

### Neutrino Mass Measurements

### $\beta$ decay



Measurement of v mass from kinematics of  $\beta$  decay.



# Requirements for experiment



•The number of electrons close to the endpoint should be small

Good (and well-understood) electron energy resolution
No (or minimal) electron energy loss within the source
Minimal atomic and nuclear final state effects, of excited transitions

Gaseous Tritium: 
$${}^{3}H \rightarrow {}^{3}He^{+} + e^{-} + \overline{v_{e}}$$

Endpoint is at 18574 eV No molecular excitation above 18547 eV Only 10<sup>-9</sup> electrons in this region Gaseous so you can have a very large source



# History of tritium-β decay results

ITEP	m <sub>v</sub>		
T <sub>2</sub> in complex molecule magn. spectrometer (Tret'yakov)	17-40 eV		experimental results
Los Alamos		100	
gaseous T <sub>2</sub> - source magn. spectrometer (Tret'yakov)	< 9.3 eV	50 2	T
Tokio	< 12 1 01/	e 0	╶╎╴╴╴╴│┽┯╴┰╴┑┺╴╶╺╸
T - source magn. spectrometer (Tret'yakov)	< 15.1 ev	ີ <sub>-50</sub>	
Livermore		-100	Los Alamos
gaseous T <sub>2</sub> - source magn. spectrometer (Tret'yakov)	< 7.0 eV	-150	Tokio
Zürich		ŀ	Troitsk
T <sub>2</sub> - source impl. on carrier magn. spectrometer (Tret'yakov)	< 11.7 eV	-200	■ Troltsk (step)
Troitsk (1994-today)		-250	- electrostatic
gaseous T <sub>2</sub> - source electrostat. spectrometer	< 2.05 eV	-300	magnetic spectrometers
Mainz (1994-today)		-350	
frozen T <sub>2</sub> - source electrostat. spectrometer	< 2.3 eV		1986 1988 1990 1992 1994 1996 1998 2000 year

# Mainz Experiment



The current standard for tritium beta decay experiments



2π acceptance
High energy resolution

**Electrostatic Filter** 

$$\frac{\Delta E}{E} \sim 0.03\%$$

### **Present Status**





#### Troitsk

windowless gaseous T<sub>2</sub> source analysis 1994 to 1999, 2001  $m_v^2 = -2.3 \pm 2.5 \pm 2.0 \text{ eV}^2$  $m_v \leq 2.2 \text{ eV}$  (95% CL.) Mainz

quench condensed solid T<sub>2</sub> source

analysis 1998/99, 2001/02

 $m_v^2$  = - 1.2 ± 2.2 ± 2.1 eV<sup>2</sup>

 $m_{\nu} \leq 2.2 \text{ eV} (95\% \text{ CL.})$ 

Both experiments have reached the intrinsic limit of their sensitivity.

### KATRIN







#### KATRIN on the move /ΔR\ DENMARK LIPUANIA ULADO UNITED Vilnius. Sea Belfas Mahilvow Copenhagen Minsk Kaliningrad RUSSIA Bornholm BELARUS Man Dublin Irish (U.K.) Hrodna Leed Gdańsk •Manchester Liverpool Homyel RELAND Hamburg KINGDOM Brest Warsaw Cardiff Birmingham Amsterdan Bremen Berlin Poznań POLAND Kiev Rotterdam Lódź Rivne London' Celtic Essen Leipzig\_ Wrocław Cologne Sea Brussels+ UKRAINE BOANGERMANY Lille ·L'viv Prague BEL. Kraków uernsey (U.K.) Frankfurt am Main Jersey (U.K.) Chernivtsi CZECH REPUBLIC Luxembourg Mykolayiv SLOVAKIA Paris Brno LUX. Chisinau Stuttgart Bratislava lasi • Strasbourg Munich Cluj-Odesa Budapest Vienna Napoca 1 MOLDOVA Nantes Láre LIECH. AUSTRIA HUNGARY Zürich<sup>®</sup> + Bern + Vaduz ROMANIA GenevaSWITZ FRANCE Bay of o Ljubljana Constanța Bucharest Biscav SLOVENIA Zagreb MilanVenice MASSIF Lyon Black Turin BOSNIA AND HERZEGOVINA Bordeaux CENTRAL Varna Belgrade Sea Genoa CROATIA Bilbao MARNO Serbia Toulouse Sarajevo BULGARIA MONACO Ligurian Florence Andorra la Vella Marseille \*Sofia PYRENI Istanbul, Porto Adriatic ITALY ★ Skopje Sea Podgorica Zaragoza ANDORRA Corsica Rome F.Y.R.O.M. Bursa Thessaloniki Tirana VATICAN Madrid Barcelona PORTUGAL CITY TURKEY Naples Balearic \*Lisbon Sea SPAIN Tyrrhenian Sardinia man Izmir Valencia GREECE Sea BALEARIC Athens 0 **ISLANDS** Cagliari Sevilla Ionian Palermo Sea Mediterranean Sea Gibraltar Málaga Rhodes Sicily Strait of Gibrali Algiers Ceuta Alborán Sea (SPAIN) Scale 1: 19,500,000 Crete Oran Tunis Melilla Lambert Conformal Conic Projection, (SPAIN) Valletta 🗯 Rabat standard parallels 40°N and 56°N MALTA TUNISIA 300 Kilometers Casablanca ALGERIA MOROCCO 300 Miles

### KATRIN on the move





### Katrin data



![](_page_10_Figure_2.jpeg)

![](_page_11_Picture_0.jpeg)

# **KATRIN Sensitivity**

![](_page_11_Figure_2.jpeg)

3 year run period

sensitivity (90% CL) m(v) < 0.2 eV

discovery potential  $m(v) = 0.35 \text{ eV} (5\sigma)$ 

Starts in 2016

![](_page_12_Figure_0.jpeg)

 $m_{\pi} = 139.56995 \pm 0.00035 \, MeV$  $m_{\mu} = 105.658358 \pm 0.000005 \, MeV$  $p_{\mu} = 29.792 \pm 0.00011 \, MeV$ 

$$m_{v}^{2} = (-0.016 \pm 0.023) MeV^{2}$$

$$m_{v} < 190 \, keV(90 \,\% \, CL)$$

![](_page_13_Figure_0.jpeg)

![](_page_13_Picture_1.jpeg)

# Cosmology

![](_page_14_Picture_1.jpeg)

![](_page_14_Picture_2.jpeg)

Density fluctuations are affected by neutrino mass in the early universe
Highly model dependent
WMAP,2dF,ACBAR, CBI

![](_page_14_Picture_4.jpeg)

# $2\nu\beta\beta$ Decay

![](_page_15_Picture_1.jpeg)

Neutrinoless double beta decay is considered a golden channel for the measurement of neutrino mass.

In some nuclei β decay is forbidden but double beta decay is not

 $(Z, A) \rightarrow (Z+2, A) + 2e^{-} + 2\overline{\nu_e}$ 

![](_page_15_Figure_5.jpeg)

![](_page_16_Picture_0.jpeg)

![](_page_16_Picture_1.jpeg)

![](_page_16_Figure_2.jpeg)

Second order process in perturbation theory

Severe test for nuclear matrix element calculation

 Nuclear structure effects cause variations in the nuclear matrix elements of factors of 10

# $2\nu\beta\beta$ Decay

![](_page_17_Picture_1.jpeg)

$2\nu\beta\beta$ mode	Half life ( $\times 10^{24}$ years)
${}^{48}_{20}Ca \rightarrow {}^{48}_{22}Ti$	4.1
$^{76}_{32}Ge \rightarrow ^{76}_{34}Se$	40.9
${}^{82}_{34}Se \rightarrow {}^{82}_{36}Kr$	9.3
${}^{96}_{40}Zr \rightarrow {}^{96}_{42}Mo$	4.4
$^{100}_{42}Mo \rightarrow ^{100}_{44}Ru$	5.7
${}^{1\overline{10}}_{46}Pd \rightarrow {}^{1\overline{10}}_{48}Cd$	18.6
${}^{116}_{48}Cd \rightarrow {}^{116}_{50}Sn$	5.3
$\tilde{124}_{50}Sn \rightarrow \tilde{124}_{52}Te$	9.5
$\overset{130}{52}Te \rightarrow \overset{130}{54}Xe$	5.9
$\tilde{136}_{54}Xe \to \tilde{136}_{56}Ba$	5.5
${}^{150}_{60}Nd \rightarrow {}^{150}_{62}Sm$	1.2

# Neutrinoless $\beta\beta$ Decay

![](_page_18_Picture_1.jpeg)

![](_page_18_Figure_2.jpeg)

$$\mathbf{v}_L = \mathbf{v}_{h=-1} + \frac{m}{E} \mathbf{v}_{h=+1}$$

helicity states

- Neutrino must have mass
- Neutrino is a *Majorana* particle
- Violation of lepton number conservation
- Experiments are crucial to understanding

$$\Gamma_{0\nu} = G_{0\nu} |M_{0\nu}|^2 |\sum_i |U_{ei}|^2 m_i|^2 \Rightarrow T_{1/2} \sim 10^{27} years$$

![](_page_19_Picture_0.jpeg)

# What is the signal?

![](_page_19_Figure_2.jpeg)

## Experimental Requirements

![](_page_20_Picture_1.jpeg)

![](_page_20_Figure_2.jpeg)

(0vββ T<sub>1/2</sub> ~ 10<sup>26</sup> - 10<sup>27</sup> years)

#### Requires

Large, highly efficient source mass

- detector as source

Best possible energy resolution

- minimize 0vββ peak ROI to maximize S/B

- separate from 0vββ from irreducible 2vββ (~ T<sub>1/2</sub> ~ 10<sup>19</sup> - 10<sup>21</sup> years)

Extremely low (near-zero) backgrounds in the 0vßß peak region

- requires ultra-clean radiopure materials
- the ability to discriminate signal from background

![](_page_20_Figure_14.jpeg)

# Types of experiments

![](_page_21_Picture_1.jpeg)

![](_page_21_Figure_2.jpeg)

#### 1. the source is inserted as thin foil inside a tracking detector

- 2e<sup>-</sup> are detected separately
  - $\rightarrow$  different channels of 0vDBD can be distinguished
- particle identification
  - $\rightarrow$  background suppression
- poor energy resolution
  - $\rightarrow$  important 2vDBD background
    - (limitation on isotope choice)

![](_page_21_Picture_11.jpeg)

#### 2. the detector is itself the source

- solid state detectors
  - → several candidates, high resolution no info on kinematic techniques for background suppression
- gaseous detectors for Xe

![](_page_22_Picture_0.jpeg)

### Passive Source - NEMO3

![](_page_22_Picture_2.jpeg)

<u>Source</u>: 10 kg of ββ isotopes cylindrical, S = 20 m<sup>2</sup>, 60 mg/cm<sup>2</sup>

Tracking detector:

drift wire chamber operating in Geiger mode (6180 cells) Gas: He + 4% ethyl alcohol + 1% Ar + 0.1% H<sub>2</sub>O

<u>Calorimeter</u>: 1940 plastic scintillators coupled to low radioactivity PMTs

Magnetic field: 25 Gauss Gamma shield: Pure Iron (18 cm) Neutron shield: borated water + Wood

**Background:** n **Able to identify e<sup>-</sup>, e<sup>+</sup>, \gamma and \alpha** 2.6 MeV)

#### Typical ββ2ν event observed from <sup>100</sup>Mo

![](_page_23_Figure_1.jpeg)

Side view

Advantage : electron tracking Disadvantage : less source material and worse energy resolution

# Cuoricino/Cuore

![](_page_24_Picture_1.jpeg)

![](_page_24_Figure_2.jpeg)

![](_page_25_Figure_0.jpeg)

Energy

 $T_{1/2}^{0\nu} > 3.0 \times 10^{24} years \Rightarrow \langle m_{\nu} \rangle < 0.68 eV$ 

### Heidelberg-Moscow

![](_page_26_Picture_1.jpeg)

11 kg of Ge enriched to 86% of <sup>76</sup>Ge in the form of 5 Ge diodes surrounded by Cu,Pb,Bn shielding  $0\nu\beta\beta$  electrons detected by Ge detectors themselves Only sum of electron energy measured

![](_page_26_Picture_3.jpeg)

![](_page_26_Figure_4.jpeg)

### Heidelberg-Moscow

![](_page_27_Picture_1.jpeg)

11 kg of Ge enriched to 86% of <sup>76</sup>Ge in the form of 5 Ge diodes surrounded by Cu,Pb,Bn shielding  $0\nu\beta\beta$  electrons detected by Ge detectors themselves Only sum of electron energy measured

![](_page_27_Picture_3.jpeg)

![](_page_27_Figure_4.jpeg)

![](_page_28_Picture_0.jpeg)

### Heidelberg-Moscow

![](_page_28_Figure_2.jpeg)

### GERDA

![](_page_29_Picture_1.jpeg)

Designed to test Heidelberg-Moscow

Uses the same Ge-76 isotope and technique

Been running since 2010

![](_page_29_Picture_5.jpeg)

### GERDA

![](_page_30_Picture_1.jpeg)

![](_page_30_Figure_2.jpeg)

Inconsistent with HdM, but not definitive (yet)

![](_page_31_Picture_0.jpeg)

### Future Program

![](_page_31_Figure_2.jpeg)

![](_page_32_Picture_0.jpeg)

### Current global limit

![](_page_32_Figure_2.jpeg)

## Absolute mass status

![](_page_33_Picture_1.jpeg)

•Tritium β decay	$< m_{\beta} >$	< 2.3 eV
	Katrin extends	sensitivity to 0.2 eV
•0v2β decay	$\left \sum_{i} U_{ei}^2 m_i\right  <$	<0.3-1.2 eV mββ><300 meV
<ul> <li>Cosmology</li> </ul>	$\sum_{i=e,\mu,\tau} m_i <$	< 0.7 eV
<ul> <li>Pion decay</li> </ul>	$m_{v_{\mu}} < 190  keV$	
•Tau decay	$m_{_{ m v} au} < 18.2  MeV$	-

# Requirements for

![](_page_34_Picture_1.jpeg)

•The number of the small

Good (and well-understood) electron energy resolution
No (or minimal) electron energy loss within the source
Minimal atomic and nuclear final state effects, of excited transitions

Gaseous Tritium: 
$${}^{3}H \rightarrow {}^{3}He^{+} + e^{-} + \overline{v_{e}}$$

Endpoint is at 18574 eV No molecular excitation above 18547 eV Only 10<sup>-9</sup> electrons in this region Gaseous so you can have a very large source

![](_page_35_Picture_0.jpeg)

![](_page_35_Figure_1.jpeg)

![](_page_35_Figure_2.jpeg)

From Zdesenko, Danevich, Tretyak, J. Phys. G 30 (2004) 971

![](_page_36_Picture_0.jpeg)

### The Ideal Detector

Source serves as the detector

Elemental (enriched) source to minimize active material.

- Large Q value faster 0 rate and also places the region of interest above many potential backgrounds.
- Relatively slow  $2\nu\beta\beta$  rate helps control this irreducible background. Direct identification of the decay progeny in coincidence with the  $0\nu\beta\beta$

decay eliminates all potential backgrounds except  $2\nu\beta\beta$ .

 Full Event reconstruction, providing kinematic data such as opening angle and individual electron energy aids in the elimination of backgrounds and demonstration of signal (can possibly use 2νββ)
 Spatial resolution and timing information to reject background

processes.

Demonstrated technology at the appropriate scale.

The nuclear theory is better understood in some isotopes than others. The interpretation of limits or signals might be easier to interpret for some isotopes.

### NFM03

![](_page_37_Picture_1.jpeg)

- X Source serves as the detector
- Elemental (enriched) source to minimize active material.
  - Large Q value faster 0 rate and also places the region of interest above many potential backgrounds.

Relatively slow  $2\nu\beta\beta$  rate helps control this irreducible background.

- Direct identification of the decay progeny in coincidence with the 0vββ decay eliminates all potential backgrounds except 2vββ.
- Full Event reconstruction, providing kinematic data such as opening angle and individual electron energy aids in the elimination of backgrounds and demonstration of signal (can possibly use 2νββ)
- Spatial resolution and timing information to reject background processes.

Demonstrated technology at the appropriate scale.

The nuclear theory is better understood in some isotopes than others. The interpretation of limits or signals might be easier to interpret for some isotopes.

![](_page_38_Picture_0.jpeg)

# Dirac vs Majorana

A Dirac particle is different from it's antiparticle (e.g. electron and positron). A majorana particle is the same as it's antiparticle.

![](_page_38_Figure_3.jpeg)

![](_page_39_Picture_0.jpeg)

![](_page_39_Figure_1.jpeg)

![](_page_39_Figure_2.jpeg)

From Zdesenko, Danevich, Tretyak, J. Phys. G 30 (2004) 971

![](_page_40_Picture_0.jpeg)

•Electromagnetic, strong and weak forces have very different strengths

If supersymmetry is valid their strengths are the same at around  $10^{16}$  GeV

•To explain light neutrino masses through the see-saw mechanics, we need a heavy neutrino with mass 10<sup>16</sup> GeV

 Probing of GUT scale physics using light neutrinos!

![](_page_40_Figure_5.jpeg)

![](_page_41_Picture_0.jpeg)

### SN1987A

![](_page_41_Picture_2.jpeg)

![](_page_42_Picture_0.jpeg)

### Neutrino detectors operating at the time Kamiokande II, IMB, BST, Mont Blanc

![](_page_42_Figure_2.jpeg)

Relative Time (seconds)

### Mass from Velocity

The neutrinos had travelled 180,000 light years – enough for small mass differences to show up as a difference in arrival times

 $\Lambda/ARM/ICK$ 

$$t_{F} = t - t_{0} = \frac{L}{v} = \frac{L}{c} \frac{E_{v}}{p_{v}} c \sim \frac{L}{c} \left( 1 + \frac{m_{v}^{2}c^{4}}{2E^{2}} \right)$$

$$\delta t = t_j - t_i = \delta t_0 + \frac{L m_v^2}{2c} \left( \frac{1}{E_j^2} - \frac{1}{E_i^2} \right)$$

Estimate dependent on models of supernova process (emission intervals, size of the neutrino shell etc)

$$m_{\overline{v_e}} < 5.7 \, eV(95 \, CL)$$

![](_page_44_Figure_0.jpeg)