

## Some Generalizations of Fourier Theory

see also [Dudok de Wit review](#)

Higher Order Spectra (summary)

Consider the nonlinear system:

$$\frac{\partial u(x,t)}{\partial x} = f(u(x,t))$$

decompose –

$$\begin{aligned} \frac{\partial u(x,t)}{\partial x} &= \int g(\tau_1)u(x,t-\tau_1)d\tau_1 \\ &+ \iint g(\tau_1,\tau_2)u(x,t-\tau_1)u(x,t-\tau_2)d\tau_1 d\tau_2 \\ &+ \iiint g(\tau_1,\tau_2,\tau_3)u(x,t-\tau_1)u(x,t-\tau_2)u(x,t-\tau_3)d\tau_1 d\tau_2 d\tau_3 \\ &+ \dots \end{aligned}$$

Taking the DFT we obtain the Volterra series:

$$\frac{\partial u_p}{\partial x} = \Gamma_p u_p + \sum_{k,l} \Gamma_{kl} u_k u_l \delta_{k+l,p} + \sum_{k,l,m} \Gamma_{klm} u_k u_l u_m \delta_{k+l+m,p} + \dots$$

with

$$u_p = u(x, \omega_p)$$

The leading term is a linear (cf Fourier) decomposition. The rest are mode coupling. Ensemble average over  $x$  and consider a homogeneous medium:

$$\begin{aligned} \text{setting } \frac{\partial}{\partial x} &= 0 \\ \Gamma_p \langle u_p^* u_p \rangle &+ \sum_{k+l=p} \Gamma_{kl} \langle u_k u_l u_{k+l}^* \rangle + \sum_{k+l+m=p} \Gamma_{klm} \langle u_k u_l u_m u_{k+l+m}^* \rangle + \dots = 0 \end{aligned}$$

Now recall convolution:

$$g_k * h_k = \sum_{u=0}^{N-1} g_u h_{k-u}$$

The DFT is  $G_m H_m$  where  $G_m$  is the DFT of  $g_k$  etc.,

This relates to cross correlation:  $C_\tau = \sum_{k=0}^{N-1} g_k h_{k+\tau}$  DFT is  $G_m^* H_m$

auto correlation:  $R_\tau = \sum_{k=0}^{N-1} x_k x_{k+\tau}$  DFT is  $S_m^* S_m$  (the power spectrum)

**generalise these to:**

bispectrum  $B_{kl} = S_k S_l S_{k+l}^*$

trispectrum  $T_{klm} = S_k S_l S_m S_{k+l+m}^*$

There are normalized versions, eg:

bicoherence =  $\frac{|B_{kl}|^2}{|S_k S_l|^2 |S_{k+l}|^2} = b_{kl}$

One can obtain averaged bispectra in the same way as averaged power spectra – average over  $M$  consecutive intervals.

Alternatively, use the 2nd order autocorrelation:  $M_{\tau_1 \tau_2} = \sum_{k=0}^{N-1} x_k x_{k+\tau_1} x_{k+\tau_2}$

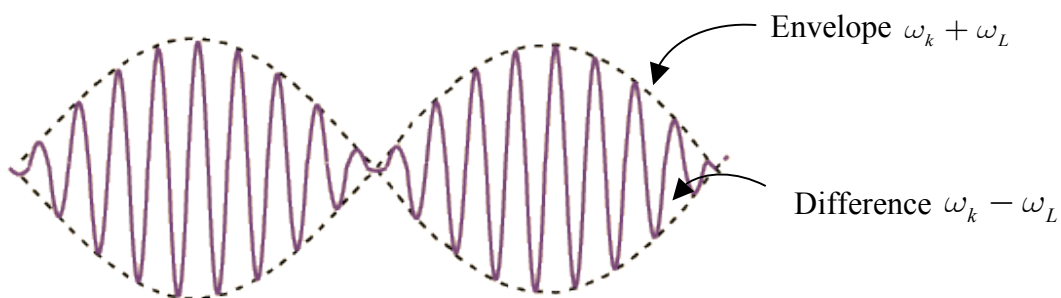
The bispectrum is the twice applied DFT on  $M$  (cf convolution theorem to prove).

Physical meaning (see also Dudok de Wit review)

Recall frequency  $f_m = \frac{m}{N \Delta t} \equiv \omega_m$

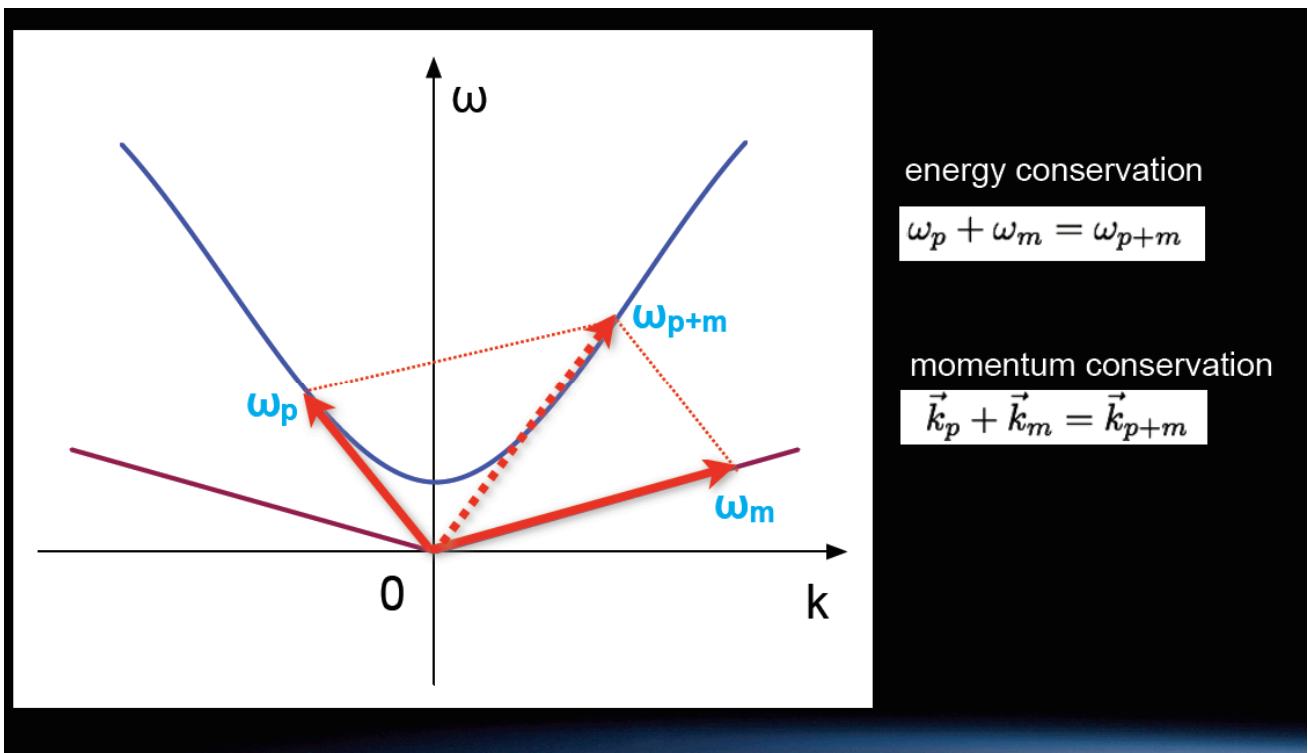
Thus :  $S_k \rightarrow S(\omega_k), S_l \rightarrow S(\omega_l), S_{k+l} \rightarrow S(\omega_k + \omega_l)$

This tests for coherence between beating oscillations

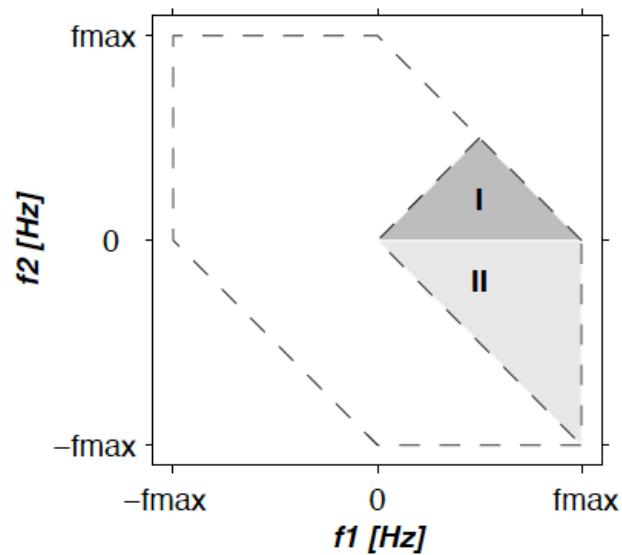


where  $\omega_{k,l} = \omega_k + \omega_l$ , for strongest mode we see strong bicoherence.

For wave- wave coupling we need to satisfy physical constraints in both space and time:

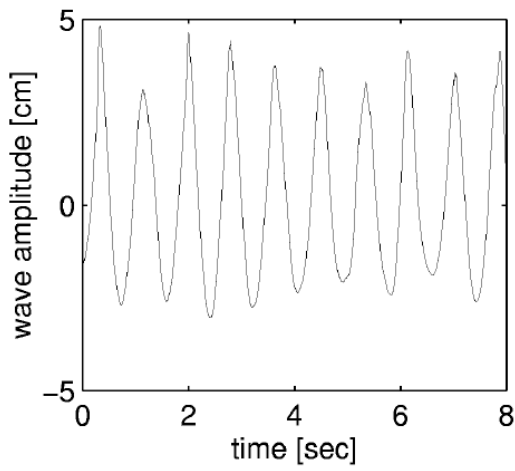


Principal domain of bicoherence- due to symmetries:

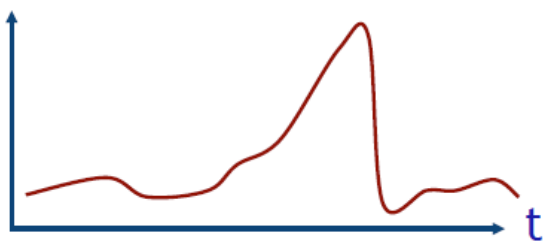
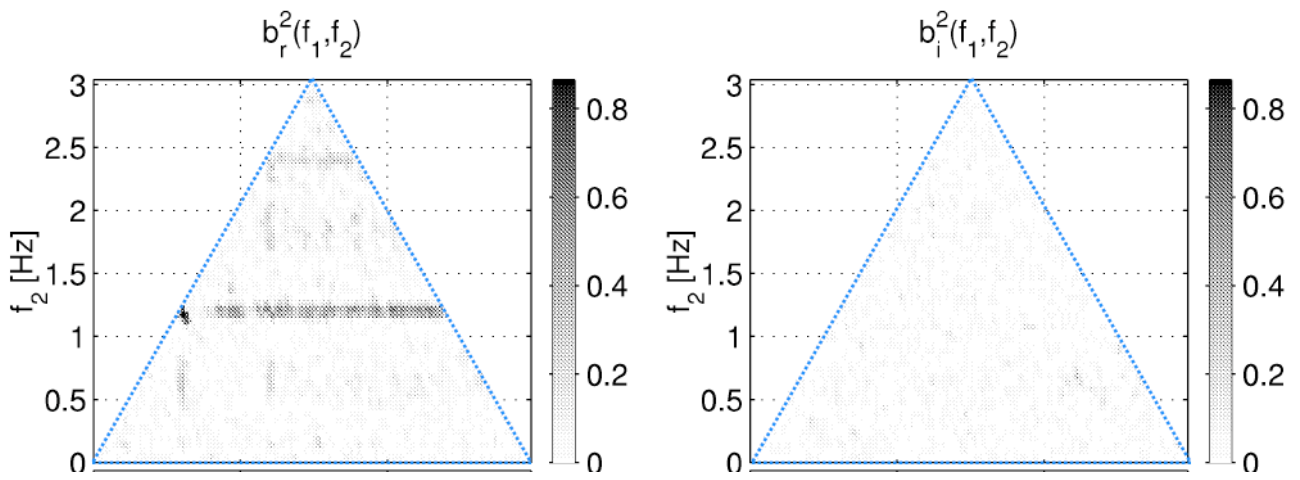


**Fig. 3.** Principal domain of the bicoherence. The Nyquist theorem restricts the display to the area enclosed by a dashed line. For the autobicoherence, the principal domain is I, for the cross-bicoherence it is I and II.

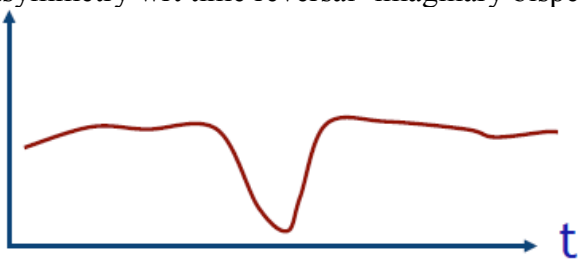
**Example- water waves: (courtesy T. Dudok De Wit)**



Real part and imaginary part of bicoherence



asymmetry wrt time reversal- imaginary bispectra



up- down asymmetry- real bispectra

so wave is not steepening- simply have a ‘wobble’ on amplitude.

### Linear Time Invariant (LTI) Filters

A way to generalize the Fourier world to other transforms.

So far – everything flowed from the idea:

$$x(t) = \sum_{m=-\infty}^{\infty} S_m e^{2\pi i f_m t} \quad f_m = \frac{m}{T}$$

with

$$S_m = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-2\pi i f_m t} dt$$

Discrete version:

$$x_k = \frac{1}{N\Delta t} \sum_{m=0}^{N-1} S_m e^{2\pi i m k / N} \quad f_m = \frac{m}{N\Delta t}$$

$$S_m = \Delta t \sum_{k=0}^{N-1} x_k e^{-2\pi i k m / N} \quad t_k = k\Delta t$$

and orthogonality property of the Fourier kernel  $\Phi = e^{2\pi i f t}$

A framework to consider other  $\Phi$  -

Write filter as:  $L[x(t)] = y(t)$

where  $L$  is a linear operator, with the properties:

- |    |                              |                                  |
|----|------------------------------|----------------------------------|
| 1. | scale preserving             | $L[ax] = aL[x]$                  |
| 2. | distributive (superposition) | $L[x_1 + x_2] = L[x_1] + L[x_2]$ |
| 3. | time invariant               | $L[x(t)] = y(t)$                 |
|    | $\Rightarrow$                | $L[x(t + \tau)] = y(t + \tau)$   |

Then in general:  $L\left[\sum_{p=1}^N \alpha_p x_p(t)\right] = \sum_{p=1}^N \alpha_p L[x_p(t)]$

Now consider for input  $x$  to the filter, the Fourier kernel:

$$\Phi_f(t) = e^{2\pi i f t} \quad f = \text{const}$$

$$y_f = L[\Phi_f] = y_f(t)$$

then:  $y_f(t + \tau) = L[\Phi_f(t + \tau)] = L[e^{2\pi i f \tau} \Phi_f(t)] = e^{2\pi i f \tau} y_f(t)$

This is just the "shift theorem" from Fourier theory.

Now let  $t = 0$ .

$$y_f(\tau) = e^{2\pi i f \tau} y_f(0)$$

true for any  $\tau$  so let  $\tau \dots t$

$$y_f(t) = e^{2\pi i f t} y_f(0)$$

so  $y_f(0)$  is a constant – there is one value for each  $f$ .

We can consider a spectrum of values of  $y_f(0)$

$$y_f(0) = G_f = A_f e^{i\theta_f}$$

Then

$$y_f(t) = G_f \cdot \Phi_f(t) = L[\Phi_f(t)] = G_f e^{2\pi i f t}$$

generally:

$$\left. \begin{array}{l} \Phi_f - \text{eigenvalues} \\ G_f - \text{eigenvectors} \end{array} \right\} \text{of } L$$

We can think of any transform in this way.

We need to identify an appropriate  $\Phi_f(t)$ . Two methods:

1. Choose  $\Phi_f(t)$  as basis on which we expand, ie:  $y(t) = \sum_f y_f(t) = \sum_f G_f \Phi_f(t)$

$\Phi_f$  may be orthogonal – chosen for "appropriate" properties.

This is equivalent to the transform:  $y(t) = \int_{-\infty}^{\infty} G(f) \Phi(f, t) df$

again,  $\Phi(f, t) = e^{2\pi i f t}$  for the Fourier transform.

2. Perform an SVD (single value decomposition, or principle component analysis) on  $L$ , so that the  $\Phi_f$  are generated by the data (beyond the scope of this course).

Application of LTI filters – coloured noises

Consider some  $x(t) = \sum_{m=-\infty}^{\infty} S_m e^{2\pi i f_m t}$

The discrete version of this is

$$x_k = \frac{1}{N\Delta t} \sum_{m=0}^{N-1} S_m e^{2\pi i m k / N} \quad f_m = \frac{m}{N\Delta t}$$

We now have:

$$y(t) = L[x(t)] = \sum_{m=-\infty}^{\infty} G_m S_m e^{2\pi i f_m t}$$

$$y_k = \frac{1}{N\Delta t} \sum_{m=0}^{N-1} G_m S_m e^{2\pi i m k / N}.$$

Now for a stochastic process we have

$$x_k = \frac{1}{N\Delta t} \sum_{m=0}^{N-1} D_m e^{2\pi i m k / N}$$

where

$$D_m = S_m e^{i\phi_m} \quad D_m, \phi_m \text{ are random iid processes (these are stationary)}$$

then as above

$$y_f(t) = G_f x_f(t)$$

or

$$y_{k,m} = G_m x_{k,m}$$

$$\text{where } x_{k,m} = S_m e^{2\pi i m k / N}$$

The  $G_f$  are just constants.

Then for process  $x_k$  stochastic, there will be a process  $y_k$  also stochastic, with spectral components  $G_f D_f$ .

The  $x_k, y_k$  should share the statistical properties .

More formally, since the R-S integral gives, for stochastic  $dz_x$  (see notes on stationarity)

$$x(t) = \int_{-1/2}^{1/2} e^{2\pi i f t} dz_x(f)$$

then

$$y(t) = \int_{-1/2}^{1/2} e^{2\pi i f t} G(f) dz_x(f)$$

$$= \int_{-1/2}^{1/2} e^{2\pi i f t} dz_y(f)$$

ie  $\langle |dz_y(f)|^2 \rangle = G^2(f) \langle |dz_x(f)|^2 \rangle$  so this filter generates "coloured" noises if  $G \sim \frac{1}{f^\beta}$