# **Universal Fluctuation Formulae** for one-cut β-ensembles with a combinatorial touch

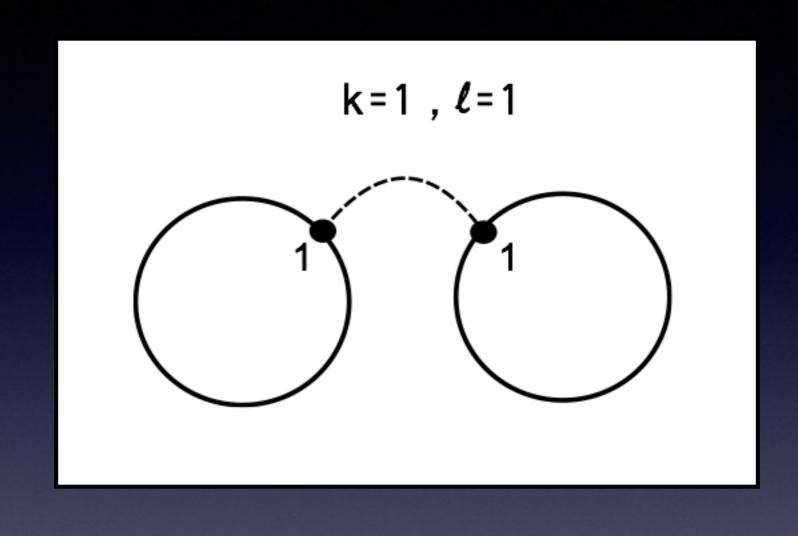
Pierpaolo Vivo

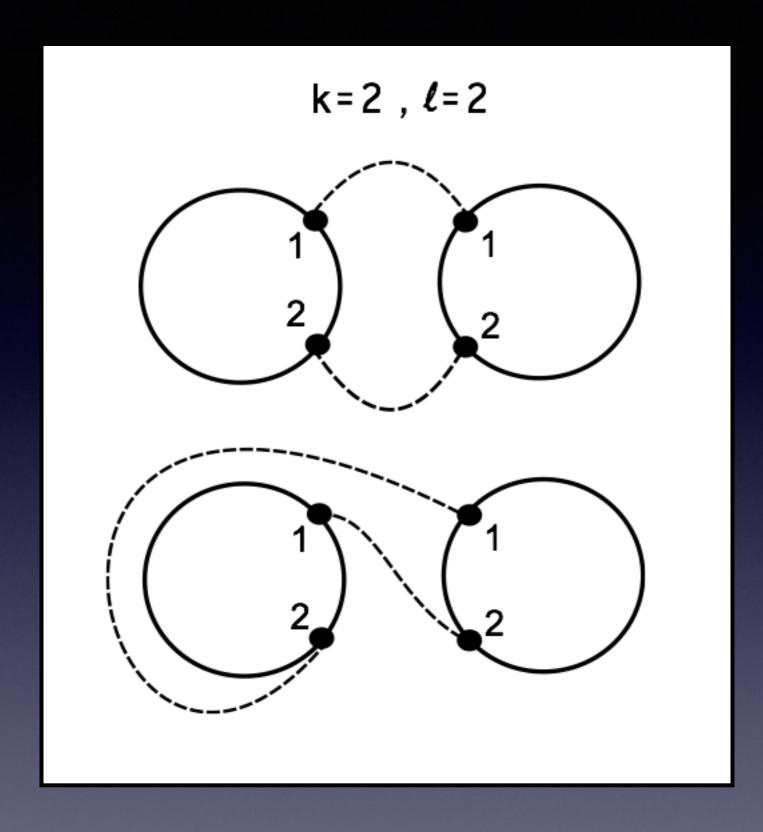


with F. D. Cunden

Phys. Rev. Lett. 113, 070202 (2014)

with F.D. Cunden and F. Mezzadri J. Phys. A 48, 315204 (2015)





# Non-crossing pairings of two circles

```
\begin{pmatrix} 1 & 0 & 3 & 0 & 10 & 0 & 35 & 0 \\ 0 & 2 & 0 & 8 & 0 & 30 & 0 & 112 \\ 3 & 0 & 12 & 0 & 45 & 0 & 168 & 0 \\ 0 & 8 & 0 & 36 & 0 & 144 & 0 & 560 \\ 10 & 0 & 45 & 0 & 180 & 0 & 700 & 0 \\ 0 & 30 & 0 & 144 & 0 & 600 & 0 & 2400 \\ 35 & 0 & 168 & 0 & 700 & 0 & 2800 & 0 \\ 0 & 112 & 0 & 560 & 0 & 2400 & 0 & 9800 \end{pmatrix}
```

No pairings are possible if 'k' and 'l' have different parity

• Formula for these numbers?

• Connection to random matrices?

## • Formula for these numbers?

#### A NEW BRANCH OF ENUMERATIVE GRAPH THEORY

BY W. T. TUTTE Communicated by Walter Rudin, April 6, 1962



W. T. TUTTE

$$\begin{cases} \frac{2\kappa\ell}{\kappa+\ell} \binom{\kappa-1}{\left\lfloor\frac{\kappa}{2}\right\rfloor} \binom{\ell-1}{\left\lfloor\frac{\ell}{2}\right\rfloor} & \text{if } \kappa=\ell \mod 2 \\ 0 & \text{if } \kappa\neq\ell \mod 2, \end{cases}$$

### • Formula for these numbers?

#### A NEW BRANCH OF ENUMERATIVE GRAPH THEORY

```
In[30] := T = Table \Big[ If \Big[ (EvenQ[k] && EvenQ[1]) \mid | (OddQ[k] && OddQ[1]), \\ \frac{2 k 1}{k+1} Binomial[k-1, Floor[k/2]] Binomial[1-1, Floor[1/2]], 0 \Big], \{k, 1, 8\}, \{1, 1, 8\} \Big]; \\ T // MatrixForm \\ \hline \\ Out[31] // MatrixForm \\ \begin{bmatrix} 1 & 0 & 3 & 0 & 10 & 0 & 35 & 0 \\ 0 & 2 & 0 & 8 & 0 & 30 & 0 & 112 \\ 3 & 0 & 12 & 0 & 45 & 0 & 168 & 0 \\ 0 & 8 & 0 & 36 & 0 & 144 & 0 & 560 \\ 10 & 0 & 45 & 0 & 180 & 0 & 700 & 0 \\ 0 & 30 & 0 & 144 & 0 & 600 & 0 & 2400 \\ 35 & 0 & 168 & 0 & 700 & 0 & 2800 & 0 \\ 0 & 112 & 0 & 560 & 0 & 2400 & 0 & 9800 \\ \end{bmatrix}
```

$$\begin{cases} \frac{2\kappa\ell}{\kappa+\ell} \binom{\kappa-1}{\left\lfloor \frac{\kappa}{2} \right\rfloor} \binom{\ell-1}{\left\lfloor \frac{\ell}{2} \right\rfloor} & \text{if } \kappa=\ell \mod 2 \\ 0 & \text{if } \kappa\neq\ell \mod 2, \end{cases}$$

Connection to random matrices?

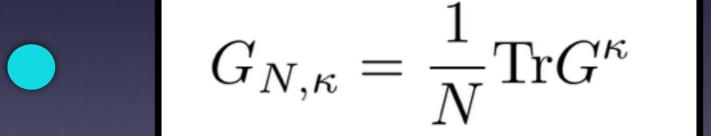
$$G_{ij} = \frac{1}{\sqrt{N}} x_{ij}$$
  $x_{ij} \sim \mathcal{N}_{\mathbb{C}}(0,1)$  for  $i < j$  and  $x_{ii} \sim \mathcal{N}_{\mathbb{R}}(0,1)$ 

GUE

Connection to random matrices?

$$G_{ij} = \frac{1}{\sqrt{N}} x_{ij}$$
  $x_{ij} \sim \mathcal{N}_{\mathbb{C}}(0,1)$  for  $i < j$  and  $x_{ii} \sim \mathcal{N}_{\mathbb{R}}(0,1)$ 

### GUE



Connection to random matrices?

$$G_{ij} = \frac{1}{\sqrt{N}} x_{ij}$$
  $x_{ij} \sim \mathcal{N}_{\mathbb{C}}(0,1)$  for  $i < j$  and  $x_{ii} \sim \mathcal{N}_{\mathbb{R}}(0,1)$ 

#### **GUE**

$$G_{N,\kappa} = \frac{1}{N} \text{Tr} G^{\kappa}$$

$$\alpha_{\kappa,\ell}^{\mathcal{G}} = \lim_{N \to \infty} \beta N^2 \operatorname{Cov} (G_{N,\kappa}, G_{N,\ell})$$

#### • Connection to random matrices?

(0,1)

 $(1/2)\alpha_{\kappa,\ell}^{\mathfrak{G}}=\#\{\text{Non-crossing pairings of two circles with }\kappa \text{ points on the 1st circle}$  and  $\ell$  points on the 2nd circle such that the two circles are connected $\}$ .

$$\alpha_{\kappa,\ell}^{\mathcal{G}} = \lim_{N \to \infty} \beta N^2 \operatorname{Cov} (G_{N,\kappa}, G_{N,\ell})$$

• Communities interested in this?

Combinatorics Free Probability **RMT** 

## Communities interested in this?

#### A NEW BRANCH OF ENUMERATIVE GRAPH THEORY

BY W. T. TUTTE Communicated by Walter Rudin, April 6, 1962

Real Second-Order Freeness and the Asymptotic Real Second-Order Freeness of Several Real Matrix Models

Catherine Emily Iska Redelmeier

## Free Probability

Combinatorics

JOURNAL OF MATHEMATICAL PHYSICS 47, 063302 (2006)

Global spectrum fluctuations for the  $\beta$ -Hermite and  $\beta$ -Laguerre ensembles via matrix models

Ioana Dumitriu<sup>a)</sup> and Alan Edelman University of California, Berkeley, Berkeley, California 94720

**RMT** 

GLOBAL FLUCTUATIONS FOR LINEAR STATISTICS OF  $\beta$ -JACOBI ENSEMBLES

IOANA DUMITRIU<sup>∗</sup> and ELLIOT PAQUETTE<sup>†</sup>

## Communities interested in this?

#### A NEW BRANCH OF ENUMERATIVE GRAPH THEORY

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Real Second-Order Freeness and the Asymptotic Real Second-Order Freeness of Several Real Matrix Models

### Combinatorics

Catherine Emily Iska Redelmeier

satisfy this definition. Second-order freeness then has the role for fluctua-Free tions that first-order freeness has for moments. In particular, if random matrices A and B are asymptotically second-order free, the asymptotic fluctuations of A + B and AB can be calculated from the asymptotic moments and fluctuations of A and B.

> Asymptotic second-order freeness as defined in [18] is not generally satis field by real ensembles of random matrices. If random matrices  $A_{k,N}$  and  $B_{l,N}, k, l = 1, \ldots, p$  are elements of the algebra generated by a model studied in this paper, then the relation satisfied instead is

$$\lim_{N \to \infty} \operatorname{cov} \left( \operatorname{Tr} \left( \mathring{A}_{1,N} \cdots \mathring{A}_{p,N} \right), \operatorname{Tr} \left( \mathring{B}_{1,N} \cdots \mathring{B}_{p,N} \right) \right)$$

$$= \sum_{k=0}^{p-1} \prod_{i=1}^{p} \lim_{N \to \infty} \mathbb{E} \left( \operatorname{tr} \left( \mathring{A}_{i,N} \mathring{B}_{k-i,N} \right) \right) + \sum_{k=0}^{p-1} \prod_{i=0}^{p} \lim_{N \to \infty} \mathbb{E} \left( \operatorname{tr} \left( \mathring{A}_{i,N} \mathring{B}_{k+i,N}^T \right) \right)$$

## General problem

$$d\mathbb{P}_{N,\beta}(\{\lambda_k\}) = \frac{1}{\mathcal{Z}_{N,\beta}} e^{-\beta\left(-\frac{1}{2}\sum_{i\neq j}\ln|\lambda_i - \lambda_j| + N\sum_i V(\lambda_i)\right)} \prod_{i=1}^N d\lambda_i.$$

Can be interpreted as the joint probability density of eigenvalues of a  $\beta$ -ensemble of random matrices

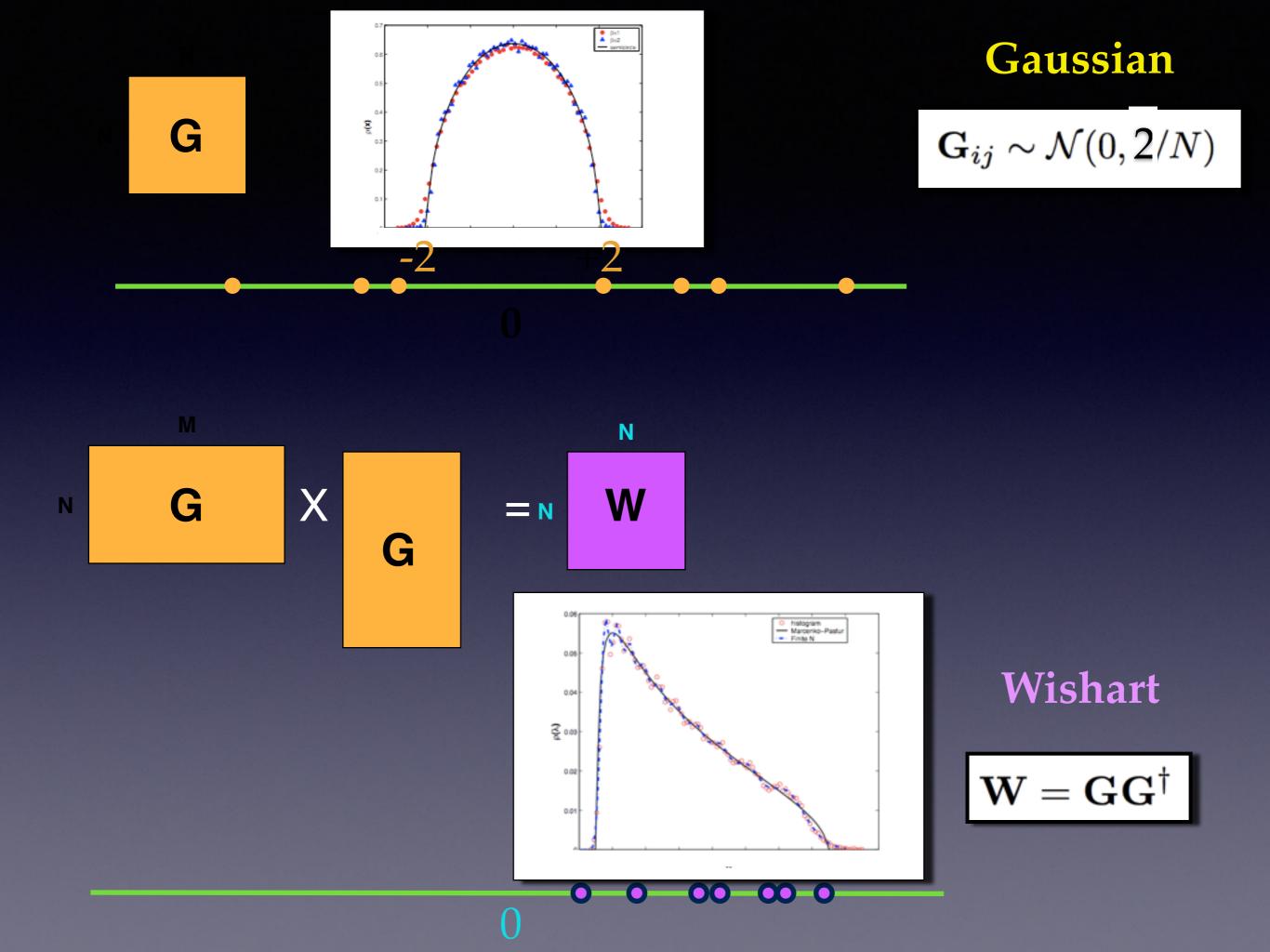
# General setting

$$d\mathbb{P}_{N,\beta}(\{\lambda_k\}) = \frac{1}{\mathcal{Z}_{N,\beta}} e^{-\beta \left(-\frac{1}{2}\sum_{i\neq j}\ln|\lambda_i - \lambda_j| + N\sum_i V(\lambda_i)\right)} \prod_{i=1}^N d\lambda_i.$$

"Inverse Temperature" "External potential"

We further assume a "one-cut" property: the external potential is such that the spectral density in the large-N limit is supported on a single interval of the real line

Examples: β-Gaussian β-Wishart β-Jacobi



$$X_{N,\kappa} = N^{-1} \operatorname{Tr} \mathfrak{X}_N^{\kappa}$$

$$\lim_{N \to \infty} N^2 \operatorname{Cov}(X_{N,\kappa}, X_{N,\ell}) = \frac{1}{\beta} \alpha_{\kappa,\ell}.$$

$$X_{N,\kappa} = N^{-1} \operatorname{Tr} \mathfrak{X}_N^{\kappa}$$

$$\lim_{N \to \infty} N^2 \operatorname{Cov}(X_{N,\kappa}, X_{N,\ell}) = \frac{1}{\beta} \alpha_{\kappa,\ell}.$$

## Our main result:

$$F(z,\zeta) = \sum_{\kappa,\ell=0}^{\infty} \alpha_{\kappa,\ell} z^{\kappa} \zeta^{\ell} = \frac{z\zeta}{(z-\zeta)^2} \left[ \frac{2\lambda_{-}\lambda_{+}z\zeta - (\lambda_{-} + \lambda_{+})(z+\zeta) + 2}{2\sqrt{(1-\lambda_{-}z)(1-\lambda_{+}z)(1-\lambda_{-}\zeta)(1-\lambda_{+}\zeta)}} - 1 \right]$$

Generating function of the covariances of any one-cut  $\beta$ -ensemble

universal!



# Global spectrum fluctuations for the $\beta$ -Hermite and $\beta$ -Laguerre ensembles via matrix models

Ioana Dumitriu and Alan Edelman

Claim 3.15.2: For any fixed integers k and l,

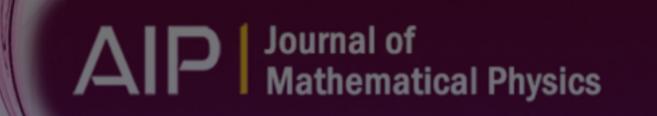
$$\lim_{n\to\infty} \operatorname{Cov}(n_{i,\gamma}(n), \eta_{j,\gamma}(n)) = \frac{2}{\beta} (\operatorname{Sum}_1(i,j,\gamma) + \operatorname{Sum}_2(i,j,\gamma)),$$

where

$$\mathrm{Sum}_{1}(i,j,\gamma) = \sum_{q=1}^{i+j-1} (-1)^{q+1} \gamma^{i+j-q} \frac{\binom{i+j}{q}}{i+j} \sum_{\substack{j=q+1\\ j-1}}^{i+j} \frac{(-1)^{j}}{\binom{i+j-1}{j-1}} \sum_{\substack{r+s=j\\ 1 \leq r \leq i\\ 1 \leq s \leq j}} rs \binom{i}{r}^{2} \binom{j}{s}^{2},$$

$$\operatorname{Sum}_{2}(i,j,\gamma) = \sum_{q=0}^{i+j-2} (-1)^{q} \gamma^{j+j-q} \frac{\binom{i+j}{q}}{i+j} \sum_{j=q}^{i+j-2} \frac{(-1)^{j}}{\binom{i+j-1}{j}} \sum_{\substack{r+s=j\\0 \leqslant r \leqslant i-1\\0 \leqslant s \leqslant j-1}} (i-r)(j-s) \binom{i}{r}^{2} \binom{j}{s}^{2}.$$

$$\alpha_{\kappa,\ell}^{W} = \lim_{N \to \infty} \beta N^{2} \operatorname{Cov}\left(W_{N,\kappa}, W_{N,\ell}\right) = 4(1+c)^{\kappa+\ell} \sum_{\substack{0 \le p \le \kappa \\ 0 \le q \le \ell \\ p = q \bmod 2}} \left(\frac{\sqrt{c}}{1+c}\right)^{p+q} \frac{pq}{p+q} \binom{\kappa}{p} \binom{\ell}{q} \binom{p-1}{\left\lfloor \frac{p}{2} \right\rfloor} \binom{q-1}{\left\lfloor \frac{q}{2} \right\rfloor}.$$



Global spectrum fluctuations for the  $\beta$  -Hermite and  $\beta$  -Laguerre ensembles via matrix models

Ioana Dumitriu and Alan Edelman

Claim 3.15.2: For any fixed integers k and l,

$$\lim_{N \to \infty} N^2 \text{Cov} (W_{N,\kappa}, W_{N,\ell}) = \frac{1}{\beta} \begin{bmatrix} 2c & 4(c+c^2) & 6(c+3c^2+c^3) \\ 4(c+c^2) & 4(2c+5c^2+2c^3) & 12(c+5c^2+5c^3+c^4) \\ 6(c+3c^2+c^3) & 12(c+5c^2+5c^3+c^4) & 6(3c+24c^2+46c^3+24c^4+3c^5) \end{bmatrix}$$

$$\operatorname{Sum}_{1}(i,j,\gamma) = \underbrace{\sum_{q=1}^{n} (-1)^{q+1} \gamma^{+j-q}}_{q=1} \underbrace{\sum_{j=q+1}^{n} \underbrace{\left(i+j-1\right)}_{\substack{j=q+1\\ j-1}} \underbrace{\sum_{\substack{r+s=j\\1 \leq r \leq i\\1 \leq s \leq j}} rs \binom{r}{s} \binom{s}{s}}_{,s},$$

$$\operatorname{Sum}_{2}(i,j,\gamma) = \sum_{q=0}^{i+j-2} (-1)^{q} \gamma^{j+j-q} \frac{\binom{i+j}{q}}{i+j} \sum_{j=q}^{i+j-2} \frac{(-1)^{j}}{\binom{i+j-1}{j}} \sum_{\substack{r+s=j\\0 \leqslant r \leqslant i-1\\0 \leqslant s \leqslant j-1}} (i-r)(j-s) \binom{i}{r}^{2} \binom{j}{s}^{2}.$$

$$\alpha_{\kappa,\ell}^{\mathcal{W}} = \lim_{N \to \infty} \beta N^2 \operatorname{Cov}\left(W_{N,\kappa}, W_{N,\ell}\right) = 4(1+c)^{\kappa+\ell} \sum_{\substack{0 \le p \le \kappa \\ 0 \le q \le \ell \\ p = q \bmod 2}} \left(\frac{\sqrt{c}}{1+c}\right)^{p+q} \frac{pq}{p+q} \binom{\kappa}{p} \binom{\ell}{q} \binom{p-1}{\left\lfloor \frac{p}{2} \right\rfloor} \binom{q-1}{\left\lfloor \frac{q}{2} \right\rfloor}.$$

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## GLOBAL FLUCTUATIONS FOR LINEAR STATISTICS OF $\beta$ -JACOBI ENSEMBLES

IOANA DUMITRIU\* and ELLIOT PAQUETTE†

The matrix  $C_{k,l}$  for  $k,l \geq 1$  can now be defined by

$$C_{k,l} := \frac{\alpha}{4} \int_{-a}^{0} \frac{1}{1+2\sigma} [(\partial_x p_k \partial_x p_m + \partial_y p_k \partial_y p_m)(1-x^2-y^2) - (\partial_x p_k \partial_y p_m + \partial_y p_k \partial_x p_m)(2xy)] d\sigma.$$
(2.14)

$$\alpha_{\kappa,\ell}^{\mathcal{J}} = \left[ \frac{\gamma_1^2 + \gamma_1 \gamma_2 + 2(\gamma_1 + \gamma_2 + 1)}{(\gamma_1 + \gamma_2 + 2)^2} \right]^{\kappa + \ell} \sum_{\substack{0 \le p \le \kappa \\ 0 \le q \le \ell}} \left[ \frac{\sqrt{(\gamma_1 + 1)(\gamma_2 + 1)(\gamma_1 + \gamma_2 + 1)}}{\gamma_1^2 + \gamma_1 \gamma_2 + 2(\gamma_1 + \gamma_2 + 1)} \right]^{p+q} \binom{\kappa}{p} \binom{\ell}{q} \alpha_{p,q}^{\mathcal{G}},$$

$$\lim_{N \to \infty} N^2 \text{Cov} \left( J_{N,\kappa}, J_{N,\ell} \right) = \frac{1}{\beta} \begin{bmatrix} \frac{1}{8} & \frac{1}{8} & \frac{15}{128} & \frac{7}{64} & \frac{105}{1024} \\ \frac{1}{8} & \frac{9}{64} & \frac{9}{64} & \frac{35}{256} & \frac{135}{1024} \\ \frac{15}{128} & \frac{9}{64} & \frac{75}{512} & \frac{75}{512} & \frac{4725}{32768} \\ \frac{7}{64} & \frac{35}{256} & \frac{75}{512} & \frac{1225}{8192} & \frac{1225}{8192} \\ \frac{105}{1024} & \frac{135}{1024} & \frac{4725}{32768} & \frac{1225}{8192} & \frac{19845}{131072} \end{bmatrix}$$

$$\alpha_{\kappa,\ell}^{\mathcal{J}} = \left[ \frac{\gamma_1^2 + \gamma_1 \gamma_2 + 2(\gamma_1 + \gamma_2 + 1)}{(\gamma_1 + \gamma_2 + 2)^2} \right]^{\kappa + \ell} \sum_{\substack{0 \le p \le \kappa \\ 0 \le q \le \ell}} \left[ \frac{\sqrt{(\gamma_1 + 1)(\gamma_2 + 1)(\gamma_1 + \gamma_2 + 1)}}{\gamma_1^2 + \gamma_1 \gamma_2 + 2(\gamma_1 + \gamma_2 + 1)} \right]^{p+q} \binom{\kappa}{p} \binom{\ell}{q} \alpha_{p,q}^{\mathcal{G}},$$

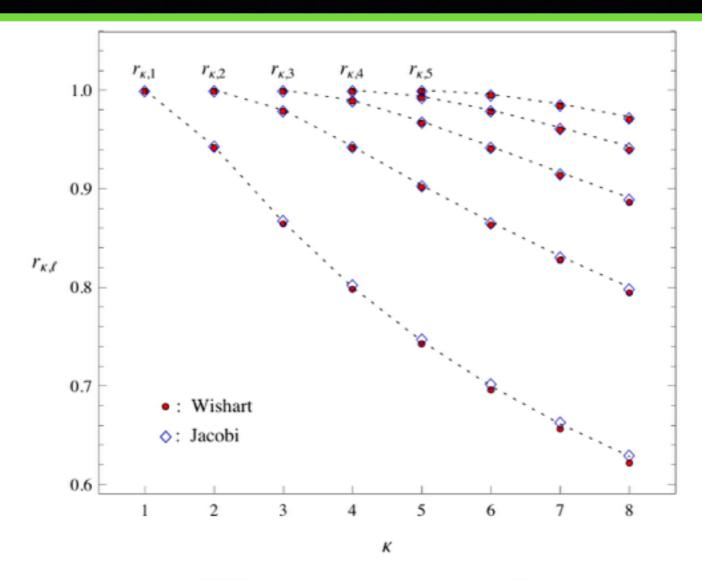


FIG. 2: Limiting correlations  $r_{\kappa,\ell} = \alpha_{\kappa,\ell}/\sqrt{\alpha_{\kappa,\kappa}\alpha_{\ell,\ell}}$  of the moments for the Wishart ensemble  $\mathcal{W}_N$  with c=1 and the Jacobi ensemble  $\mathcal{J}_N$  with  $\gamma_1=\gamma_2=0$ . The densities of states  $\rho_{\mathcal{W}}$  and  $\rho_{\mathcal{J}}$  are supported on [a,b]=[0,4] and [a,b]=[0,1], respectively. According to Remark III, the correlation coefficients  $r_{\kappa,\ell}$  are the same for these ensembles. Here we show a comparison of our prediction (50) with numerical simulations.

More general setting

2 linear statistics, A and B (differentiable)

$$A = \sum_{i=1}^{N} a(\lambda_i)$$

$$B = \sum_{i=1}^{N} b(\lambda_i)$$

on eigenvalues of random matrices distributed as

$$\mathcal{P}_{\beta}(\boldsymbol{\lambda}) = \frac{1}{\mathcal{Z}} e^{-\beta E(\boldsymbol{\lambda})},$$

$$E(\boldsymbol{\lambda}) = -\sum_{i < j} \log |\lambda_i - \lambda_j| + N \sum_i V(\lambda_i).$$

such that the spectral density has single support

# Problem

Behaviour of Covariance of A and B for N>>1



Constant value (dependent on \beta)... dependence on potential \text{V}?

# Solution

$$\operatorname{Cov}(A, B) \to \frac{1}{\beta \pi^2} \int_0^{+\infty} dk \, \varphi(k) \operatorname{Re} \left[ \tilde{a}(k) \tilde{b}^*(k) \right] + \mathcal{O}(N^{-1})$$

where

$$\varphi(k) = \begin{cases} k & \text{for } (\lambda_{-}, \lambda_{+}) = (-\infty, \infty) \\ k \tanh(\pi k) & \text{otherwise} \end{cases}$$

$$\tilde{f}(k) = \begin{cases} \int_{-\infty}^{+\infty} dx \, e^{ikx} f(x) & \text{for } (\lambda_{-}, \lambda_{+}) = (-\infty, \infty) \\ \int_{-\infty}^{+\infty} dx \, e^{ikx} f(T_{\lambda_{\pm}}(e^{x})) & \text{otherwise} \end{cases}$$

... depends only on the edges of the support of the spectral density

# Usually, interest is for a single linear statistics...

- Politzer (1989).
- Basor and Tracy (1993).
- Chen and Manning (1994).
- Costin and Lebowitz (1995).
- Baker and Forrester (1997).
- Pastur and Shcherbina (1997).
- Johansson (1998).
- Soshnikov (2000).
- Wieand (2002).
- Pastur (2006).
- Lytova and Pastur (2009).
- **●** + ......

#### Statistical Theory of the Energy Levels of Complex Systems. IV

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PHYSICAL REVIEW B

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#### Random-matrix theory of mesoscopic fluctuations in conductors and superconductors

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NUCLEAR PHYSICS B [FS]

Nuclear Physics B 422 [FS] (1994) 515-520

Universality of Brézin and Zee's spectral correlator

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#### 1. Dyson-Mehta (DM) [1]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \int_0^\infty \mathrm{d}k \ k |\hat{a}(k)|^2 + \mathcal{O}(N^{-1})$$

$$\hat{a}(k) = \int_0^\infty \mathrm{d}x \ \mathrm{e}^{\mathrm{i}kx} a(x)$$
2. Recordshort

$$A = \sum_{i=1}^{N} a(\lambda_i)$$

#### 2. Beenakker (B1) [2]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \int_0^\infty \mathrm{d}k \ k \tanh(\pi k) |\hat{F}(k)|^2 + \mathcal{O}(N^{-1})$$

$$\hat{F}(k) = \int_{-\infty}^{\infty} dx \, e^{ikx} a \left( \frac{1}{1 + e^x} \right)$$

#### 3. Beenakker (B2) [3]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \operatorname{Pr} \iint_{\lambda_{-}}^{\lambda_{+}} d\lambda d\mu \, \Theta(\lambda, \mu) \frac{a(\lambda)}{\lambda - \mu} \frac{d}{d\mu} a(\mu) + \mathcal{O}(N^{-1})$$

$$\Theta(\lambda,\mu) = \sqrt{\frac{(\mu - \lambda_{-})(\lambda_{+} - \mu)}{(\lambda - \lambda_{-})(\lambda_{+} - \lambda)}}$$

# Goals

- Generalize these formulae to the covariance case (2 linear statistics)
- Clarify when and where the classical variance formulae should/can be used

# Ingredients:

- (i) Conformal change of variables
- ii) Definition of covariance in terms of 2-point kernel
- iii) Functional derivative identity + electrostatic equation
- iv) Universality of the smoothed kernel on  $(0,\infty)$

(i)

$$\mathcal{P}_{\beta}(\boldsymbol{\lambda}) = \frac{1}{\mathcal{Z}} e^{-\beta \left[-\sum_{i < j} \ln |\lambda_i - \lambda_j| + N \sum_i V(\lambda_i)\right]} \equiv \frac{e^{-\beta E(\boldsymbol{\lambda})}}{\mathcal{Z}}$$

$$Cov(A, B) = \int_{\Lambda^N} d\boldsymbol{\lambda} \mathcal{P}_{\beta}(\boldsymbol{\lambda}) \left( A(\boldsymbol{\lambda}) - \langle A \rangle \right) \left( B(\boldsymbol{\lambda}) - \langle B \rangle \right) .$$

Let us consider the conformal map

$$\lambda = Tx$$
 :  $\lambda_i = \frac{ax_i + b}{cx_i + d}$  (20)

with  $ad - bc \neq 0$ . Then

$$\tilde{\mathcal{P}}_{\beta}(\boldsymbol{x}) = \frac{1}{\tilde{\mathcal{Z}}} e^{-\beta \tilde{E}(\boldsymbol{x})} , \qquad (21)$$

$$\tilde{E}(\boldsymbol{x}) = -\sum_{i < j} \log |x_i - x_j| + N \sum_i \tilde{V}(x_i) + \mathcal{O}(N)$$

$$Cov(A, B) = \int_{\tilde{\Lambda}^{N}} d\boldsymbol{x} \frac{1}{\tilde{\mathcal{Z}}} e^{-\beta \tilde{E}(\boldsymbol{x})} A(T(\boldsymbol{x})) B(T(\boldsymbol{x})) - \langle A \rangle \langle B \rangle$$

$$\tilde{\rho}_{N}(x) = \frac{1}{N} \sum_{i=1}^{N} \delta(x - x_{i}) \qquad \text{(random measure)}$$

$$\tilde{\rho}(x) = \lim_{N \to \infty} \langle \tilde{\rho}_{N}(x) \rangle \qquad \text{(single support)}$$

$$\tilde{\mathcal{K}}_{N}(x, x') = -\langle \tilde{\rho}_{N}(x) \tilde{\rho}_{N}(x') \rangle + \langle \tilde{\rho}_{N}(x) \rangle \langle \tilde{\rho}_{N}(x') \rangle$$

$$\tilde{\mathcal{K}}(x, x') = \lim_{N \to \infty} N^{2} \tilde{\mathcal{K}}_{N}(x, x')$$

(ii) 
$$\operatorname{Cov}(A, B) = -N^2 \iint dx dx' \tilde{\mathcal{K}}_N(x, x') a(T(x)) b(T(x'))$$
$$\to -\iint_{\sigma} dx dx' \tilde{\mathcal{K}}(x, x') a(T(x)) b(T(x')) + \mathcal{O}(1/N)$$

We can choose a,b,c,d in a convenient way...



# The support $\sigma$ in the large-N limit is moved to $(0,\infty)$

$$Cov(A, B) = -N^2 \iint dx dx' \mathcal{K}_N(x, x') a(T(x)) b(T(x'))$$

$$\rightarrow -\iint_{\sigma} dx dx' \mathcal{K}(x, x') a(T(x)) b(T(x')) + \mathcal{O}(1/N)$$

What is the kernel in this case?

# iii) Functional derivative identity...

$$\left| \tilde{\mathcal{K}}_{N}(x, x') = \frac{1}{\beta N^{2}} \frac{\delta \left\langle \tilde{\rho}_{N}(x) \right\rangle}{\delta \tilde{V}(x')} \left( 1 + \mathcal{O}(N^{-1}) \right) \right|$$

[Beenakker 1993]

$$\begin{split} &\frac{\delta \left\langle \tilde{\rho}_{N}(x) \right\rangle}{\delta \tilde{V}(x')} = \frac{\delta}{\delta \tilde{V}(x')} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x) \mathcal{P}_{\beta}(\boldsymbol{x}) \\ &= \frac{\delta}{\delta \tilde{V}(x')} \frac{1}{\tilde{\mathcal{Z}}} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x) \mathrm{e}^{-\beta \tilde{E}(\boldsymbol{x})} \\ &= \frac{-\beta}{\tilde{\mathcal{Z}}} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x) \mathrm{e}^{-\beta \tilde{E}(\boldsymbol{x})} \frac{\delta \tilde{E}(\boldsymbol{x})}{\delta \tilde{V}(x')} + \\ &+ \frac{\beta}{\tilde{\mathcal{Z}}^{2}} \iint \mathrm{d}\boldsymbol{x} \mathrm{d}\boldsymbol{x}'' \, \tilde{\rho}_{N}(x) \mathrm{e}^{-\beta (\tilde{E}(\boldsymbol{x}) + \tilde{E}(\boldsymbol{x}''))} \frac{\delta \tilde{E}(\boldsymbol{x}'')}{\delta \tilde{V}(x')} \\ &= -\beta N^{2} \left[ \frac{1}{\tilde{\mathcal{Z}}} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x) \tilde{\rho}_{N}(x') \mathrm{e}^{-\beta \tilde{E}(\boldsymbol{x})} - \right. \\ &\left. \left( \frac{1}{\tilde{\mathcal{Z}}} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x) \mathrm{e}^{-\beta \tilde{E}(\boldsymbol{x})} \right) \left( \frac{1}{\tilde{\mathcal{Z}}} \int \mathrm{d}\boldsymbol{x} \, \tilde{\rho}_{N}(x') \mathrm{e}^{-\beta \tilde{E}(\boldsymbol{x})} \right) \right] + \mathcal{O}(N) \; . \end{split}$$

## ... + electrostatic equation

for 
$$\tilde{\sigma} = (0, +\infty)$$
 support

$$\int_0^\infty \mathrm{d}y \tilde{\rho}(y) \ln|x-y| = \tilde{V}(x), \quad x > 0$$

$$ilde{
ho}(x) = \int_0^\infty \mathrm{d}y \; \Phi(x,y) ilde{V}(y)$$

$$ilde{\mathcal{K}}(x,x') = rac{1}{eta} rac{\delta ilde{
ho}(x)}{\delta ilde{V}(x')} = \Phi(x,x')$$

[Beenakker 1993] 
$$\Phi(x,x') = \partial_x \partial_{x'} \ln \left| \frac{\sqrt{x} + \sqrt{x'}}{\sqrt{x} - \sqrt{x'}} \right|$$

#### Multiloop correlators for two-dimensional quantum gravity

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# Universality of the correlations between eigenvalues of large random matrices

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# Summary of the main steps

- Covariance: double integral over finite-N kernel
- Conformal transformation maps the original support to  $(0,\infty)$  [unless it was  $(-\infty,\infty)$ ]
- Limit of the finite-N kernel is universal upon smoothing (it depends only on the edge points of the support of average density)
- Inserting the smoothed kernel for a support (0,∞) inside the double integral gives the result (after simple algebra in Fourier space)

$$\operatorname{Cov}(A, B) \to \frac{1}{\beta \pi^2} \int_0^{+\infty} dk \, \varphi(k) \operatorname{Re} \left[ \tilde{a}(k) \tilde{b}^{\star}(k) \right] + \mathcal{O}(N^{-1})$$

$$\tilde{f}(k) = \begin{cases} \int_{-\infty}^{+\infty} dx \, \mathrm{e}^{\mathrm{i}kx} f(x) & \text{for } (\lambda_{-}, \lambda_{+}) = (-\infty, \infty) \\ \int_{-\infty}^{+\infty} dx \, \mathrm{e}^{\mathrm{i}kx} f(T_{\lambda_{\pm}}(\mathrm{e}^{x})) & \text{otherwise} \end{cases}$$

- Invariant under the exchange  $A \longrightarrow B$
- It can be zero (decorrelation to leading order)

# Recovering classical variance formulas (1)

#### 1. Dyson-Mehta (DM) [1]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \int_0^\infty dk \ k |\hat{a}(k)|^2 + \mathcal{O}(N^{-1})$$
$$\hat{a}(k) = \int_{-\infty}^\infty dx \ e^{ikx} a(x)$$

Recovered when A=B for a matrix model whose spectral density has support over the full real axis [example: Cauchy ensemble]

# Recovering classical variance formulas (11)

#### 2. Beenakker (B1) [2]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \int_0^\infty dk \ k \tanh(\pi k) |\hat{F}(k)|^2 + \mathcal{O}(N^{-1})$$

$$\hat{F}(k) = \int_{-\infty}^{\infty} dx \, e^{ikx} a \left( \frac{1}{1 + e^x} \right)$$

Recovered when A=B for a matrix model whose spectral density has support over [0,1] [example: Jacobi ensemble]

# Recovering classical variance formulas (III)

### 3. Beenakker (B2) [3]

$$\operatorname{Var}(A) = \frac{1}{\beta \pi^2} \operatorname{Pr} \iint_{\lambda_{-}}^{\lambda_{+}} \mathrm{d}\lambda \mathrm{d}\mu \, \Theta(\lambda, \mu) \frac{a(\lambda)}{\lambda - \mu} \frac{\mathrm{d}}{\mathrm{d}\mu} a(\mu) + \mathcal{O}(N^{-1})$$

$$\Theta(\lambda,\mu) = \sqrt{\frac{(\mu - \lambda_{-})(\lambda_{+} - \mu)}{(\lambda - \lambda_{-})(\lambda_{+} - \lambda)}}$$

Recovered when A=B for a matrix model whose spectral density has finite support

# Conclusions

- Covariance formula for 2 (differentiable) linear statistics: it depends only on the edge points of the support of the average density
- Applicable to "Coulomb-gas" ensembles whose level density is supported on a single interval
- Conformal transformation maps the support to  $(0,\infty)$
- Smoothed 2-point kernel for (0,∞) is known and completely universal

$$\operatorname{Cov}(A, B) \to \frac{1}{\beta \pi^2} \int_0^{+\infty} dk \, \varphi(k) \operatorname{Re} \left[ \tilde{a}(k) \tilde{b}^{\star}(k) \right] + \mathcal{O}(N^{-1})$$

# Outlook

- Extensions to multi-cut ensembles?
- ...to non-invariant ensembles?
- ...to non-Hermitian ensembles?
- ...to biorthogonal ensembles?
- Non-differentiable linear statistics
- Central Limit Theorem for joint linear statistics?