

Flavour Permutation Symmetry, Mixing and Jarlskog Invariance

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Outline of Talk

- The Flavour Problem
- Jarlskog Covariance
- Plaquette Invariance and Flavour Permutation Symmetry
- Flavour Symmetric Mixing Observables
- Applications
- Conclusions

The Flavour Problem

- SM has 20 (22) low energy flavour parameters (out of 25(27))
- They show tantalising structure, eg.

$$\begin{array}{c}
 \begin{array}{c}
 u \\
 c \\
 t
 \end{array}
 \begin{array}{c}
 d \\
 s \\
 b
 \end{array}
 \begin{array}{c}
 \mathcal{O}(1) \\
 \mathcal{O}(\lambda) \\
 \mathcal{O}(\lambda^3)
 \end{array}
 \begin{array}{c}
 \mathcal{O}(\lambda) \\
 \mathcal{O}(1) \\
 \mathcal{O}(\lambda^2)
 \end{array}
 \begin{array}{c}
 \mathcal{O}(\lambda^3) \\
 \mathcal{O}(\lambda^2) \\
 \mathcal{O}(1)
 \end{array}
 \end{array}
 \begin{array}{c}
 \nu_1 \\
 \nu_2 \\
 \nu_3
 \end{array}
 \end{array}
 \begin{array}{c}
 \xrightarrow{q_i} \\
 \xrightarrow{l_i}
 \end{array}
 \begin{array}{c}
 W^+ \\
 W^-
 \end{array}
 \begin{array}{c}
 q_\alpha \\
 \nu_\alpha
 \end{array}$$

$|V_{CKM}| \sim$

$|U_{MNS}| \sim$

- Are not predictable in the SM - profoundly unsatisfactory.

Masses and Mixings have Common Origin

- in Quark and Lepton mass matrices: $M_{\frac{2}{3}}, M_{-\frac{1}{3}}, M_\ell, M_\nu$
- Each is product of Higgs vev and Yukawa coupling matrix
- Work with Hermitian Squares.
 - eg. for neutrinos: $N = M_\nu M_\nu^\dagger$
 - for charged leptons: $L = M_\ell M_\ell^\dagger$ etc.
- Diagonalise $U_\nu N U_\nu^\dagger = D_\nu \rightarrow$ eigenvalues (ie. masses-squared)
- Diagonalisation different for L and $N \Rightarrow$ mixing, ie. $U_{MNS} \equiv U_\ell U_\nu^\dagger$.
NB. Analogous for quarks.

Jarlskog (ie. Weak Basis) Covariance

- However L and N are not observable:

$$L' = U_J L U_J^\dagger \text{ and } N' = U_J N U_J^\dagger$$

have same eigenvalues and same mixing matrix:

$$U'_{MNS} = U_\ell (U_J^\dagger U_J) U_\nu^\dagger = U_{MNS}$$

- ie. Masses and mixing angles are “Jarlskog-invariant”, even though mass matrices transform.
- New basis is just as good a starting point...

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- But usually assume “special” basis, in which texture applies
- “Specialness” of this basis is not explained.

eg. H. Fritzsch:

$$M_{\frac{2}{3}} = \begin{pmatrix} 0 & a & 0 \\ a^* & 0 & b \\ 0 & b^* & c \end{pmatrix}$$

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- Before and since, many models enforce “texture” zeroes, hierarchical mass matrices, other textures.
- But usually assume “special” basis, in which texture applies
- “Specialness” of this basis is not explained.
- Since Einstein, have been alert to excessive reliance on a given basis.
- Jarlskog reiterates (hep-ph/0606050): “you should be able to formulate it in an invariant form”
- We adopt Jarlskog’s prescription as a principle!

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The Jarlskogian and Plaquette Invariance

Jarlskog's CP -violating invariant:

$$J = \text{Im}(U_{\alpha i} U_{\alpha j}^* U_{\beta i}^* U_{\beta j})$$

Fascinating properties:

- Parameterises CP violation
- Does not depend on which plaquette is used
- May be simply related to the lepton mass matrices:

$$\begin{pmatrix} U_{e1} & U_{e2}^* & U_{e3} \\ U_{\mu 1}^* & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix}$$

$$J = -i \frac{\text{Det}[L, N]}{2L_{\Delta} N_{\Delta}}$$

L_{Δ}, N_{Δ} are simple polynomials in m_{ℓ} and m_{ν} .

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- Clearly, *any* function of the $U_{\alpha i}$, symmetrised over all *flavour* labels, and reduced to a function of only elements of a single plaquette is plaquette-invariant.
- Working with observables, find that, like J , FS functions of the mixing matrix can always be expressed as simple functions of the mass matrices.
- Will introduce an elemental set
 - can be used for the flavour-symmetric description of *any* mixing scheme.

The $S3_\ell \times S3_\nu$ Flavour Permutation Group

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- Introduce the observable P matrix:

$$P = \begin{pmatrix} |U_{e1}|^2 & |U_{e2}|^2 & |U_{e3}|^2 \\ |U_{\mu1}|^2 & |U_{\mu2}|^2 & |U_{\mu3}|^2 \\ |U_{\tau1}|^2 & |U_{\tau2}|^2 & |U_{\tau3}|^2 \end{pmatrix}$$

- Parameterises mixing (up to the sign of J).

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- Parameterises mixing (up to the sign of J).
- Transforms as a (reducible) 3×3 (natural representation) of FPG.
- Rows and columns each sum to unity (a magic square)
 \Rightarrow completely specified by elements of *any* P -plaquette.
- Each P -plaquette transforms as (irreducible) 2×2 of FPG.

Singlets Under FPG?

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 - Find simple polynomials of elements of P
 - (anti-)symmetrise over flavour labels
- Simple representation theory →
 - 1st order in P : \exists no non-trivial singlets
 - 2nd order: one each of $\bar{\mathbf{1}} \times \bar{\mathbf{1}}$ and $\mathbf{1} \times \mathbf{1}$
 - 3rd order: one each of all four singlets
 - \geq 4th order: multiple instances of each
- Will stay at \leq 3rd order. Clearly four are sufficient.

Elemental Set of FS Observables

Define themselves, up to normalisation (and “offset” in 1×1 case):

$$\mathcal{G} = \frac{1}{2} \left[\sum_{\alpha i} (P_{\alpha i})^2 - 1 \right] \quad \mathcal{F} = \text{Det} P$$

$$\mathcal{C} = \frac{3}{2} \left[\sum_{\alpha i} (P_{\alpha i})^3 - \sum_{\alpha i} (P_{\alpha i})^2 \right] + 1 \quad \mathcal{A} = \frac{1}{18} \sum_{\gamma k} (L_{\gamma k})^3$$

where $L_{\gamma k} = (P_{\alpha i} + P_{\beta j} - P_{\beta i} - P_{\alpha j})$.

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where $L_{\gamma k} = (P_{\alpha i} + P_{\beta j} - P_{\beta i} - P_{\alpha j})$.

- \mathcal{F} and \mathcal{A} need no offset (they are anti-symmetric).
 - Reach extremum for no mixing
 - 0 for trimaximal mixing
- All normalised to maximum value = 1 (no mixing).
- \mathcal{G} and \mathcal{C} offset to zero for maximal mixing



$$P = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



$$P = \begin{pmatrix} \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \end{pmatrix}$$

Properties and Values

Observable Name	Order in P	Symmetry: $S3_\ell \times S3_\nu$	Theor. Range	Exptl. Range for Leptons	Exptl. Range for Quarks
\mathcal{F}	2	$\bar{\mathbf{1}} \times \bar{\mathbf{1}}$	$(-1, 1)$	$(-0.14, 0.12)$	$(0.893, 0.896)$
\mathcal{G}	2	$\mathbf{1} \times \mathbf{1}$	$(0, 1)$	$(0.15, 0.23)$	$(0.898, 0.901)$
\mathcal{A}	3	$\bar{\mathbf{1}} \times \bar{\mathbf{1}}$	$(-1, 1)$	$(-0.065, 0.052)$	$(0.848, 0.852)$
\mathcal{C}	3	$\mathbf{1} \times \mathbf{1}$	$(-\frac{1}{27}, 1)$	$(-0.005, 0.057)$	$(0.848, 0.852)$

Properties and values of FS observables. Experimentally allowed ranges estimated (90% CL) from compilations of current experimental results (neglect any correlations between the input quantities).

FSOs in Terms of Mass Matrices

Define reduced (ie. traceless) powers of mass matrices: $\widetilde{L}^m := L^m - \frac{1}{3}\text{Tr}(L^m)$
 (similarly for \widetilde{N}^m).

Now define Jarlskog-invariant:

$$\widetilde{T}_{mn} := \text{Tr}(\widetilde{L}^m \widetilde{N}^n), \quad m, n = 1, 2.$$

\widetilde{T} Completely equivalent to P (for known lepton masses).

Find:

$$\mathcal{F} \equiv \text{Det } P = 3 \frac{\text{Det } \widetilde{T}}{L_\Delta N_\Delta}; \quad \left[\text{cf. } J = -i \frac{\text{Det}[L, N]}{2L_\Delta N_\Delta} \right]$$

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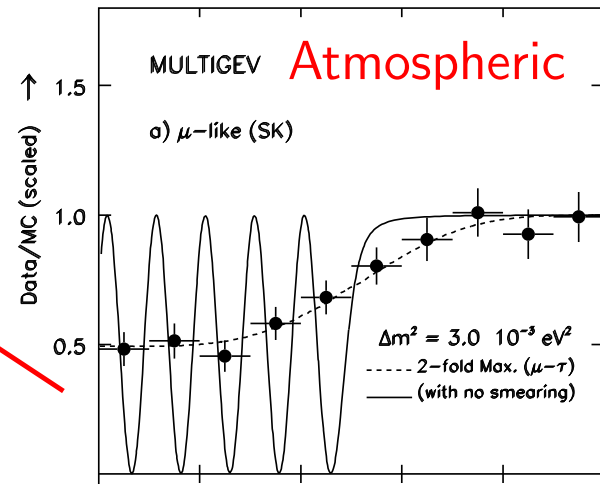
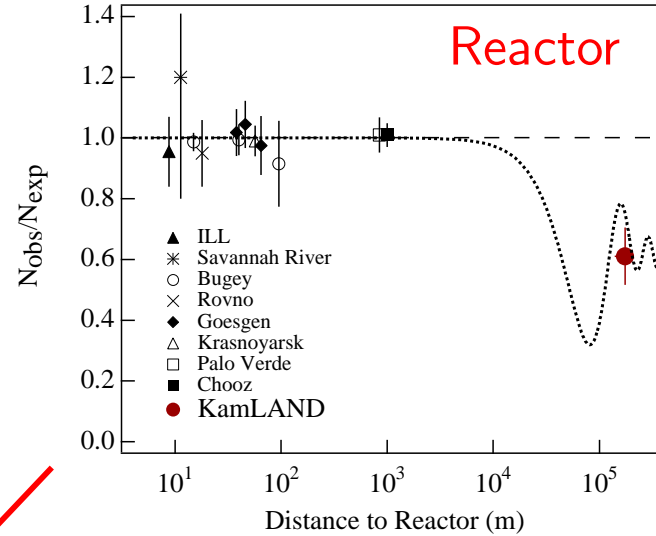
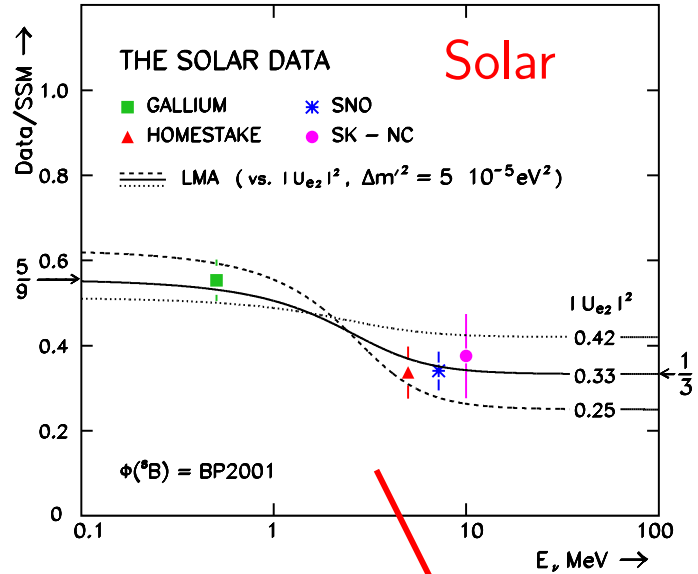
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$$\mathcal{G} = \frac{\widetilde{T}_{mn} \widetilde{T}_{pq} \mathcal{L}^{mp} \mathcal{N}^{nq}}{(L_{\Delta} N_{\Delta})^2}; \quad \mathcal{C}, \mathcal{A} = \frac{\widetilde{T}_{mn} \widetilde{T}_{pq} \widetilde{T}_{rs} \mathcal{L}_{\mathcal{C}, \mathcal{A}}^{(mpr)} \mathcal{N}_{\mathcal{C}, \mathcal{A}}^{(nqs)}}{(L_{\Delta} N_{\Delta})^{K_{\mathcal{C}, \mathcal{A}}}}$$

The \mathcal{L} (\mathcal{N}) are simple functions of traces of \widetilde{L}^m (\widetilde{N}^m). $K_{\mathcal{C}} (K_{\mathcal{A}}) = 2(3)$.

Application: Flavour-symmetric Descriptions of Mixing



$$|U| \sim \begin{matrix} e \\ \mu \\ \tau \end{matrix} \begin{pmatrix} 2/\sqrt{6} & 1/\sqrt{3} & \epsilon \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

$\Rightarrow \mathcal{F} = 0, \quad \mathcal{C} = 0, \quad \mathcal{A} = 0, \quad \mathcal{G} = \frac{1}{6}(1 - 3\epsilon^2)^2.$ Where is flavour-symmetry?

A Further Application

- What feature do quark and lepton mixing matrices have in common?
- Each has at least one “small” element.
- What is the flavour-symmetric expression of this?

$$|V_{CKM}| \sim \begin{pmatrix} \sim 1 & \sim \lambda & \sim \lambda^3 \\ \sim \lambda & \sim 1 & \sim \lambda^2 \\ \sim \lambda^3 & \sim \lambda^2 & \sim 1 \end{pmatrix}$$

$$|U_{MNS}| \sim \begin{pmatrix} 2/\sqrt{6} & 1/\sqrt{3} & \epsilon \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

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≥ 1 zero \Rightarrow CP-conservation $\implies J = 0$.

But need two constraints - find:

$$2\mathcal{A} + \mathcal{F}(\mathcal{F}^2 - 2\mathcal{C} - 1) = 0 \text{ and } J = 0 \implies \text{at least one zero}$$

A Further Application (Cont.)

- NB. Consider K -matrix: $K_{\gamma k} = \text{Re}(U_{\alpha i} U_{\alpha j}^* U_{\beta i}^* U_{\beta j})$ (cf. J).
- Noticed that $2\mathcal{A} + \mathcal{F}(\mathcal{F}^2 - 2\mathcal{C} - 1) \equiv \text{Det}K$
- Both MNS and CKM mixing matrices satisfy $\text{Det}K = 0$ (within errors)
- Conjecture: MNS *and* CKM constrained according to the same FS JI condition:

$$\text{Det}K = 0; \quad J \text{ small.}$$

- \implies FS prediction: if $\text{Det}K = 0$, then as $J \rightarrow 0$, at least one UT angle $\rightarrow 90^\circ$.

For quarks: $\text{Det}K = 0 \implies (90^\circ - \alpha) = \bar{\eta}\lambda^2 = 1^\circ \pm 0.2^\circ$

cf. $(90^\circ - \alpha) = 0^{+3^\circ}_{-7^\circ}$ experimentally.

For leptons: $\text{Det}K = 0 \implies (90^\circ - \delta) = \frac{|U_{e3}|}{\sqrt{2}} \lesssim 8^\circ$ (good for CPV).

Summary

- Models of masses and mixings should be weak-basis invariant
- Have defined flavour-symmetric mass and/or mixing observables.
- “Simplest” set defines itself up to normalisation
- Remarkably, leptonic mixing is consistent with 3 of these = 0!
- Can use them to implement FS and JI constraints on flavour observables, and make testable conjectures