



### Exotic meson spectroscopy at LHCb

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### www2.warwick.ac.uk

### Introduction



- In the quark model we think of hadrons as  $q\overline{q}$  or qqq
- But there is nothing preventing other combinations

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1 February 1964

### A SCHEMATIC MODEL OF BARYONS AND MESONS \*

M.GELL-MANN California Institute of Technology, Pasadena, California

We then refer to the members  $u^{\frac{2}{3}}$ ,  $d^{-\frac{1}{3}}$ , and  $s^{-\frac{1}{3}}$  of the triplet as "quarks" 6) q and the members of the anti-triplet as anti-quarks  $\bar{q}$ . Baryons can now be constructed from quarks by using the combinations (qqq),  $(qqqq\bar{q})$ , etc., while mesons are made out of  $(q\bar{q})$ ,  $(qq\bar{q}\bar{q})$ , etc. It is assuming that the lowest

 $\rightarrow$  Where are all those combinations with more than 3 quarks or anti-quarks?

### Introduction - molecules

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• There are lot of objects composed of baryons

### Introduction - molecules

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- There are lot of objects composed of baryons
- Where are similar objects from mesons?

### Scalar mesons



### NOTE ON SCALAR MESONS BELOW 2 GEV

Revised September 2013 by C. Amsler (Univ of Bern), S. Eidelman (Budker Institute of Nuclear Physics, Novosibirsk), T. Gutsche (University of Tübingen), C. Hanhart (Forschungszentrum Jülich), S. Spanier (University of Tennessee), and N.A. Törnqvist (University of Helsinki).

V. Interpretation of the scalars below 1 GeV: In the literature, many suggestions are discussed, such as conventional  $q\bar{q}$  mesons,  $q\bar{q}q\bar{q}$  or meson-meson bound states. In addition one expects a scalar glueball in this mass range. In reality, there can be superpositions of these components, and one often depends on models to determine the dominant one. Although we have seen progress in recent years, this question remains open. Here, we mention some of the present conclusions.

The f<sub>0</sub>(980) and a<sub>0</sub>(980) are often interpreted as multiquark states [140–144] or KK bound states [145]. The insight into
Candidates beyond qq mesons exist, but real trouble is how to decide whether they are qq or something else

### Charmonium

- Back in 2003, many expected states still missing
- Belle started to search for them and quickly found one
- Did not fit into expected spectra

















# What is X(3872)?

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- Lot has been done for X(3872)
- Despite all effort during 10 years, our understanding of what X(3872) is still about same
- To find convincing case of non-conventional meson is hard

 $Z(4430)^+$  history

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- Seen by Belle, but not Babar
- Data consistent
- Charged state
- $\rightarrow\,{\rm Cannot}\,\,{\rm be}\,\,c\overline{c}$ 
  - Latest Belle result uses 4D analysis
  - Is it real and if yes, is it resonance?



# $\mathbf{Z}(\mathbf{4430})^+$ history

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## **Issue of reflections**



- Look to bit simpler system of  $D^0 \rightarrow K_S \pi^+ \pi^-$  (only 2D rather than 4D system)
- Inspecting  $\pi^+\pi^-$  invariant mass, is there state around 1.8 GeV?



# **Issue of reflections**



- No new  $\pi^+\pi^-$  resonance
- What is seen in  $\pi^+\pi^-$  is result of reflection from  $K_S\pi$  and its angular structure
- Need to be careful when making claims in this type of systems



PRD 86, 032007 (2012

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# LHCb detector



# LHCb detector

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10h

12h 14h

02h

00h

### Data sample

- Use  $B^0 \to \psi(2S) K \pi$  decays
- Large statistics (> 25k), about 10 times what B-factories had
- $\blacksquare$  Very clean signal, background  $4\,\%$  of events (about 8% at B-factories)
- Perform both model-independent analysis (BABAR) and amplitude fit (Belle)



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# Amplitude analysis







- Mass described by relativistic Breit-Wigner
- Angular part using helicity formalism
- Imposes model how invariant mass distribution should look like

# $|M|^{2} = \sum_{\Delta\lambda_{\mu}} \left| \sum_{\lambda_{\psi}} \sum_{k} A_{k,\lambda_{\psi}}(\Omega | m_{0k}, \Gamma_{0k}) + \sum_{\lambda_{\psi}^{Z}} A_{Z,\lambda_{\psi}^{Z}}(\Omega^{Z} | m_{0Z}, \Gamma_{0Z}) e^{i\Delta_{\mu}\alpha} \right|^{2}$

**Amplitude** analysis

$$A_{k,\lambda_{\psi}}(\Omega|m_{R},\Gamma_{R}) = F_{B}^{L_{B}} \left(\frac{p_{B}}{m_{B}}\right)^{L_{B}} R(m|m_{R},\Gamma_{R})F_{R}^{L_{R}} \left(\frac{p_{R}}{m_{R}}\right)^{L_{R}} Z(\Omega)$$
  
Blatt-Weisskopf form factor Angular distribution (Helicity)  
Orbital momentum part

$$R(m|m_R, \Gamma_R) = \frac{1}{m_R^2 - m^2 - im_R \Gamma(m, \Gamma_R)}$$
$$\Gamma(m, \Gamma_R) = \Gamma_R \left(\frac{p_R}{p_{R0}}\right)^{2L_R + 1} \frac{m_R}{m} F_R^2$$

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## Model independent method

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- Test whether contributions in  $K\pi$  system can describe data
- Do not impose specific model for resonances
- $\rightarrow$  Model independent test

- Try to build up model which has proper behaviour for  $K\pi$ resonances
- But avoid imposing assumptions on the shape of  $m(K\pi)$  for resonances
- Construct Dalitz plot for pure  $K\pi$  activity and project on  $\psi(2S)\pi$  axis
- See whether model and data agree

# Model independent method



• Look to  $\cos(\theta_K)$  in bins of  $K\pi$  mass

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 Allows to find out which spins contribute

$$\sum_{i} \frac{1}{\epsilon_i} P_l(\cos \theta_{Ki})$$

Take only moments corresponding to  $J \leq 2$ 

Construct Dalitz plot and project on  $\psi(2S)\pi$  axis

### Why Legendre moments?



$$\langle P_1^U \rangle = S_0 P_0 \cos(\delta_{S_0} - \delta_{P_0}) + 2\sqrt{\frac{2}{5}} P_0 D_0 \cos(\delta_{P_0} - \delta_{D_0}) + \sqrt{\frac{6}{5}} [P_{+1} D_{+1} \cos(\delta_{P_{+1}} - \delta_{D_{+1}}) + P_{-1} D_{-1} \cos(\delta_{P_{-1}} - \delta_{D_{-1}})]$$

$$\langle P_2^U \rangle = \sqrt{\frac{2}{5}} P_0^2 + \frac{\sqrt{10}}{7} D_0^2 + \sqrt{2} S_0 D_0 \cos(\delta_{S_0} - \delta_{D_0}) \\ - \left( \frac{1}{\sqrt{10}} \left( P_{+1}^2 + P_{-1}^2 \right) + \frac{5\sqrt{10}}{28} \left( D_{+1}^2 + D_{-1}^2 \right) \right)$$

$$\langle P_3^U \rangle = 3 \sqrt{\frac{6}{35}} P_0 D_0 \cos(\delta_{P_0} - \delta_{D_0}) - 3 \sqrt{\frac{2}{35}} (P_{+1} D_{+1} \cos(\delta_{P_{+1}} - \delta_{D_{+1}}) + P_{-1} D_{-1} \cos(\delta_{P_{-1}} - \delta_{D_{-1}}) )$$

## Why Legendre moments?



$$\langle P_1^U \rangle = S_0 P_0 \cos(\delta_{S_0} - \delta_{P_0}) + 2\sqrt{\frac{2}{5}} P_0 D_0 \cos(\delta_{P_0} - \delta_{D_0}) + \sqrt{\frac{6}{5}} [P_{+1} D_{+1} \cos(\delta_{P_{+1}} - \delta_{D_{+1}}) + P_{-1} D_{-1} \cos(\delta_{P_{-1}} - \delta_{D_{-1}})]$$

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- Allows to cut expansion in physically meaningful way
- If we cut expansion, we select maximal spin which can contribute
- You might wonder why this is important

### $\psi'\pi$ reflection





 $\begin{array}{c} \text{Stars} \\ \text{22} \\ \text{20} \\ 18 \\ 16 \\ 14 \\ 12 \\ 10 \\ 8 \\ 6 \\ 4 \\ 2 \\ 0 \end{array}$ 

- Example of  $B^0 \to Z^+ (\to \psi' \pi^+) K^-$
- Such contribution reflect to whole  $K\pi$  mass range
- Helicity angle distribution peaking
- $\rightarrow$  Moments will receive contributions from reflection



# Model independent result





Clearly, pure kaon resonances cannot explain M(\u03c6(2S)\u03c0) spectrum
 Understanding details difficult

- Resonances in  $\psi(2S)\pi$  will contribute to  $K\pi$  and its moments
- Any fit to  $\psi(2S)\pi$  on top of reflections neglects interference between two axes

## Only $K^{\ast}\xspace$ resonances

Candidates / ( $0.2 \text{ GeV}^2$ )	000	LH 16	Cb ++	<sup>+</sup> <sup>+</sup> <sup>+</sup> <sup>+</sup> <sup>+</sup> <sup>+</sup> <sup>+</sup> <sup>+</sup>	$0 \\ Candidates / (0.2 \text{ GeV}^2) \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ $	LHCb -1.0 < $m_{K^+\pi^-}^2$ < 1.8 G	$eV^2$ $eV^2$	PRL 112 (2014) 222002
F	Resonance	$J^P$	Likely n <sup>25+1</sup> LJ	Mass (MeV)	Width (MeV) $547 \pm 24$	$\mathcal{B}(K^{*0} \rightarrow K^+\pi^-)$	- data	
,	$K^{*}(892)^{0}$	$\binom{\kappa}{1}^{-}$	$1^{3}S_{1}$	$895.94 \pm 0.26$	$347 \pm 24$ 48.7 + 0.7	$\sim 100\%$ $\sim 100\%$	total fit	
	$K_0^* (1430)^{C}$	<b>0</b> <sup>-</sup>	$1^{3}P_{0}$	$1425 \pm 50$	$270 \pm 80$	$(93 \pm 10)\%$	—— K <sup>*</sup> (902)	
ŀ	$K_1^*(1410)^0$	) 1 <sup>-</sup>	$2^{3}S_{1}$	1414 $\pm$ 15	$232\pm21$	$(6.6 \pm 1.3)\%$	- K (892)	
	$K_{2}^{*}(1430)^{0}$	2+	$1^{3}P_{2}$	$1432.4\pm1.3$	$109\pm5$	$(49.9 \pm 1.2)\%$	── K <sup>*</sup> S-wa	ave
B	$B^0  o \psi(2, \phi)$	$S)K^+\pi$	phase space limit	1593				
ŀ	$K_1^*(1680)^{\circ}$		$1^{3}D_{1}$	$1717 \pm 27$	$322 \pm 110$	$(38.7 \pm 2.5)\%$	$ K_2(1430)$	J)
/	K <sub>3</sub> (1780) <sup>6</sup>	$3^{-}$	$1^{3}D_{3}$	$1776 \pm 7$	$159 \pm 21$	$(18.8 \pm 1.0)\%$	backgro	ound
I	$\kappa_0^{\pi} (1950)^{\circ}$	$0^+$	$2^{\circ}P_{0}$	$1945 \pm 22$	$201 \pm 78$	$(52 \pm 14)\%$		
	$x_4 (2045)^{\circ}$	4' K+		$2045 \pm 9$	$198 \pm 30$	$(9.9 \pm 1.2)\%$	<b>→</b> К (1680	))
	$\psi^* \rightarrow J/\psi$	$\kappa'\pi$	phase space limit	2183	170   20	(6.1   1.0)0/	<u> — к</u> *(1410	))
_	$(2380)^{\circ}$	5	1° G5	$2382 \pm 9$	$178 \pm 32$	$(0.1 \pm 1.2)\%$		·)

## Only $K^{\ast}$ resonances

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 $\textbf{Adding} \ \mathbf{Z}^+$ 



## **Dalitz plot slices**



### Results





- Data are described well with  $1^+ Z(4430)^+$  contribution ( $\chi^2$  p-value 12%)
- Parameters extracted consistent with Belle
- Large interference effects seen
- Adding additional  $K^*$  resonances to model does not alter conclusion

$$f_Z = \frac{\int A_Z(\Omega) d\Omega}{\int A(\Omega) d\Omega} \qquad f_Z^I = 1 - \frac{\int A_{\text{no}Z}(\Omega) d\Omega}{\int A(\Omega) d\Omega}$$

$\mathbf{Z}(4430)^+$ spin							THE UNIVERSITY OF	
PRL 112 (2014) 222002	0.300 250 50 50 -200	Simulated experiments $J_Z^P = 0^-$	Simulated experiments $J_Z^P = 1^+$ 200	HCb ata 400 $\Delta(-2 \ln L)$	_	$\begin{array}{c} \text{Hypothesis} \\ 0^- \\ 1^- \\ 2^+ \\ 2^- \end{array}$	$\begin{array}{c} {\sf Rejection} \\ 9.7\sigma \\ 15.8\sigma \\ 16.1\sigma \\ 14.6\sigma \end{array}$	·

- As we use full kinematic information, we have sensitivity to quantum numbers
- Test spins 0,1 and 2 with both parities
- Based on likelihood ratio
- Quote exclusion based on asymptotic formula (lower bound)
- Verified by simulation
- All rejections relative to  $1^+$
- $Z(4430)^+$  is  $1^+$  state without any doubts

# Is $Z(4430)^+$ resonance?



$$\frac{1}{m_R^2 - m^2 - im_R\Gamma(m,\Gamma_R)}$$

- Data are consistent with BW for  $Z(4430)^+$
- But will they follow if BW is not imposed?
- Change BW in  $Z(4430)^+$ amplitude to 6 complex numbers in 6  $M(\psi(2S)\pi)$  bins
- Plot resulting amplitude on Argand plot

# Is $Z(4430)^+$ resonance?

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- Data are consistent with BW for  $Z(4430)^+$
- But will they follow if BW is not imposed?
- Change BW in  $Z(4430)^+$ amplitude to 6 complex numbers in 6  $M(\psi(2S)\pi)$  bins
- Plot resulting amplitude on Argand plot
- ⇒ It shows resonance behaviour without imposing it

#### Second $Z^+$ state THE UNIVERSITY OF WAR PRL 112 (2014) 222002 Candidates / ( 0.2 GeV<sup>2</sup> LHCb $4239 \pm 18^{+45}_{-10} \; {\rm MeV}$ $M(Z_0)$ 200 $-1.0 < m_{K^+\pi^-}^2 < 1.8 \text{ GeV}^2$ $220 \pm 47^{+1\bar{0}\bar{8}}_{-74} ~{\rm MeV}$ $\Gamma(Z_0)$ $1.6 \pm 0.5^{+1.9}_{-0.4}$ % $f_{Z_0}$ $2.4 \pm 1.1^{+1.7}_{-0.2}$ % $f_{Z_0}^I$ 100 Significance $6\sigma$ $22 m_{W'\pi^-}^2 [GeV^2]$ 18 20 16

- $\blacksquare$  Data can be described even better by adding second  $\psi(2S)\pi$  state
- On its own, it is significant
- Preferred  $0^-$  (but  $660 \pm 150$  MeV wide  $1^+$  option cannot be ruled out)
- Argand diagram is inconclusive
- No evidence in model-independent approach
- Will need more data to clarify situation

Second  $Z^+$  state





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### **Excitement?**



### Interpretation



- What  $Z(4430)^+$  really is?
- Large decay width  $\Rightarrow$  strong decay
- $c\overline{c}$  state in final state  $\Rightarrow c\overline{c}$  has to be also in initial state
- Charged, so cannot be conventional charmonia with  $c\overline{c}$  only
- From Argand plot it behaves as resonance
- Often there are threshold effects (cusps) when new channels opens
  - In case of  $Z(4430)^+$  we are close to  $\overline{D}^* D_1$  threshold
  - Cusp would be S-wave effect, so would have  $J^P = 0^-$ ,  $1^-$  or  $2^-$
  - We find  $J^P = 1^+$  thus excluding threshold effect
- From all this all conventional explanations fail
- We have something "exotic" like tetraquark, molecule, ...
- If there is one "exotic" state, then there should be whole spectrum of them, so lot of work ahead of us
- First natural choice to look for is neutral partner of  $Z(4430)^+$

## Conclusions



- At LHCb we collected large samples of  ${\cal B}$  decays
- Started to check various claims for new states
  - Often those analyses are difficult as we want to do careful job
- $\blacksquare$  We confirmed existence of  $Z(4430)^+$ 
  - Belle and Babar had different conclusion due to lower statistics and lower sensitivity of method at Babar
- We improved measurement of  $Z(4430)^+$  properties
- Our date show proper resonance behaviour of  $Z(4430)^+$
- We exclude non-exotic interpretation of  $Z(4430)^+$
- Exotic spectroscopy is now fully open for new states