$   $\eta = (\eta_x)_{x \in S}$ 

with 
$$\eta_x = \#$$
 particles on  $x$  s.t.  $\sum_{x \in S} \eta_x = N$ 

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• Markovian dynamics:

$$\mathcal{L}f(\eta) = \sum_{x,y\in\mathcal{S}} r(x,y)\eta_x(d_N + \eta_y) \left(f(\eta^{x,y}) - f(\eta)\right) \quad \text{generator}$$

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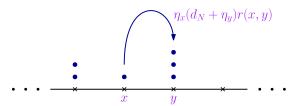
$$\mathcal{L}f(\eta) = \sum_{x,y \in S} r(x,y) \eta_x (d_N + \eta_y) \left( f(\eta^{x,y}) - f(\eta) \right) \quad \text{generator}$$

•  $r(x, y) \ge 0$  transition rates of a *m*-reversible RW on *S* 

•  $d_N > 0$  constant tuning the rates of the underlying RW

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Example:



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Example:  $\eta_x(d_N + \eta_y)r(x, y)$ 

#### Remarkable facts:

Under suitable hypotheses, e.g. letting  $N \to \infty$  and  $d_N \longrightarrow 0$ , the model displays condensation (*particles concentrate on single site*) and metastable behavior (*condensate may appears in different sites of S*)

[Grosskinsky, Redig, Vafayi 2011], [Chleboun 2012].

## Metastable behavior

- Symmetric IP [Grosskinsky, Redig, Vafayi 2013] At the timescale 1/*d*<sub>N</sub>:
  - formation of the condensate
  - condensate jumps between  $x, y \in S$  at rate r(x, y)

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 $S^* = \arg \max\{m(x) : x \in S\}$ m(x) rev. measure for r s.t.  $\max\{m(x) : x \in S\} = 1$ 

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- The RW restricted to  $S^*$  need not to be irreducible
  - $\implies$  the condensate may be trapped in subsets of  $S^*$

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[Chleboun, Grosskinsky 2014], [Cao, Chleboun, Grosskinsky 2014], [Evans, Waklaw (2014)]

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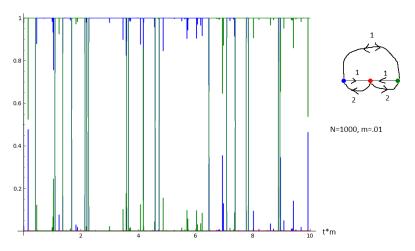
[Chleboun, Grosskinsky 2014], [Cao, Chleboun, Grosskinsky 2014], [Evans, Waklaw (2014)]

Main goal: Characterization of *further metastable timescales*, and *motion of the condensate between traps*.

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# Simulations and heuristics

### **First example**

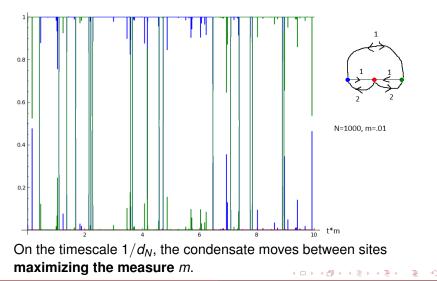


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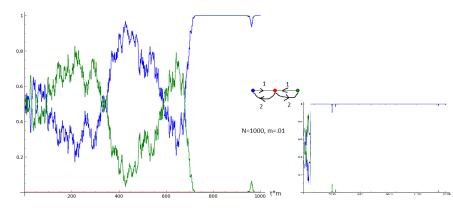
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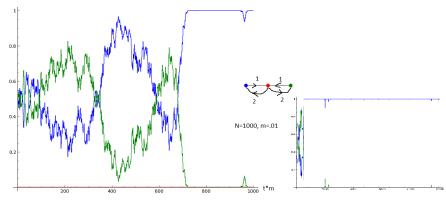
## Second example



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# Simulations and heuristics

## Second example



On the timescale  $1/d_N$ , condensation takes place (though at a long scaled time), while once created, the condensate remains trapped for very long time on a vertex of  $S^*$ .

Assume  $\{r(x, y)\}_{x, y \in S^*}$  is reducible, and let  $C_1, \ldots, C_m, m \ge 2$ , the connected components of  $(S^*, r_{|_{S^*}})$ 

$$S^* = \bigcup_{j=1}^m C_j \ , \ C_i \cup C_j = \emptyset, \text{ for } i \neq j$$

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As for the derivation of the first metastable timescale,  $1/d_N$ , we apply the martingale approach to metastability.

[Beltrán, Landim, 2010-2015]

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• Define a new set of metastable sets  $\mathcal{E}_1, \ldots, \mathcal{E}_m$ :

$$\mathcal{E}_{j} = \bigcup_{x \in \mathcal{C}_{j}} \eta^{N,x}$$
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• Verify the hypotheses  $H_0$ ,  $H_1$  and  $H_2$  of [Beltrán, Landim, 2010]  $\longrightarrow$  compute capacities  $\operatorname{Cap}_N(\mathcal{E}_i, \mathcal{E}_j)$ .

# Capacity versus Metastability

#### Capacity is a key quantity in the analysis of metastable systems

[Bovier, Eckhoff, Gayrard, Klein '01-'04]-[Beltrán, Landim '10-'15]

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Its definition comes from correspondence btw reversible dynamics and electrical networks. If  $A, B \subset \Omega$ , and  $\mu$  reversible measure

$$\mathsf{Cap}(\mathit{A}, \mathit{B}) := \sum_{\eta \in \mathit{A}} \mu(\eta) \mathbb{P}_{\eta}[ au_{\mathit{B}} < au_{\mathit{A}}^+]$$

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# **I Fact.** If *A* e *B* are metastable sets, the mean metastable time btw *A* e *B* is (roughly) $\sim \mu(A)/\text{Cap}(A,B)$ .

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 $\longrightarrow$  we start from a simple IP dynamics to understand traps.

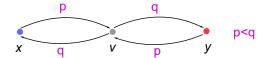
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## Analysis of a 3- sites IP

Consider the IP defined through the underlying RW on  $S = \{v, x, y\}$  with transition rates s.t.

$$\begin{cases} r(y,x) = r(x,y) = 0\\ m(x) = m(y) = 1 > m(v) \end{cases}$$

 $\implies \eta^{N,x}, \eta^{N,y}$  are disconnected components of  $(S^*, r_{|_{S^*}})$ 



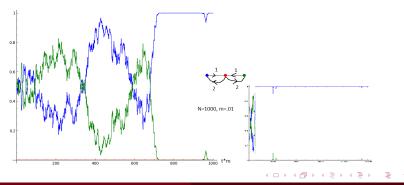
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## Capacities for the 3-sites IP

## **Proposition 1.**

In the above notation and for  $d_N \log N \rightarrow 0$ ,

$$\lim_{N\to\infty}\frac{N}{d_N^2}\cdot\operatorname{Cap}_N(\eta^{N,x},\eta^{N,y}) = \left(\frac{1}{r(v,x)} + \frac{1}{r(v,y)}\right)^{-1}\cdot\frac{m(v)}{1-m(v)}$$

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#### In particular

 $\operatorname{Cap}_{N}(\eta^{N,x},\eta^{N,y}) \sim d_{N}^{2}/N \ll d_{N} \longrightarrow \text{ second timescale } \sim N/d_{N}^{2}$ 

Projects and open problems

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Following [Beltrán, Landim 2010], hypothesis H<sub>0</sub> is verified:

$$\lim_{N \to \infty} \frac{N}{d_N^2} p_N(\eta^{N,x}, \eta^{N,y}) = \left(\frac{1}{r(v,x)} + \frac{1}{r(v,y)}\right)^{-1} \frac{m(v)}{1 - m(v)} =: p^{(2)}(x,y)$$

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## Dynamics of the condensate in the 3-sites IP

As a consequence (hypotheses  $H_1$  and  $H_2$  are easily verified), for

$$X_{\mathsf{N}}(t) = \sum_{z \in S^*} z \mathbb{1}_{\{\eta_z(t) = \mathsf{N}\}}$$

#### **Proposition 2.**

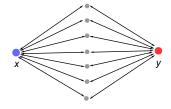
Let  $d_N \log N \to 0$  as  $N \to \infty$ , and  $\eta_z(0) = N$  for some  $z \in S^*$ . Then

 $X_N(t \cdot N/d_N^2)$  converges weakly to x(t) as  $N \to \infty$ 

where x(t) is a Markov process on  $S^*$  with symmetric rates  $p^{(2)}(x, y)$ .

# Easy extension I

## Consider the IP defined through the following underlying RW

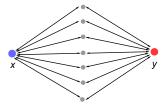


with transition rates s.t.

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In the above notation and for  $d_N \log N \rightarrow 0$ ,

$$\lim_{N \to \infty} \frac{N}{d_N^2} \cdot \operatorname{Cap}_N(\eta^{N,x}, \eta^{N,y}) = \sum_{\nu \neq \{x,y\}} \left( \frac{1}{r(\nu,x)} + \frac{1}{r(\nu,y)} \right)^{-1} \frac{m(\nu)}{1 - m(\nu)}$$

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For  $S^* = \bigcup_{j \in J} C_j$ ,  $C_j$  connected component of  $(S^*, R_{|_{S^*}})$ . For  $I \subsetneq J$  $S_1^* = \bigcup_{i \in I} C_i$ ,  $S_2^* = \bigcup_{i \in J \setminus I} C_i$  and  $\mathcal{E}_1 = \bigcup_{x \in S_1^*} \eta^{N,x}$ ,  $\mathcal{E}_2 = \bigcup_{y \in S_2^*} \eta^{N,y}$ 

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Then, for  $d_N \log N \rightarrow 0$ ,

$$\lim_{N \to \infty} \frac{N}{d_N^2} \cdot \operatorname{Cap}_N(\mathcal{E}_1, \mathcal{E}_2) = \sum_{\substack{x \in S_1^* \\ y \in S_2^*}} \sum_{v \in S \setminus S^*} \left( \frac{1}{r(v, x)} + \frac{1}{r(v, y)} \right)^{-1} \frac{m(v)}{1 - m(v)}$$
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# Analysis of a IP on on $\{1, 2, \ldots, L\}$ , $L \ge 4$

Let  $S = \{x = v_1, v_2, ..., v_L = y\}$  with  $L \ge 4$  and consider the IP defined through the following RW



with transition rates s.t.  $S^* = \{x, y\}$ 

# Capacities for the IP on $\{1, 2, \ldots, L\}$

### **Proposition 3.**

In the above notation and for  $d_N \log N \rightarrow 0$ ,

$$\lim_{N \to \infty} \frac{N^2}{d_N^{L-1}} \cdot \operatorname{Cap}_N(\eta^{N,x},\eta^{N,y}) \geq c > 0$$

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Upper bound on these capacities is still under investigation, but rough (not-matching) estimates show

 $\operatorname{Cap}_{N}(\eta^{N,x},\eta^{N,y}) \leq d_{N}^{3}/N^{2} \ll d_{N}^{2}/N \longrightarrow \text{third timescale}$ 

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Are there further metastable timescales associated to length of minimal path btw connected components? (to investigate)

## Comment on the results

Main goal of the work (*in progress*)  $\longrightarrow$  characterization of the condensate dynamics for the IP with reversible rates.

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Though results rigorously obtained only for simple underlying RW (1D RW), we expect that the mechanism highlighted here hold in generality.

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Though results rigorously obtained only for simple underlying RW (1D RW), we expect that the mechanism highlighted here hold in generality.

We conjecture the existence of many (at least three) metastable timescales:

Consider the IP on a finite set *S* with reversible rates r(x, y). Let  $\left\{C_1^{(k)}, \ldots, C_{m(k)}^{(k)}\right\}$ , for k = 1, 2, be partitions of *S*<sup>\*</sup>, of cardinality  $0 \le m(k) \le |S^*|$ , such that •  $C_i^{(1)}$ 's are the connected components of  $(S^*, r_{ls*})$ 

•  $C_j^{(2)}$ , are s.t  $x \in S^*$  if  $d(x, S^* \setminus \{x\}) \le 2$ (*d* graph distance on (S, r)).

Let 
$$T_1 = N/d_N^2$$
 and  $T_2 = N^2/d_N^3$ , and

$$\mathcal{E}_j^{(k)} = \bigcup_{x \in \mathcal{C}_j^{(k)}} \eta^{N,i}, \quad \forall j, \ k = 1, 2$$

(metastable sets at timescale  $T_k$ )

$$X_N^{(k)}(t) = \sum_{j=1}^{m^{(k)}} j \mathbb{1}_{\{\eta(t) \in \mathcal{E}_j^{(k)}\}}, \ k = 1, 2$$

(processes projected on metastable sets)

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Following the techniques of [Beltrán, Landim, 2010-2014],

### Conjecture 1.

Let  $d_N \log N \to 0$  as  $N \to \infty$ , and  $\eta_z(0) = N$  for some  $z \in S^*$ . Then, for k = 1, 2

 $X_N^{(k)}(t \cdot T_k)$  converges weakly to  $x^{(k)}(t)$  as  $N \to \infty$ 

where  $x^{(k)}(t)$  is a Markov process on  $\{1, ..., m(k)\}$  with  $x^{(k)}(0) = \sum_{j=1}^{m(k)} j \mathbb{1}_{\{\eta(0) \in \mathcal{E}_{j}^{(k)}\}}$ .

Condensation in the IP

Metastable timescales

Analysis of IP on 1D lattice

Projects and open problems

# Upper and lower bound on capacities

Recall that by the Dirichlet principle

$$Cap_N(A, B) = \inf_{f:f_{|_A}=1, f_{|_B}=0} \{D_N(f)\}$$

where the Dirichlet form of the IP is

$$D_N(f) = \frac{1}{2} \sum_{\eta} \mu_N(\eta) \sum_{x,y \in S} R(\eta, \eta^{x,y}) \left( f(\eta^{x,y}) - f(\eta) \right)^2$$

with 
$$\mu_N(\eta) = \frac{1}{Z_N} \prod_{x \in S} m(x)^{\eta_x} w_N(\eta_x)$$
$$w_N(k) = \frac{\Gamma(k + d_N)}{k! \Gamma(d_N)}$$
$$R(\eta, \eta^{x, y}) = r(x, y) \eta_x(\eta_y + d_N)$$

(see computation at the blackboard)

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# Future projects

## Formation of the condensate in the TASIP

(joint with S. Dommers, S. Grosskinsky)

Consider the IP on  $S = \mathbb{Z}/L\mathbb{Z}$  in the totally asymmetric case, with  $r(x, x - 1) = 0 \ \forall x \in S$  (jumps to the right).

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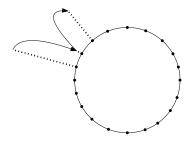
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Main goal: Provide rigourous arguments to compute the condensation time.

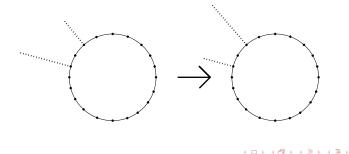
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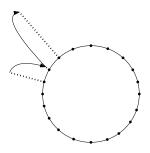


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By simulations, the slower dynamical step along nucleation is the union of two half-condensates, of size  $m_1 > m_2$ . This can be represented by a cyclic mechanism of 4 main steps: **II step** *The largest condensate loses*  $O(m_1 - m_2)$  *particles before they are absorbed by the other condensate. At the hand, the two condensates had roughly exchanged mass.* 



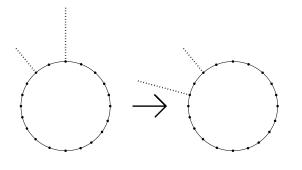
By simulations, the slower dynamical step along nucleation is the union of two half-condensates, of size  $m_1 > m_2$ . This can be represented by a cyclic mechanism of 4 main steps: **III step** *As in the first step, the smallest condensate (now on the left) loses one particle which is absorbed by the other condensate. Difference*  $m_2 - m_1$  *increases.* 



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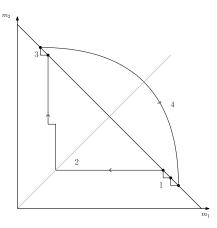
By simulations, the slower dynamical step along nucleation is the union of two half-condensates, of size  $m_1 > m_2$ . This can be represented by a cyclic mechanism of 4 main steps: **IV step** *The largest condensate (on the right) moves faster until* 

reaches the configuration at step I.



#### Reduction to a random walk

 $\begin{cases} m_1 = \# \text{particles of condensate on the left} \\ m_2 = \# \text{particles of condensate on the right} \\ m_1 + m_2 = N. \end{cases} \text{ with }$ 



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We studied the reversible IP on a finite set in the limit  $N \to \infty$  and for  $d_N \to 0$  with  $d_N \log N \to 0$  by martingale approach:

- We derive the dynamics of the condensate at timescale  $\sim 1/d_N$ ;
- We prove the existence of a longer metastable timescale
   N/d<sup>2</sup><sub>N</sub> and derive dynamics of the condensate in simple IP processes (1D RW);
- We conjecture the existence of longer metastable timescales (at least one) ~ N<sup>2</sup>/d<sup>a</sup><sub>N</sub> with a = a(r) ≥ 3 and possibly dependent on the shortest length of paths between components of (S\*, r<sub>|s\*</sub>).

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• Conjecture on metastable timescales longer than  $N/d_N^2$  and derivation of asymptotic rates

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- Thermodynamic limit  $|S| \rightarrow \infty$  with  $N/|S| \rightarrow \rho > 0$

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# Thank you for your attention!

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