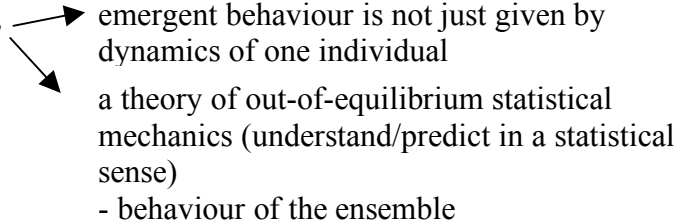


High dimensional systems

Properties

- many interacting elements
- 'microscopic' rules for element-element interaction
- order and control parameters
- emergence (a class of which is scaling)
- systems may have dynamical steady states but in general may not be in equilibrium

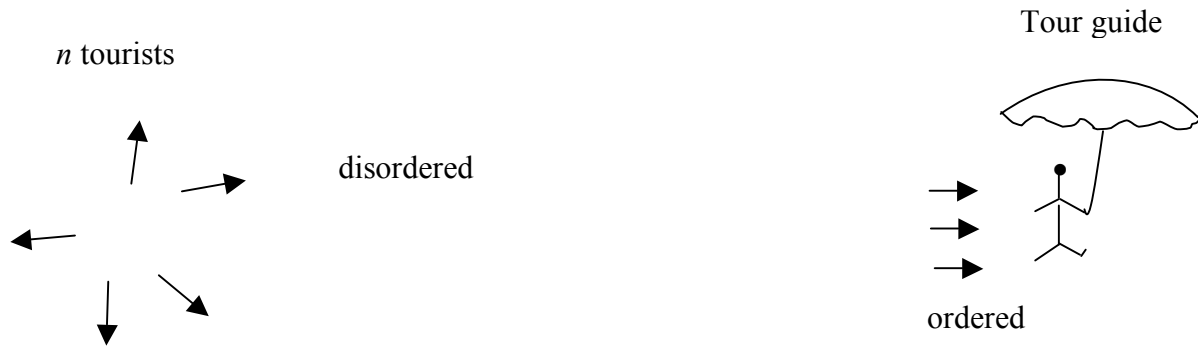
- complex systems 
 - emergent behaviour is not just given by dynamics of one individual
 - a theory of out-of-equilibrium statistical mechanics (understand/predict in a statistical sense)
 - behaviour of the ensemble

Examples- flocking birds have a microscopic rule: competition between order ('wanting to be like your neighbour') and disorder ('doing your own thing/random perturbation') and order/control parameters for the macroscopic behaviour that emerges (flocking)

High dimensional systems- we will do the following

1. Introduction – self-organisation (when a many d.o.f. system looks like a few d.o.f. system).
2. Order - disorder transitions (met one already- ferromagnet) → order/control parameters.
3. Emergence and self-similarity (fractals)
4. Dimensional analysis- Buckingham Pi theorem
 - how to identify order and control parameters.

Self organisation ("enslaving")



Motion of tourists 'enslaved' to tour guide.

Tour guide = order parameter.

- just follow the guide
- clearly desirable to treat system in this way –

just solve 1 equation – tour guide
instead of N – tourists.

How can this occur?

Consider tourist ODE:

$$\frac{dq_i}{dt} = -\alpha_i q_i + A e^{-\lambda t} \quad \begin{matrix} i = 1 \dots n \\ \alpha, \lambda > 0 \end{matrix}$$

q_i is i th tourist position driving term \equiv tour guide

each tourist solution obtained by integrating directly – use integrating factor

ie: $\frac{dy}{dt} + f(t)y = g(t) \quad I.F. \quad \exp\left[\int f(t) dt\right]$

multiply by I.F. gives a perfect integral.

So, here – dropping i subscript for the moment

$$\frac{dq}{dt} + \alpha q = A e^{-\lambda t} \quad 1.F. \quad e^{\int \alpha dt} = e^{\alpha t}.$$

So we have

$$\frac{dq}{dt} e^{\alpha t} + \alpha q e^{\alpha t} = A e^{(\alpha-\lambda)t}$$

integrate:

$$q e^{\alpha t} = \frac{A}{\alpha - \lambda} e^{(\alpha-\lambda)t} + q_0 \quad q_0 = \text{constant of integration}$$

or

$$q = \frac{A}{\alpha - \lambda} e^{-\lambda t} + q_0 e^{-\alpha t}$$

so, for each tourist:

$$q_i = \frac{A}{\alpha_i - \lambda} e^{-\lambda t} + q_{0i} e^{-\alpha_i t}$$

Now, if tourist timescale $\frac{1}{\alpha_i} \ll \frac{1}{\lambda}$ tour guide timescale, then for large t , $\alpha_i t \gg \lambda t$ (and $\alpha_i, \lambda > 0$)

$$q_i \approx \frac{A}{\alpha_i} e^{-\lambda t} - \text{tourists just follow the tour guide.}$$

Note that the competition between the terms:

$$\begin{array}{ccc} e^{-\lambda t} & \text{and} & e^{-\alpha_i t} \\ \text{order} & & \text{disorder} \end{array}$$

is competition between order and disorder – and there is a parameter $\frac{\alpha_i}{\lambda} = \pi$

that tells you whether ordered $\pi \gg 1$
 disordered $\pi \ll 1$

NB: could have instead sought stationary solution

$$\frac{dq_i}{dt} = 0 = -\alpha_i q_i + A e^{-\lambda t}$$

$$\rightarrow q_i = \frac{A}{\alpha_i} e^{-\lambda t} \quad - \text{same result!}$$

Hence, the 'enslaved' solution is the time independent solution (steady state) – tourist solution quickly relaxes to follow the tour guide.

General statement of same thing

Any set of equations

$$\frac{dq_i}{dt} = H_i(q_j(t)) + S_i(t) \quad \begin{array}{l} i = 1 \dots N \\ j = 1 \dots N \end{array}$$

ie: H_i are a function of the q_i and the S_i are the external driving terms.

IF all q_i interact on timescales \ll characteristic timescale of S_i solution is in "dynamic equilibrium"

$$H_i(q_j(t)) = -S_i(t)$$

Now just need to do algebra, not integrate ODEs.

Self organisation

No tour guide – just tourists, but they interact with each other.

Write equations in terms of tourists c of m q_c

$$\frac{dq_i}{dt} = H_i(q_j) = \alpha_i(q_i) + G(q_c)$$

where $q_c = \sum_N q_j$ say

Again, if q_c has slower timescale than all q_i

- q_i enslaved to q_c