## Maps and chaos

So far have discussed continuous 1st, 2nd order DE

- found trajectories in phase space are "well behaved"

only possibility for divergence of trajectories.


This does not permit stochastic (mixing up/shuffling) diffusion in phase space. This is only possible (in plane of paper) if we add another dimension(out of plane of paper) ie in 3 or more dimensions.

Now look at simplest systems that show route to chaos:

## Difference equations

i) Show how they "work" - what are their properties; example of chaos.
ii) Show how they are fundamentally different to the analogous differential equations.
iii) Quantify the route to chaos - Lyapunov exponents.
iv) Introduce Universality - the bifurcation route to chaos, Feigenbaum numbers.

1D Noninvertible maps - the easiest place to start.

Piecewise linear 1D maps - the tent map

$$
\begin{aligned}
x_{n+1} & =1-2\left|x_{n}-\frac{1}{2}\right| \\
& \equiv M(x)
\end{aligned}
$$

- tells us how $x+1^{\text {th }}$ step depends on

$$
x^{t h} \text { step }
$$

Write as:

$$
x_{n+1}=\begin{array}{c|c}
2 x_{n} & x_{n} \leq \frac{1}{2} \\
2\left(1-x_{n}\right) & x_{n} \geq \frac{1}{2}
\end{array}
$$

What happens when we iterate the map?
If (i) $\quad x_{n}<0 \quad x_{n+1}<0$ and $\left|x_{n+1}\right|>\left|x_{n}\right|$
ie: $x \rightarrow-\infty$
If (ii) $x_{n}>1 \quad x_{n+1}<0$, so $x_{n+2}<0 \quad\left|x_{n+2}\right|>\left|x_{n+1}\right|$ as in (i) ie: $x \rightarrow-\infty$.

So, just consider the interval $x=[0,1]$ - iterates of $x$ are bounded in this range.
We can still look for fixed points, linearise to examine stability as before $\rightarrow$ local behaviour.
Interesting difference will be in global behaviour of maps.

Fixed points are where $x_{n}$ doesn't change as $n \rightarrow \infty$ so

$$
\begin{aligned}
\bar{x}_{n+1} & =\bar{x}_{n} \text { is fixed point } \\
& =\bar{x} \quad M(x)=x
\end{aligned}
$$

Graphically:


Fixed point is in range

$$
\begin{aligned}
\frac{1}{2} & \leq x_{n} \leq 1 \quad \text { ie: } \quad x_{n+1}=2\left(1-x_{n}\right) \\
\bar{x} & =2(1-\bar{x}) \\
x & =2-2 \bar{x} \\
\bar{x} & =\frac{2}{3}
\end{aligned}
$$

so fixed point is
also in range $0 \leq x_{n} \leq \frac{1}{2} \quad$ fixed point $\quad \bar{x}=0$ here $\quad x_{n+1}=2 x_{n}$

Stability/classification. Can still write

$$
\begin{aligned}
x_{n} & =\bar{x}+\delta x_{n} \\
x_{n+1} & =\bar{x}+\delta x_{n+1}
\end{aligned}
$$

ie: $\delta x$ is small but now
discontinuous.
sub in to

$$
\begin{aligned}
& x_{n+1}=2\left(1-x_{n}\right) \\
& \bar{x}+\delta x_{n+1}=2\left(1-\bar{x}-\delta x_{n}\right) \\
& \quad \delta x_{n+1}=2-3 \bar{x}-2 \delta x_{n} \\
& \quad=-2 \delta x_{n}
\end{aligned}
$$

$$
\bar{x}=\frac{2}{3} \quad \delta x_{n+1}=2-3 \bar{x}-2 \delta x_{n}
$$

So

$$
\delta x_{n+1}=-2 \delta x_{n}
$$

also

$$
\delta x_{n}=-2 \delta x_{n-1}
$$

so

$$
\begin{aligned}
\delta x_{n+1} & =-2.2 \delta x_{n-1}=-2.2 .2 \delta x_{n-2} \\
& =[-2]^{j+1} \delta x_{n-j}=[-2]^{n+1} \delta x_{0} \text { where } \delta x_{0} \text { is the initial condition }
\end{aligned}
$$

Hence, the fixed point is unstable - oscillates.
Similarly, $\quad \bar{x}=0, \delta x_{n+1}=2 \delta x_{n}$ - unstable $\quad\left[x_{N+1}=2 x_{N}\right]$

## Consider global behaviour.

Iterate many times
Look at one interate: $\quad x_{n+1}=M\left(x_{n}\right)$
(some notation here)
two iterates:

$$
\begin{aligned}
& x_{n+2}=M^{2}\left(x_{n}\right) \\
& \quad=M\left(M\left(x_{n}\right)\right)
\end{aligned}
$$

p iterates

$$
x_{n+p}=M^{p}\left(x_{n}\right)
$$

where | $M(x)=\begin{array}{c}2 x \\ 2(1-x)\end{array}$ | $\left.\begin{array}{l}x \\ \\ x\end{array}\right) \frac{1}{2}$ |
| :---: | :---: |
| $x$ |  |

## Fixed points - notice that




No of fixed points doubles each iterate.

Another way of looking at $M^{p}(x)$ graphically:


$$
\begin{aligned}
& \text { vertical lines are } M\left(x_{n}\right) \\
& \text { horizontal } \quad x_{n+1}=M\left(x_{n}\right)
\end{aligned}
$$

note that the iterates are 'shuffled'

