

Quark flavour physics

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- Plan
 - Kaon physics and SM construction (bit of history)
 - Establishing SM experimentally
 - Looking for breakdown of SM
- Hard to cover everything in details in three lectures, some details are offloaded to exercises

Outline – lecture 3

- Why we need new physics
- Where to look for new physics
 - CPV in $B_s \rightarrow J/\psi\phi$
 - $b \rightarrow s l^+ l^-$
 - $B_s \rightarrow \mu\mu$
 - $B^+ \rightarrow \tau\nu$
 - Charm mixing and CPV
- Useful theory oriented reference: A.J.Buras, arXiv:1012.1447

Why new physics?

- What we saw up to now is extreme success of SM
- So why we are so excited about new physics if we can describe all measurements without it?
- SM does not tell why we have three generations (do we?)
- SM does not tell anything about masses of fermions
- Way SM works needs quite some fine tuning (final result is often difference between two huge numbers)
- Higgs mass in SM is not really stable
- Cosmology argument:
 - No candidate for Dark matter
 - Not possible to generate current matter-antimatter asymmetry in Universe

- In 1967 Andrei Sakharov formulated three necessary conditions to produce baryon-antibaryon asymmetry
 - Baryon number violation
 - CP violation
 - Interactions out of equilibrium
- Surprisingly SM can best cope with baryon number violation – but non-perturbative QCD effect
- CP violation exists in SM, but definitely not large enough
- In SM there is nothing which could drive system out of equilibrium
- Those are in my view strongest needs for new physics

Typical NP scenarios

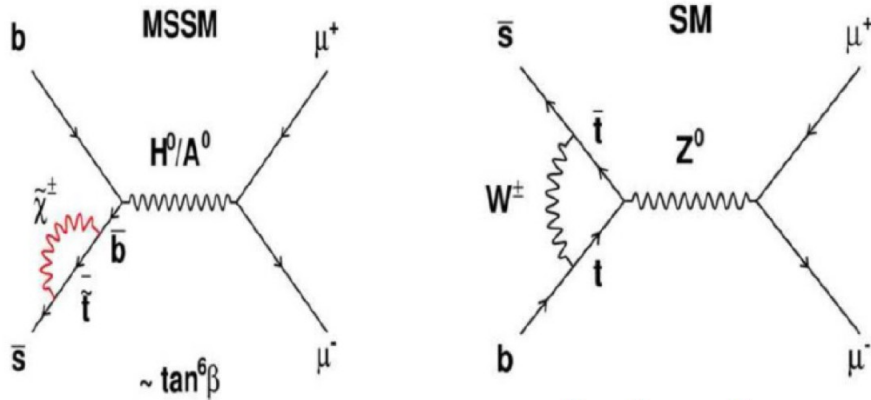
- SUSY – most popular
 - Each SM particle gets its SUSY partner
 - Squarks can also mix (similar to quarks) and provide additional phases
- Fourth generation
 - Puts in additional fermion generation
 - Not as appealing as SUSY and cannot answer all questions, but can make huge difference on CPV
- Technicolor
 - Gets more popular again as it can dynamically create Higgs like effect without need of Higgs

	AC	RVV2	AKM	δLL	FBMSSM	LHT	RS
$D^0 - \bar{D}^0$	★★★★	★	★	★	★	★★★★	?
ϵ_K	★	★★★★	★★★★	★	★	★★	★★★★
$S_{\psi\phi}$	★★★★	★★★★	★★★★	★	★	★★★★	★★★★
$S_{\phi K_S}$	★★★★	★★	★	★★★★	★★★★	★	?
$A_{CP}(B \rightarrow X_s \gamma)$	★	★	★	★★★★	★★★★	★	?
$A_{7,8}(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★★★★	★★★★	★★	?
$A_9(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★	★	★	?
$B \rightarrow K^{(*)} \nu \bar{\nu}$	★	★	★	★	★	★	★
$B_s \rightarrow \mu^+ \mu^-$	★★★★	★★★★	★★★★	★★★★	★★★★	★	★
$K^+ \rightarrow \pi^+ \nu \bar{\nu}$	★	★	★	★	★	★★★★	★★★★
$K_L \rightarrow \pi^0 \nu \bar{\nu}$	★	★	★	★	★	★★★★	★★★★
$\mu \rightarrow e \gamma$	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★	★★★★
d_n	★★★★	★★★★	★★★★	★★	★★★★	★	★★★★
d_e	★★★★	★★★★	★★	★	★★★★	★	★★★★
$(g-2)_\mu$	★★★★	★★★★	★★	★★★★	★★★★	★	★★

Table 2: “DNA” of flavour physics effects [55] for the most interesting observables in a selection of SUSY and non-SUSY models. ★★★★★ signals large effects, ★★ visible but small effects and ★ implies that the given model does not predict sizable effects in that observable.

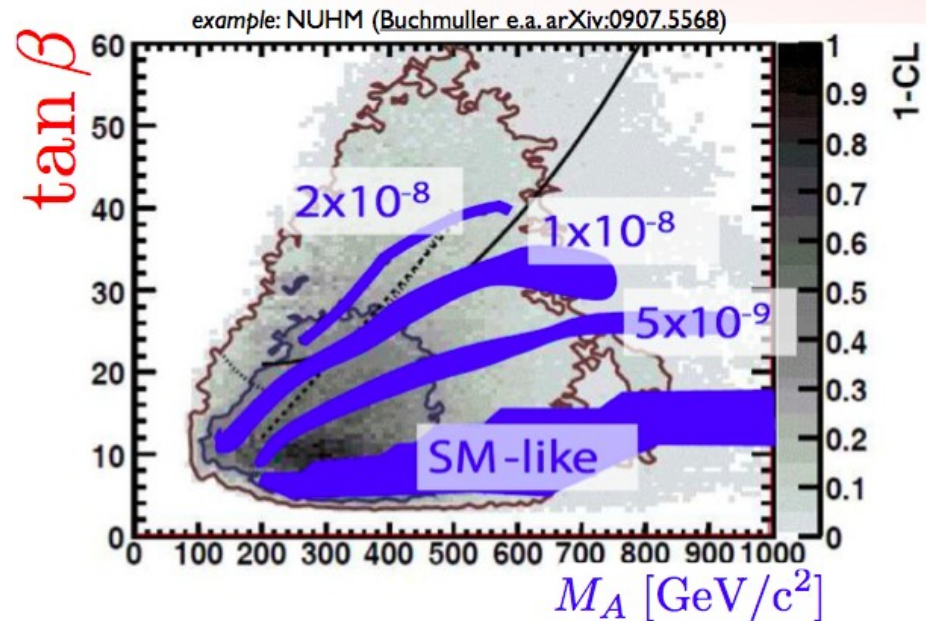
Buras, arXiv:0910.1032v1

$B_s \rightarrow \mu^+ \mu^-$



$$Br^{MSSM}(Bq \rightarrow l^+ l^-) \propto \frac{m_b^2 m_l^2 \tan^6 \beta}{M_{A0}^4}$$

- SM prediction $Br = (3.2 \pm 0.2) 10^{-9}$
- Extremely sensitive to new physics
- Already with limits 10 times above SM one could severely restrict NP models



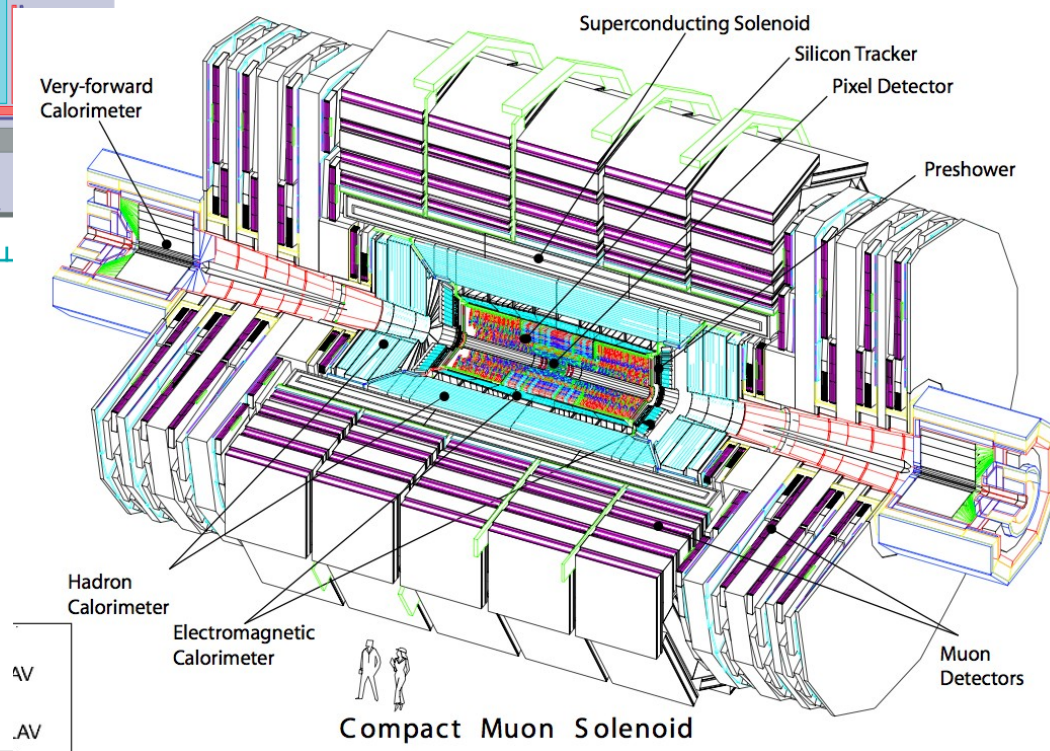
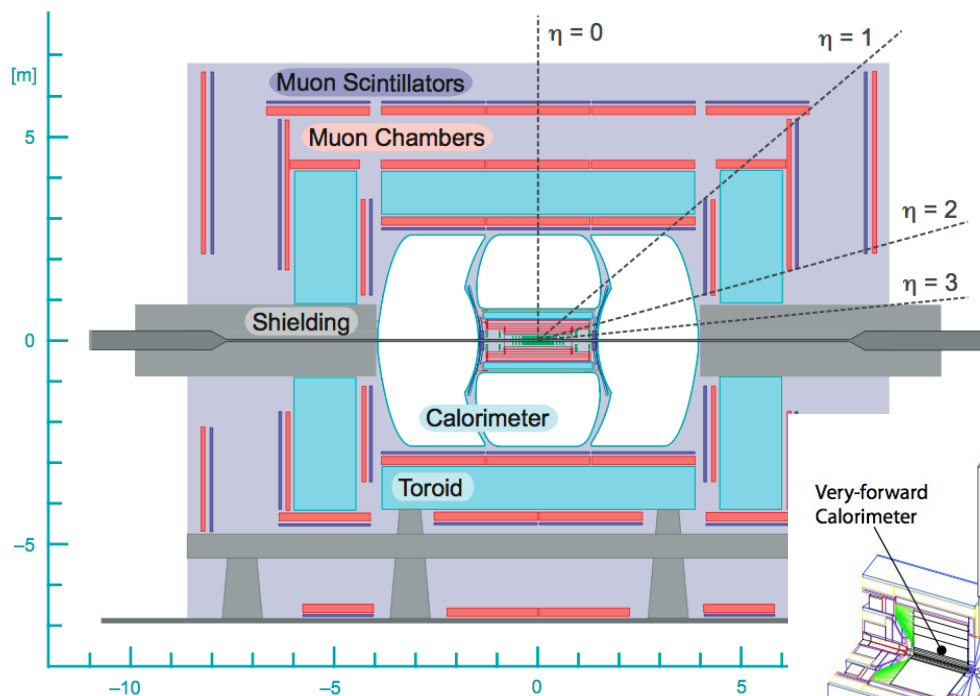
Background plot is from
Buchmuller et al arXiv:0907.5568

Blue regions are allowed regions for given
measurement.

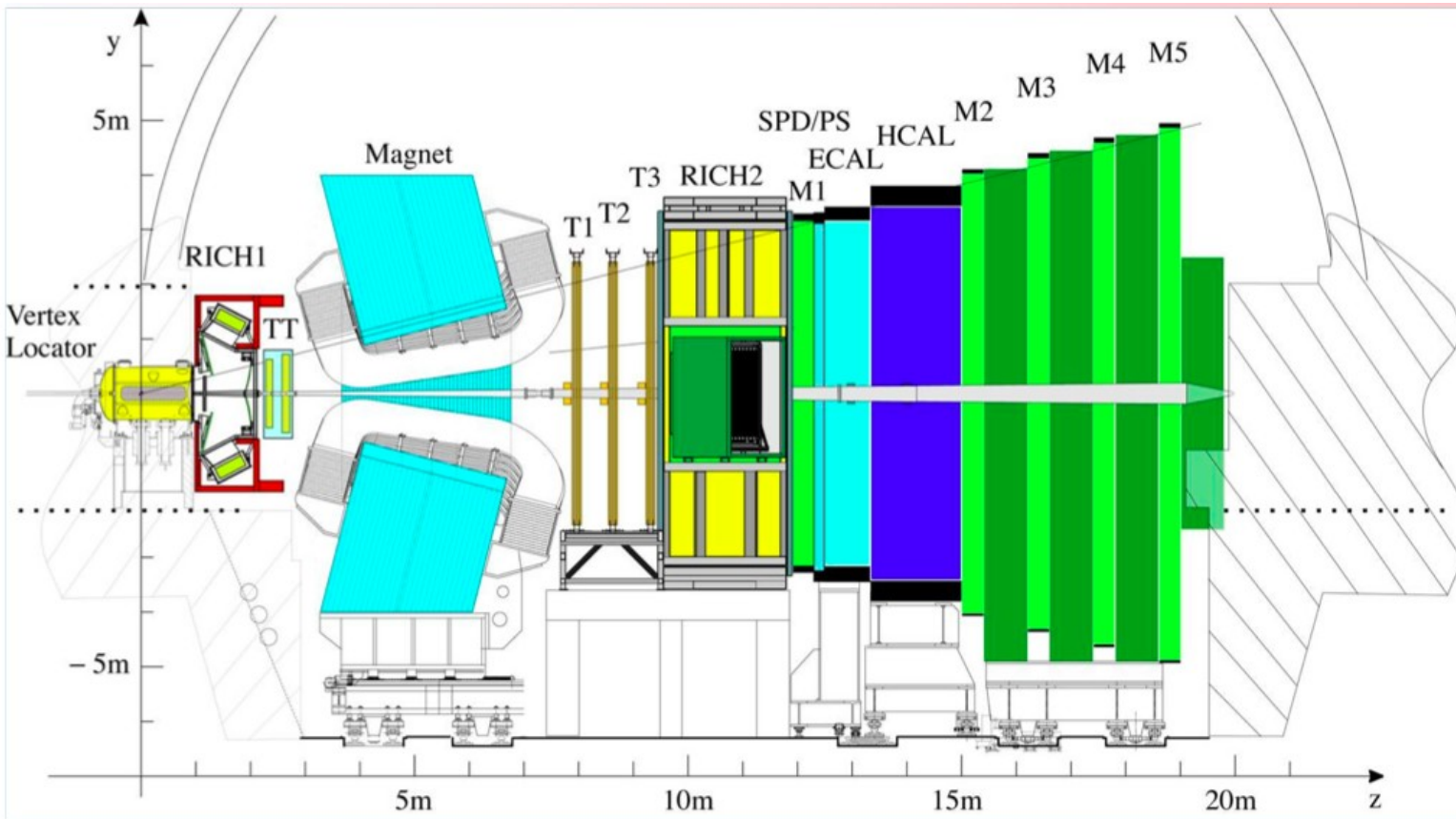
SuperIso program, (F. Mahmoudi, arXiv: 08083144)
and SoftSusy (B.C. Allanach, Comput. Phys. Commun.
143 (2002) 305-331)

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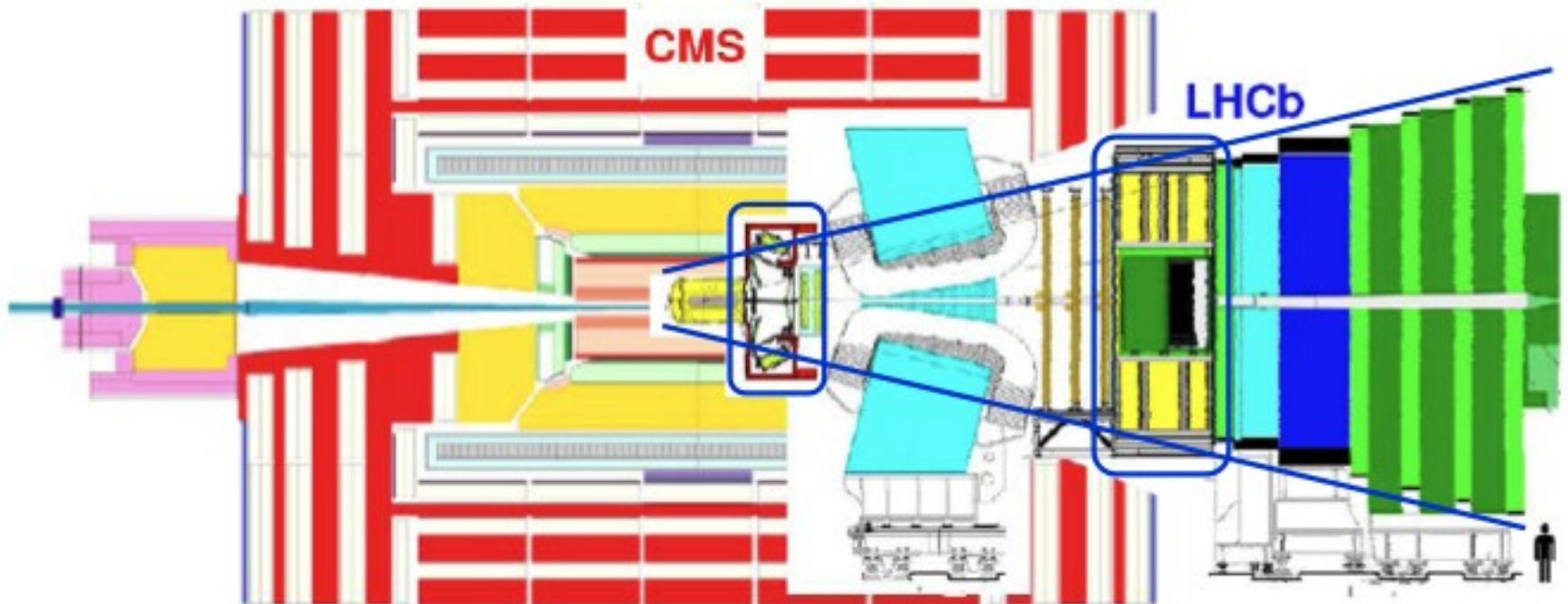
Experiments



Experiments

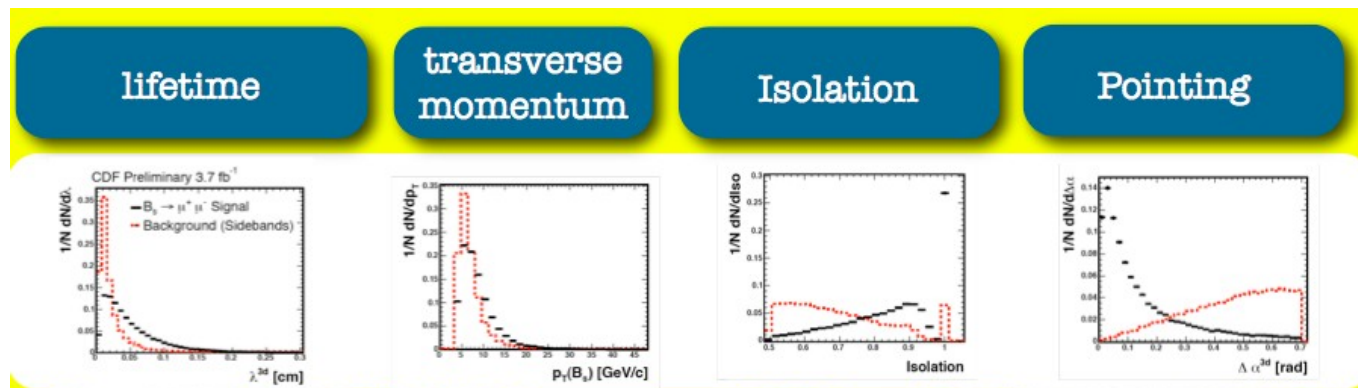


LHCb is not small one

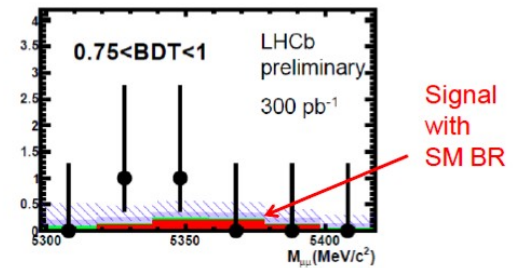
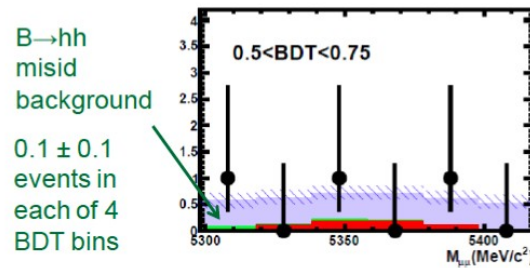
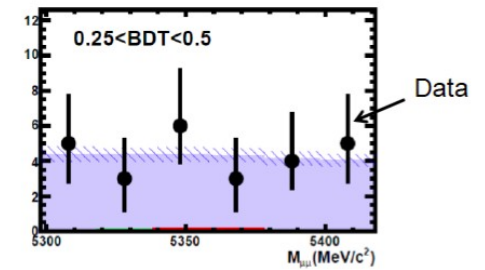
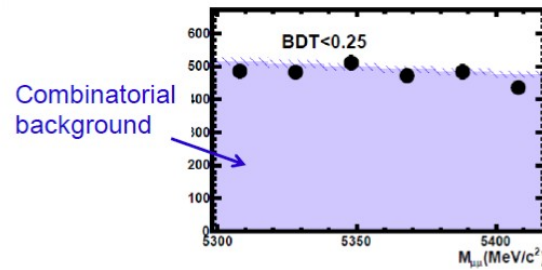
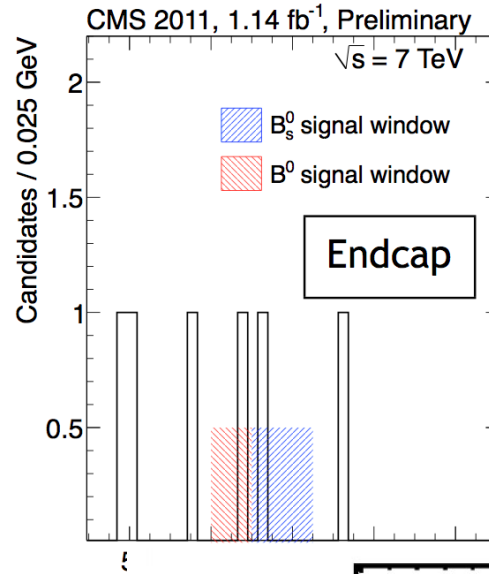
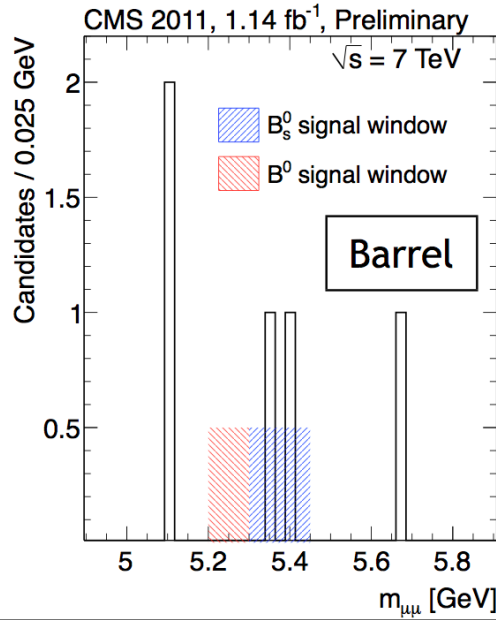


Main experimental issues

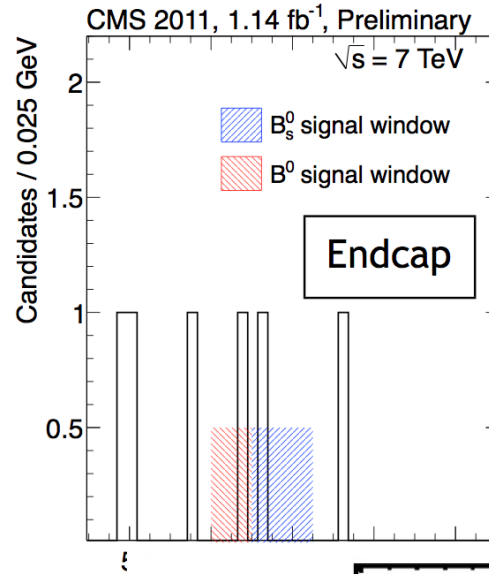
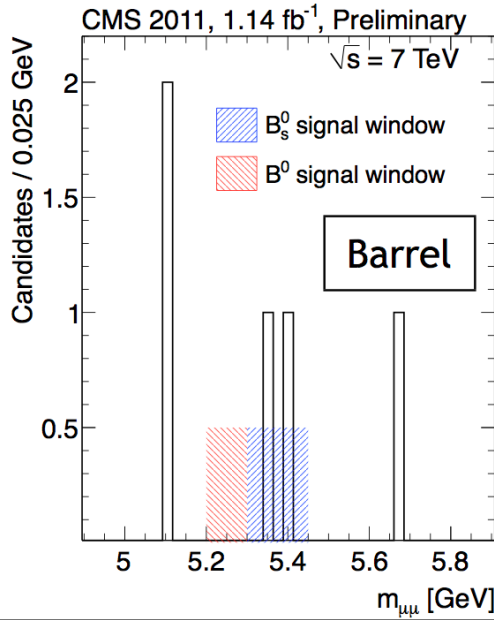
- As decay is rare, main issue is to control background
- Special care is needed for $B \rightarrow hh'$ decays which peak in same place as signal
- Selection ranges from cut based (CMS) to BDT (LHCb) to NN(CDF, D0)



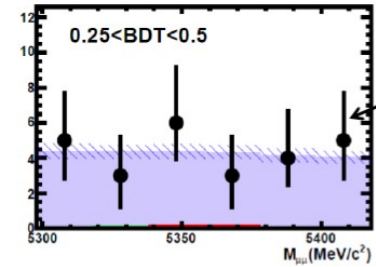
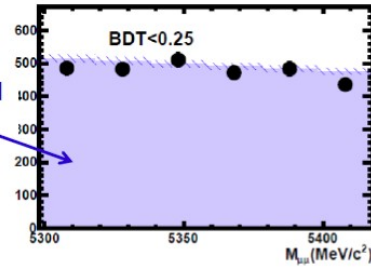
$B_s \rightarrow \mu^+ \mu^-$



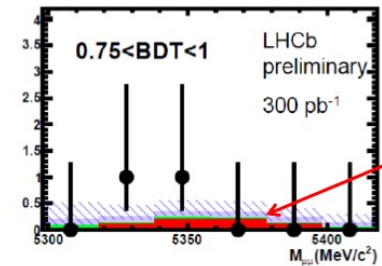
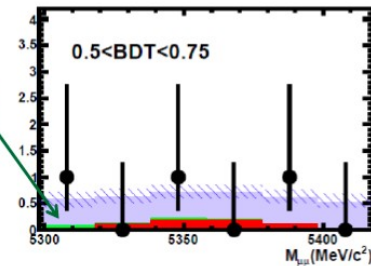
$B_s \rightarrow \mu^+ \mu^-$



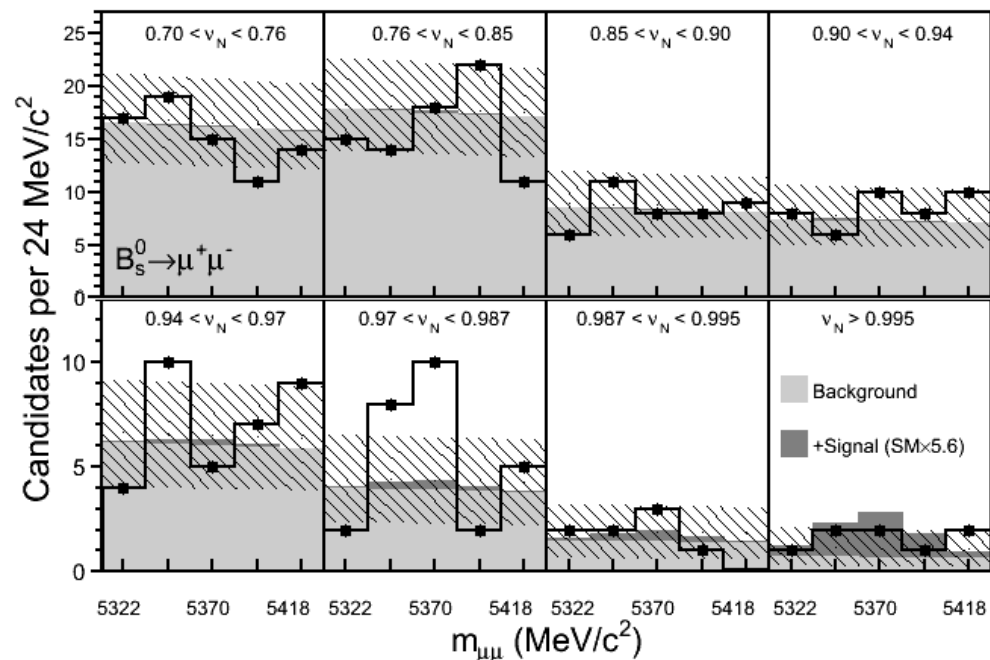
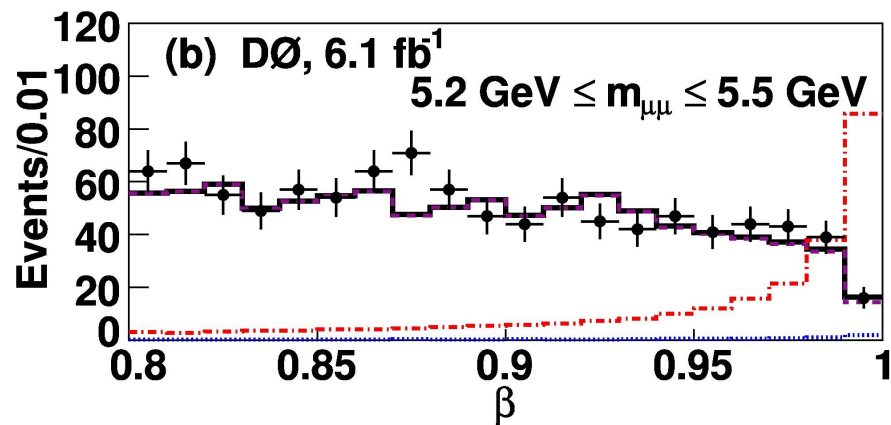
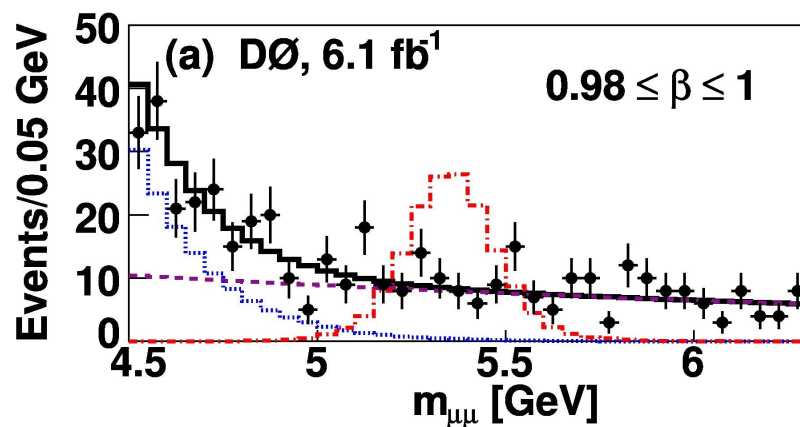
Combinatorial background



$B \rightarrow hh$ misid background
 0.1 ± 0.1 events in each of 4 BDT bins



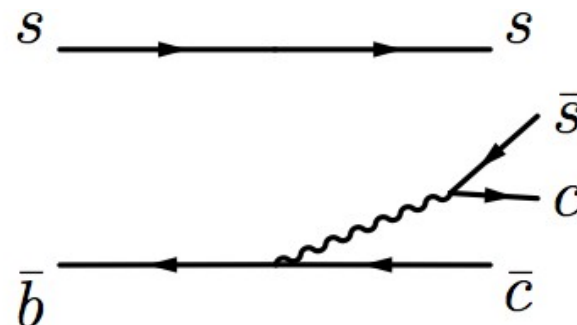
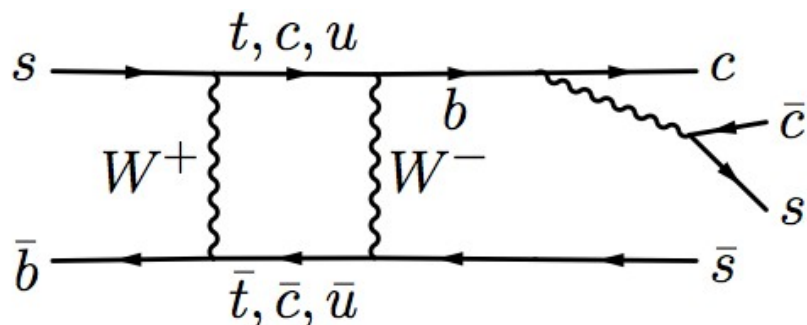
$$B_S \rightarrow \mu^+ \mu^-$$



$B_s \rightarrow \mu^+ \mu^-$

- D0, LHCb and CMS compatible with pure background (plus tiny SM signal)
- CDF sees excess of events above background
- Results:
 - $B(B_s^0 \rightarrow \mu^+ \mu^-) < 5.1 \times 10^{-8}$ @ 95% C.L. (D0)
 - $B(B_s^0 \rightarrow \mu^+ \mu^-) < 1.9 \times 10^{-8}$ @ 95% C.L. (CMS)
 - $B(B_s^0 \rightarrow \mu^+ \mu^-) < 1.5 \times 10^{-8}$ @ 95% C.L. (LHCb)
 - $B(B_s^0 \rightarrow \mu^+ \mu^-) < 3.9 \times 10^{-8}$ @ 95% C.L. (CDF)
 - $B = (1.8^{+1.1}_{-0.9}) \times 10^{-8}$ (CDF interpreting excess as signal)
- About factor 5 above SM
- Expect to come close to SM at latest end of next year

CPV induced by B_s mixing



- CPV stems from interference of decays with and w/o mixing
- In SM it is expected to be tiny – Effect of almost real V_{ts}
- Needs some care in interpretation of what is really measured
- Up to now measured with $B_s \rightarrow J/\psi\phi$ and $B_s \rightarrow J/\psi f_0(980)$
- Difficult due to large $\Delta m_s \sim 17.8 \text{ ps}^{-1}$
- Needs very good time resolution

Master formulas of time evolution

$$\Gamma(M(t) \rightarrow f) = \mathcal{N}_f |A_f|^2 e^{-\Gamma t} \left\{ \frac{1 + |\lambda_f|^2}{2} \cosh \frac{\Delta\Gamma t}{2} + \frac{1 - |\lambda_f|^2}{2} \cos(\Delta M t) - \operatorname{Re} \lambda_f \sinh \frac{\Delta\Gamma t}{2} - \operatorname{Im} \lambda_f \sin(\Delta M t) \right\}, \quad (50)$$

$$\Gamma(\bar{M}(t) \rightarrow f) = \mathcal{N}_f |A_f|^2 \frac{1}{1-a} e^{-\Gamma t} \left\{ \frac{1 + |\lambda_f|^2}{2} \cosh \frac{\Delta\Gamma t}{2} - \frac{1 - |\lambda_f|^2}{2} \cos(\Delta M t) - \operatorname{Re} \lambda_f \sinh \frac{\Delta\Gamma t}{2} + \operatorname{Im} \lambda_f \sin(\Delta M t) \right\}. \quad (51)$$

$$\Gamma(M(t) \rightarrow \bar{f}) = \mathcal{N}_f |\bar{A}_{\bar{f}}|^2 e^{-\Gamma t} (1-a) \left\{ \frac{1 + |\lambda_{\bar{f}}|^{-2}}{2} \cosh \frac{\Delta\Gamma t}{2} - \frac{1 - |\lambda_{\bar{f}}|^{-2}}{2} \cos(\Delta M t) - \operatorname{Re} \frac{1}{\lambda_{\bar{f}}} \sinh \frac{\Delta\Gamma t}{2} + \operatorname{Im} \frac{1}{\lambda_{\bar{f}}} \sin(\Delta M t) \right\}, \quad (53)$$

$$\Gamma(\bar{M}(t) \rightarrow \bar{f}) = \mathcal{N}_f |\bar{A}_{\bar{f}}|^2 e^{-\Gamma t} \left\{ \frac{1 + |\lambda_{\bar{f}}|^{-2}}{2} \cosh \frac{\Delta\Gamma t}{2} + \frac{1 - |\lambda_{\bar{f}}|^{-2}}{2} \cos(\Delta M t) - \operatorname{Re} \frac{1}{\lambda_{\bar{f}}} \sinh \frac{\Delta\Gamma t}{2} - \operatorname{Im} \frac{1}{\lambda_{\bar{f}}} \sin(\Delta M t) \right\}. \quad (54)$$

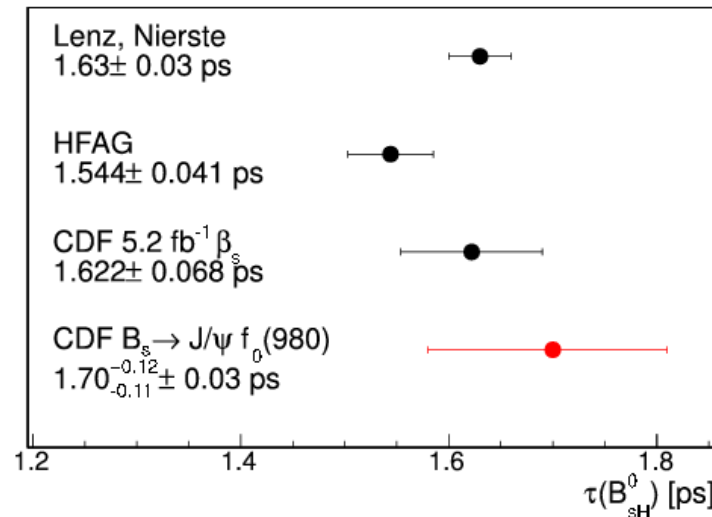
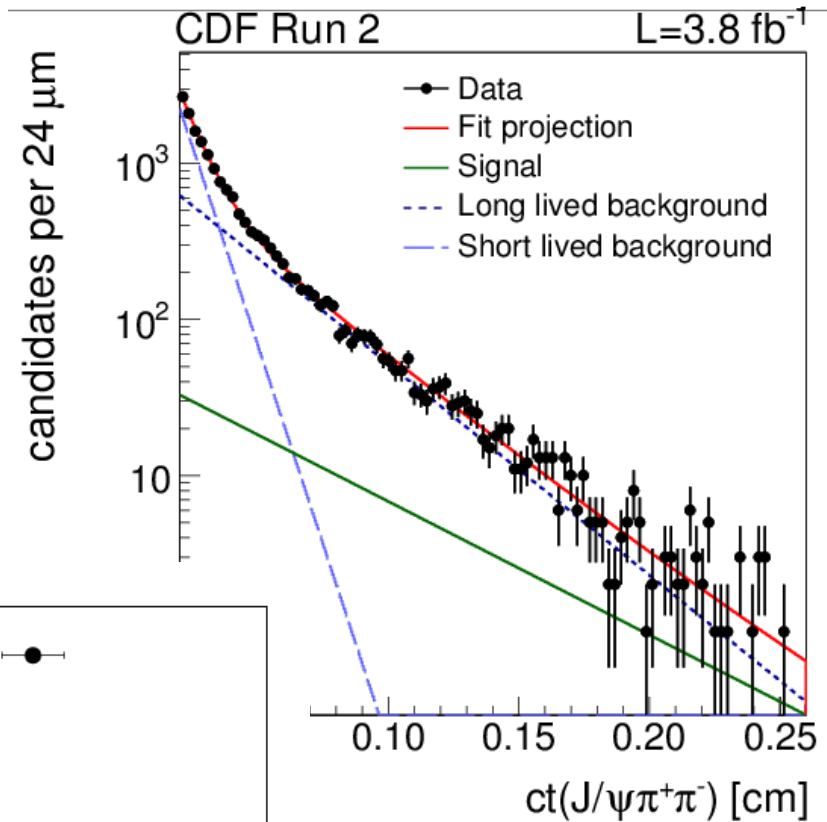
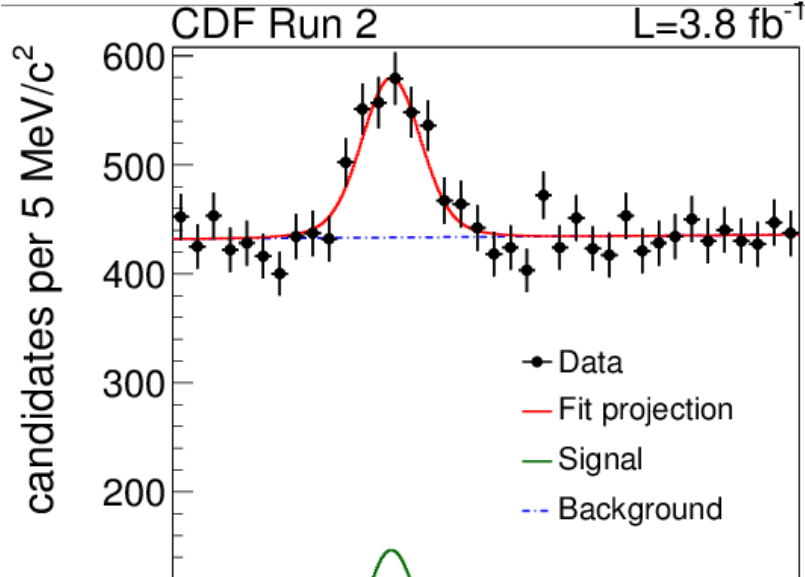
Decay $B_s \rightarrow J/\psi f_0$

- This final state is CP-eigenstate
- In master formula this (with some additional assumptions) means $|\lambda_f|=1$
- For B_s the a is at most $\sim 1\%$ and $\Delta\Gamma$ is in principle non-zero

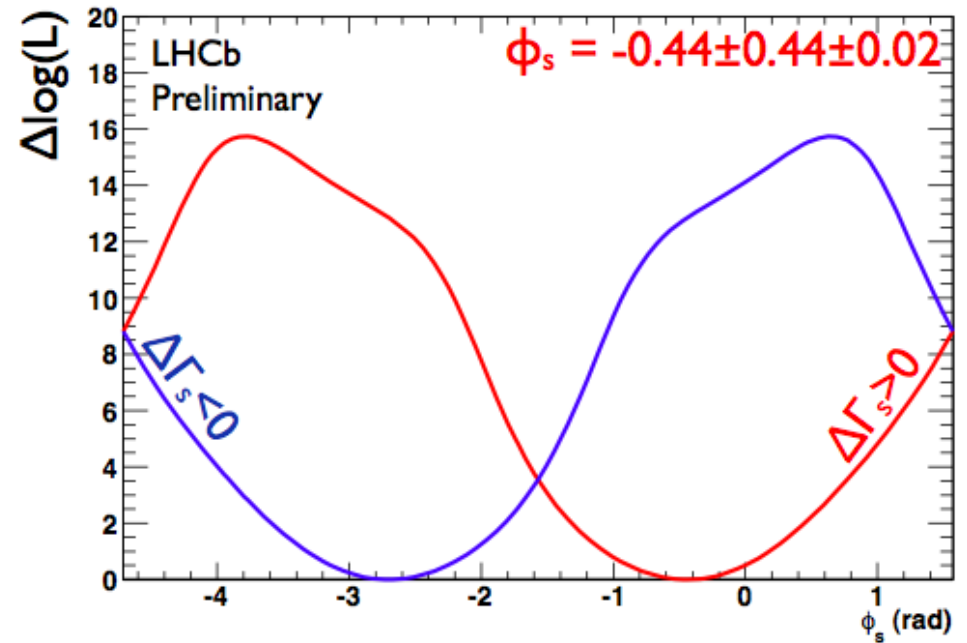
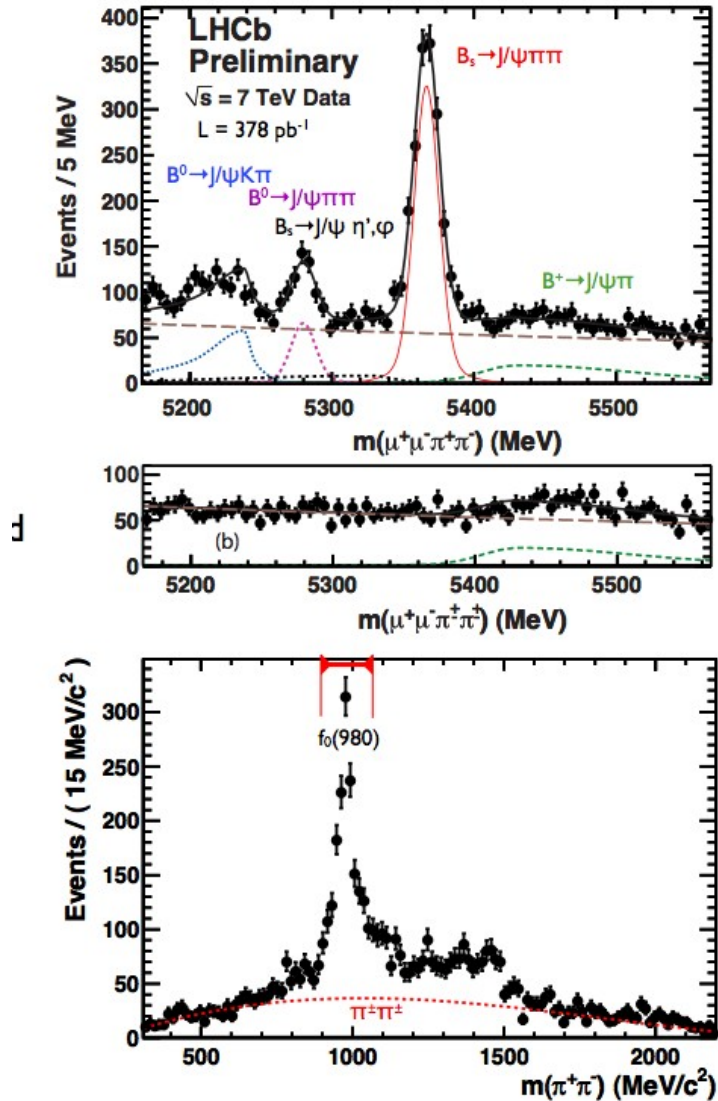
$$\frac{1 + \cos(\phi)}{2} e^{-\Gamma_H t} + \frac{1 - \cos(\phi)}{2} e^{-\Gamma_L t} - e^{-\Gamma t} \sin(\phi) \sin(\Delta m t)$$

$$\frac{1 + \cos(\phi)}{2} e^{-\Gamma_H t} + \frac{1 - \cos(\phi)}{2} e^{-\Gamma_L t}$$

Decay $B_s \rightarrow J/\psi f_0$ - CDF



Decay $B_s \rightarrow J/\psi f_0$ - LHCb

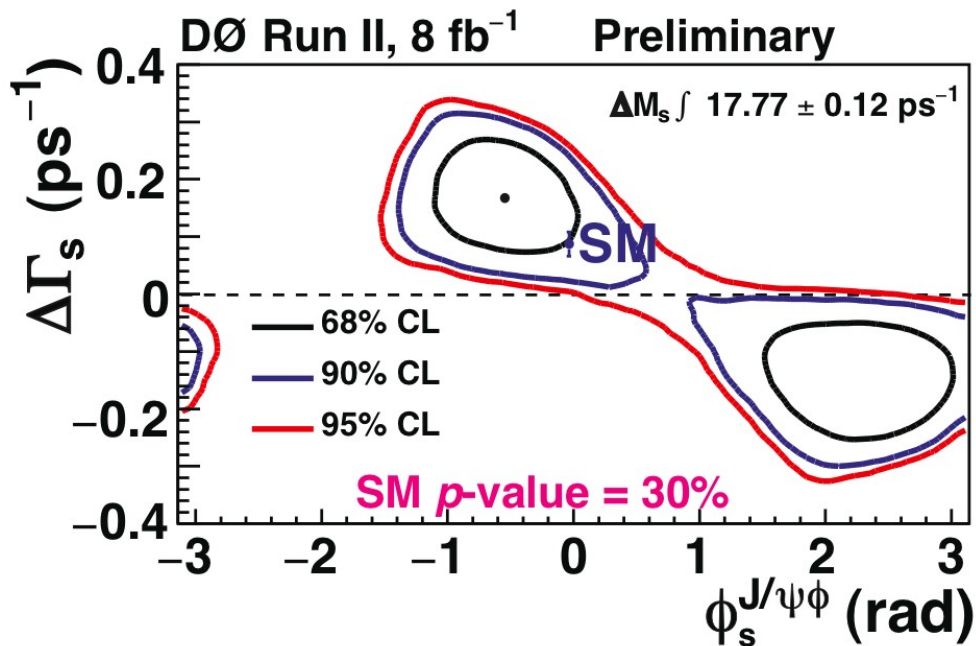
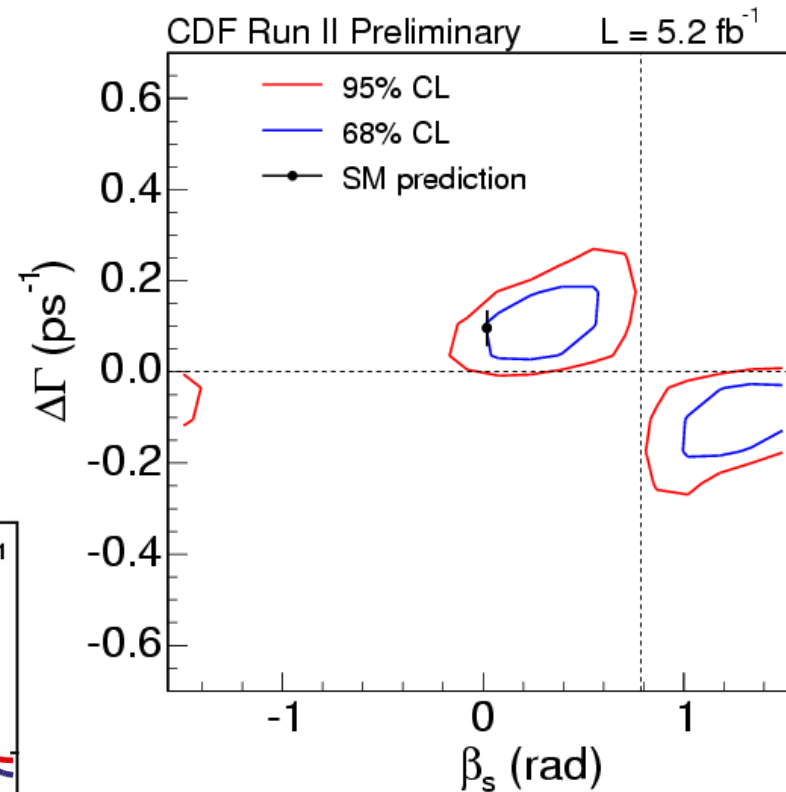
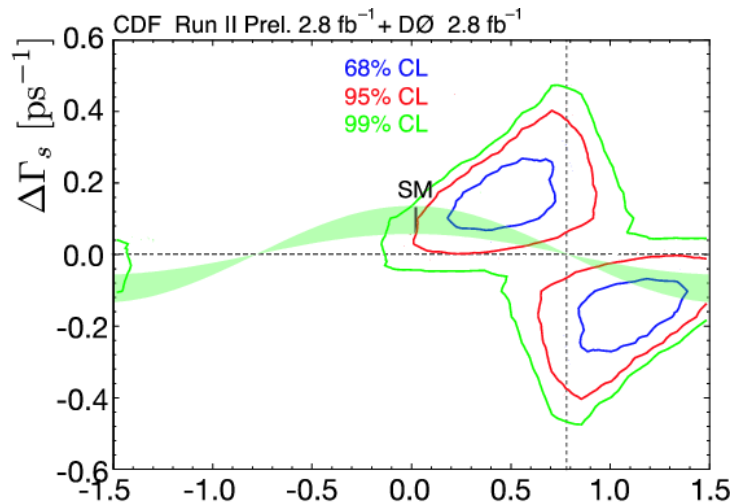


- Takes input of Γ and $\Delta\Gamma$ from LHCb $B_s \rightarrow J/\psi \phi$ result
- Consistent with SM

Decay $B_s \rightarrow J/\psi \phi$

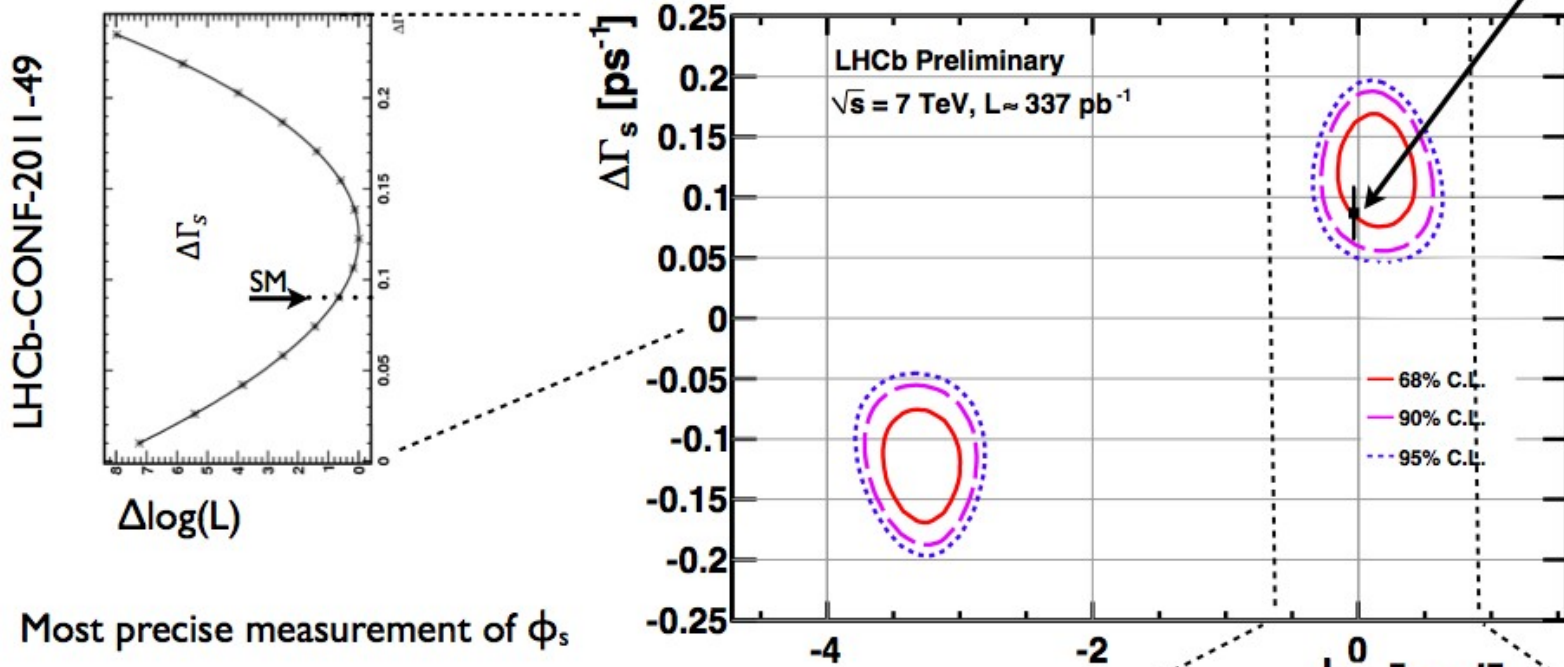
- This decay was first to study as it has higher signal yield
- Complication arises from the fact that both J/ψ and ϕ are have spin 1 => mixture of CP-even and CP-odd
- Need additional work (angular analysis) to statistically distinguish two
- On the other hand it also helps to gain additional sensitivity from interference
- There are many details to care about, much beyond scope of this lecture
- If there is interest, in some free time I can walk you through CDF analysis to full depth

Decay $B_s \rightarrow J/\psi\phi$



$B_s \rightarrow J/\psi \varphi$: $\Delta\Gamma_s$ vs. ϕ_s

Standard Model
(Lenz, Nierste: arXiv:1102.4274)

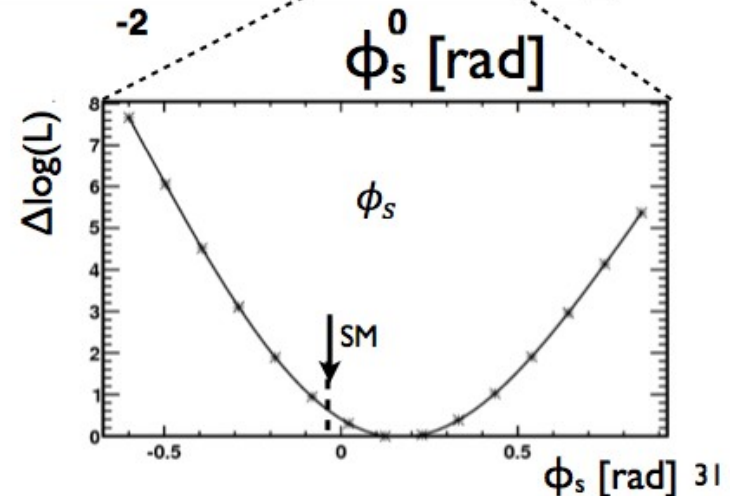


Most precise measurement of ϕ_s

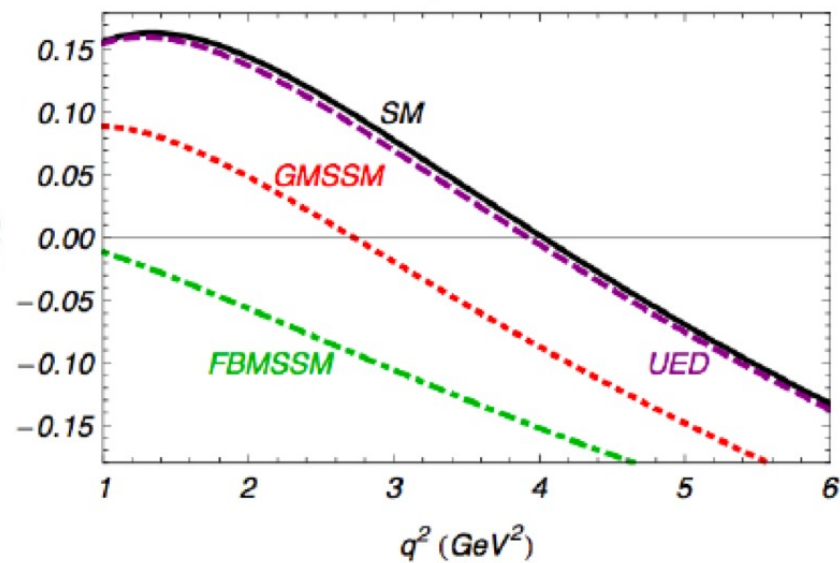
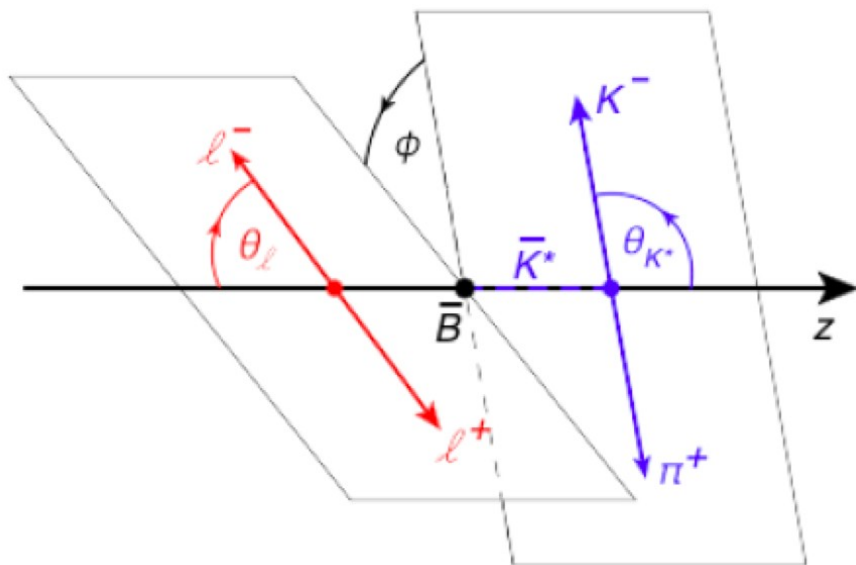
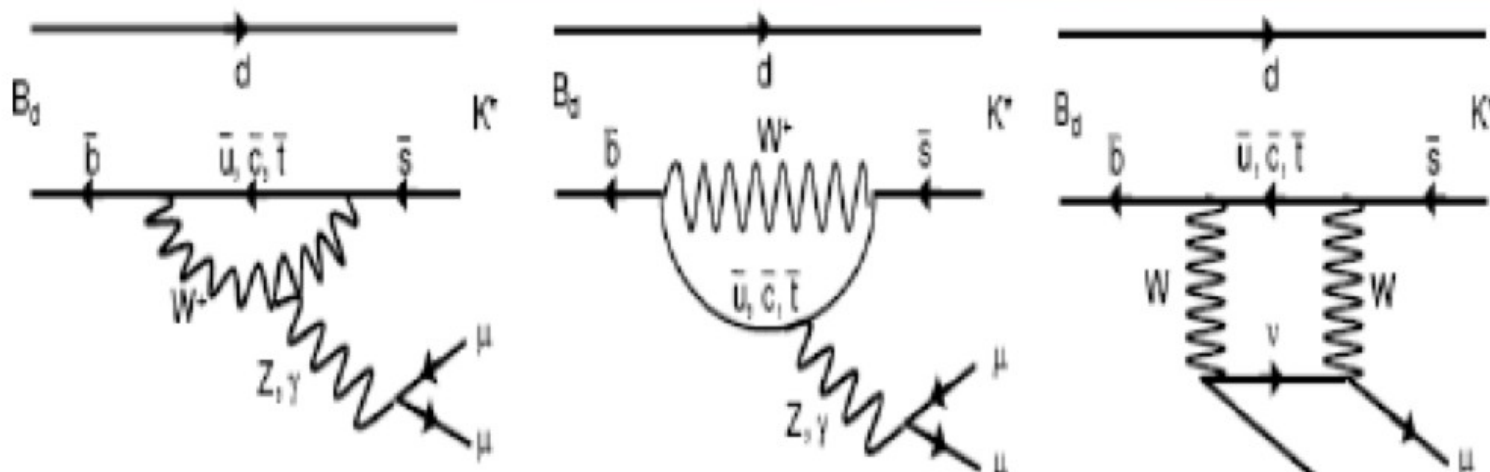
- $\phi_s = 0.13 \pm 0.18$ (stat) ± 0.07 (syst) rad
- Consistent with SM

4 σ Evidence for $\Delta\Gamma_s \neq 0$:

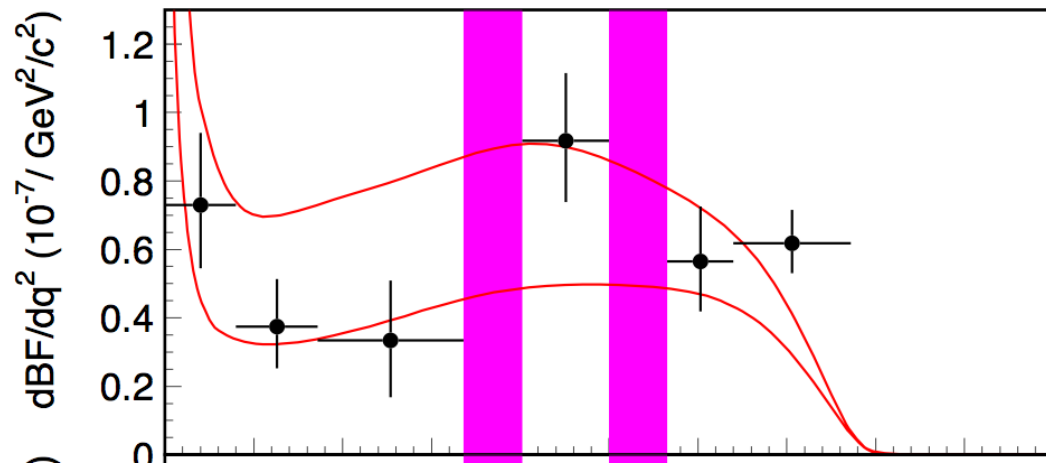
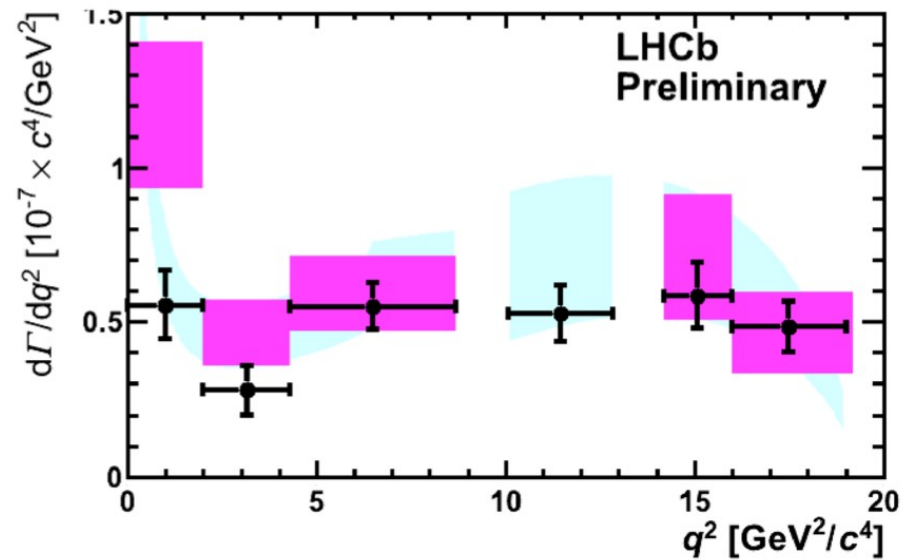
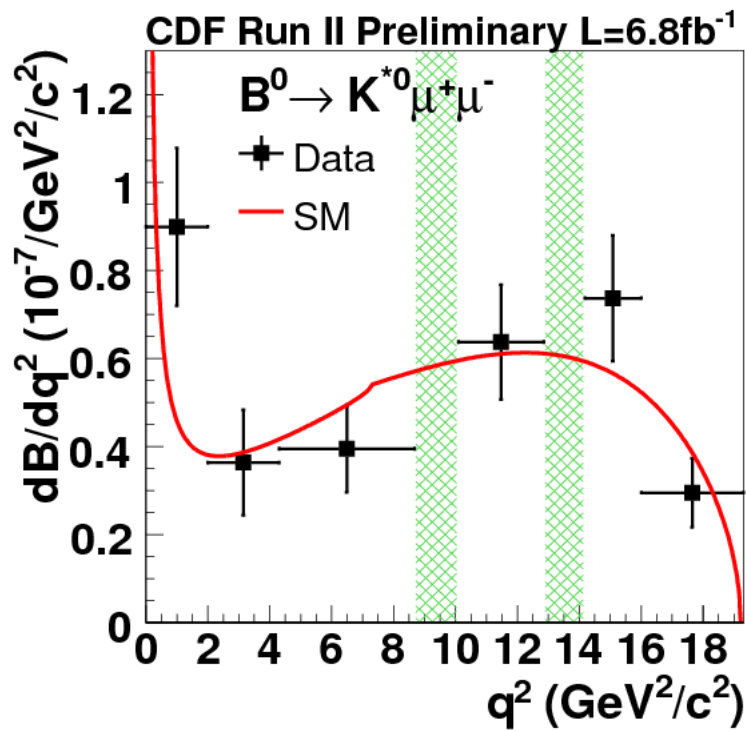
- $\Delta\Gamma_s = 0.123 \pm 0.029$ (stat) ± 0.008 (syst) ps^{-1}
- $\Gamma_s = 0.656 \pm 0.009$ (stat) ± 0.008 (syst) ps^{-1}



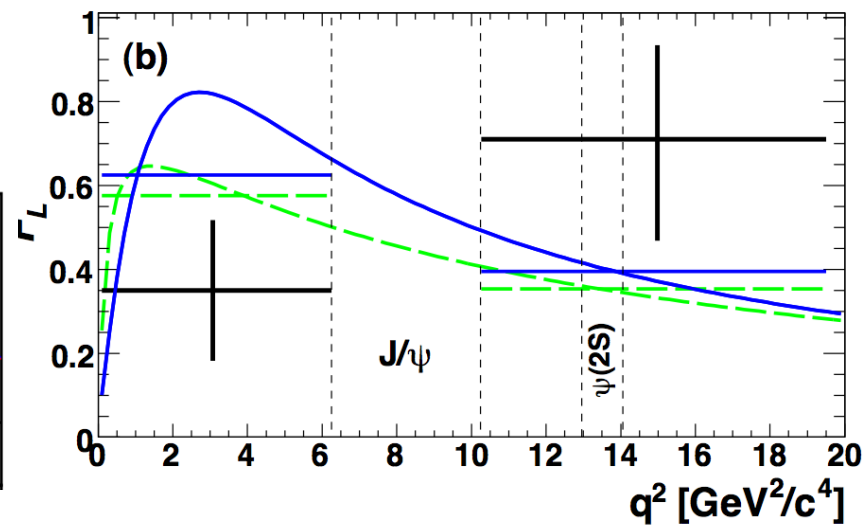
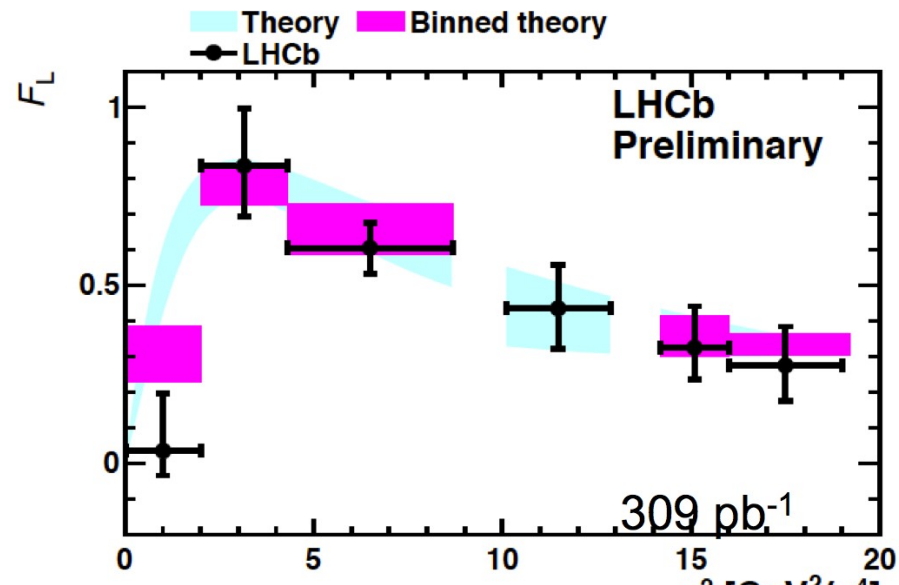
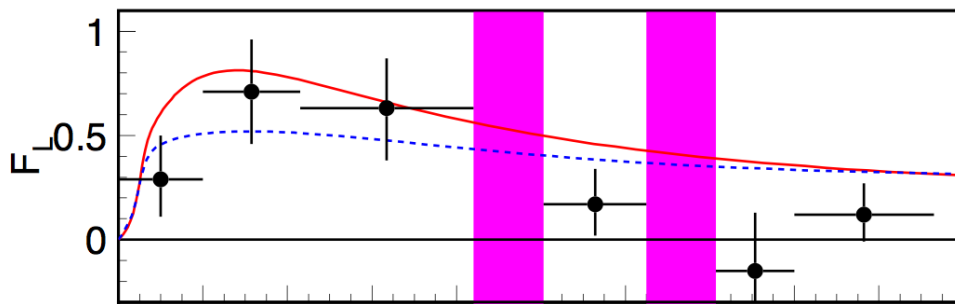
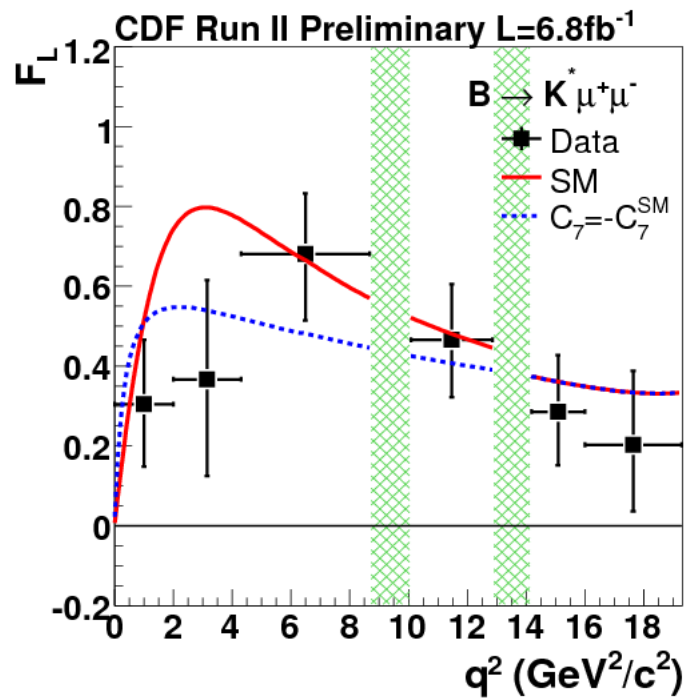
$b \rightarrow s l^+ l^-$



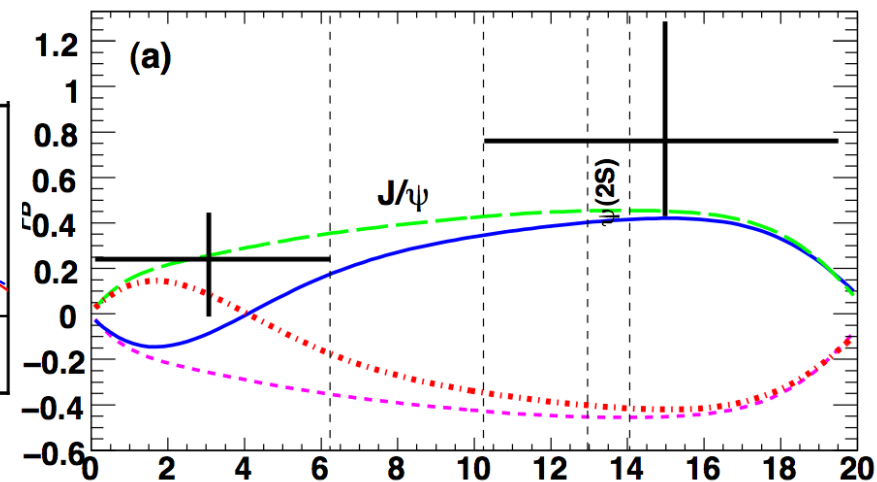
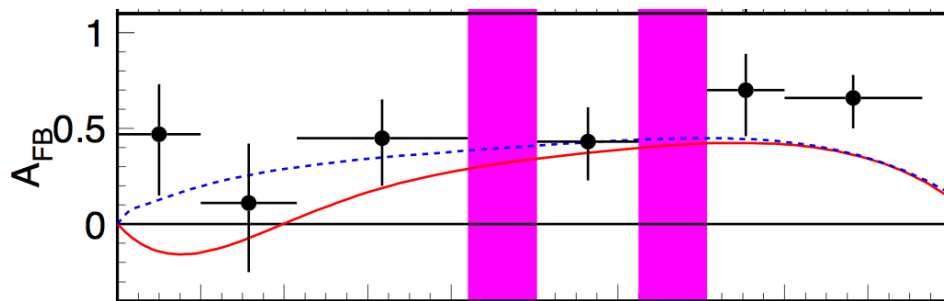
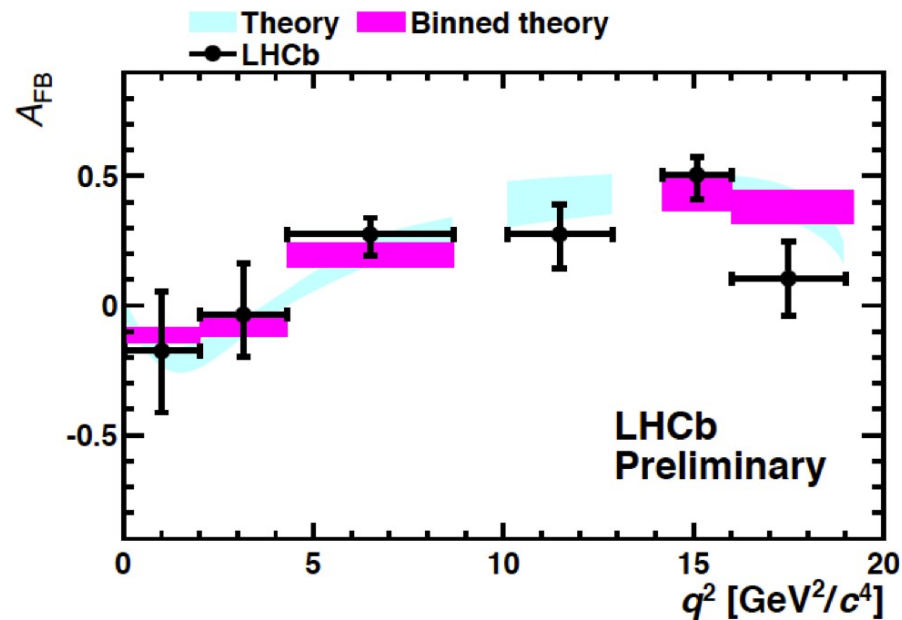
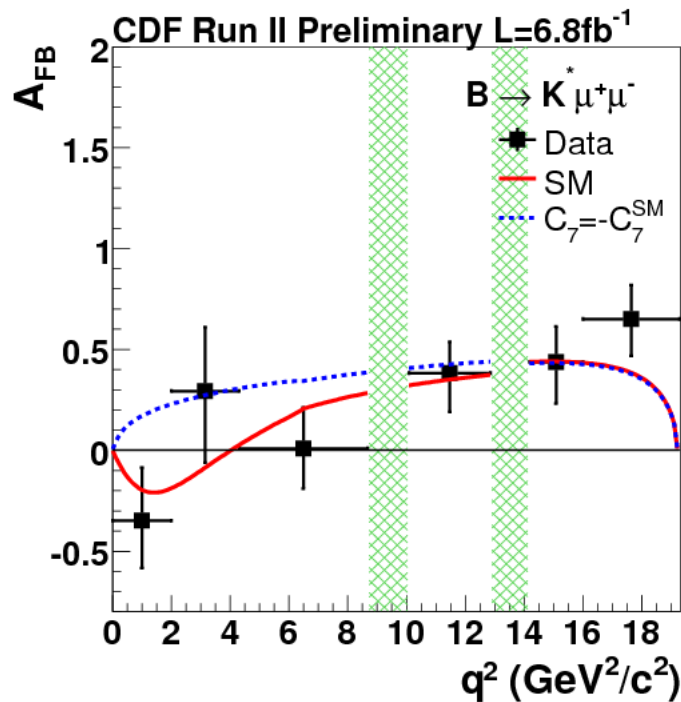
$b \rightarrow s |^+ |^-$



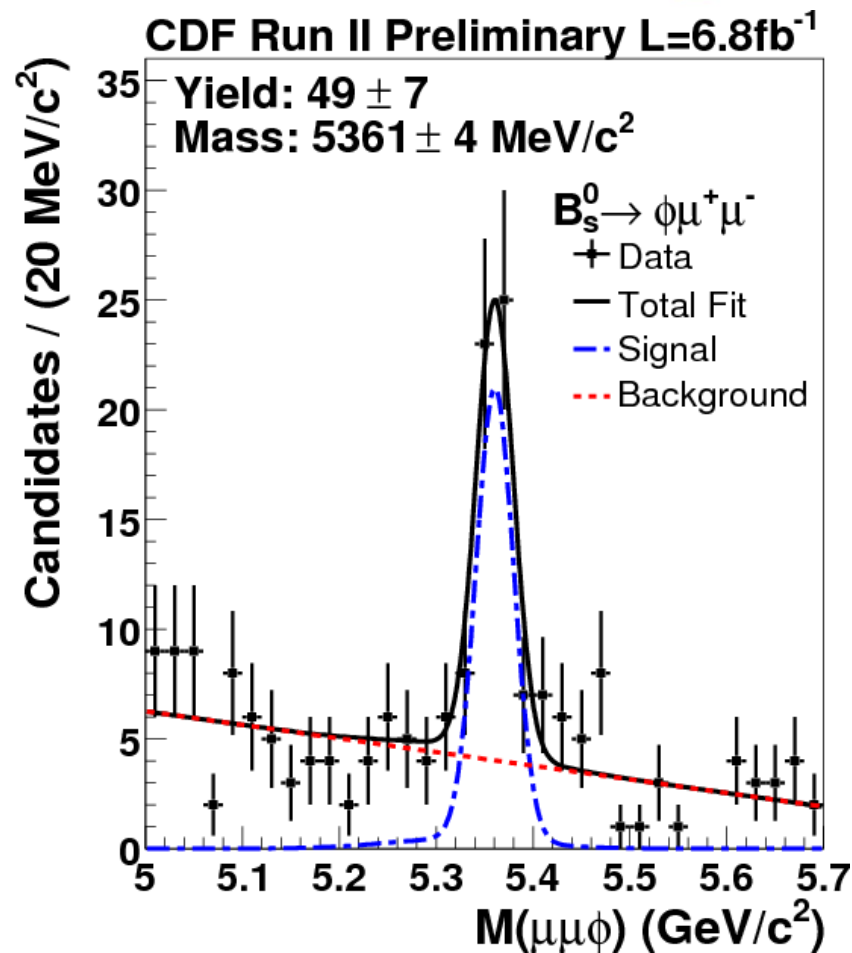
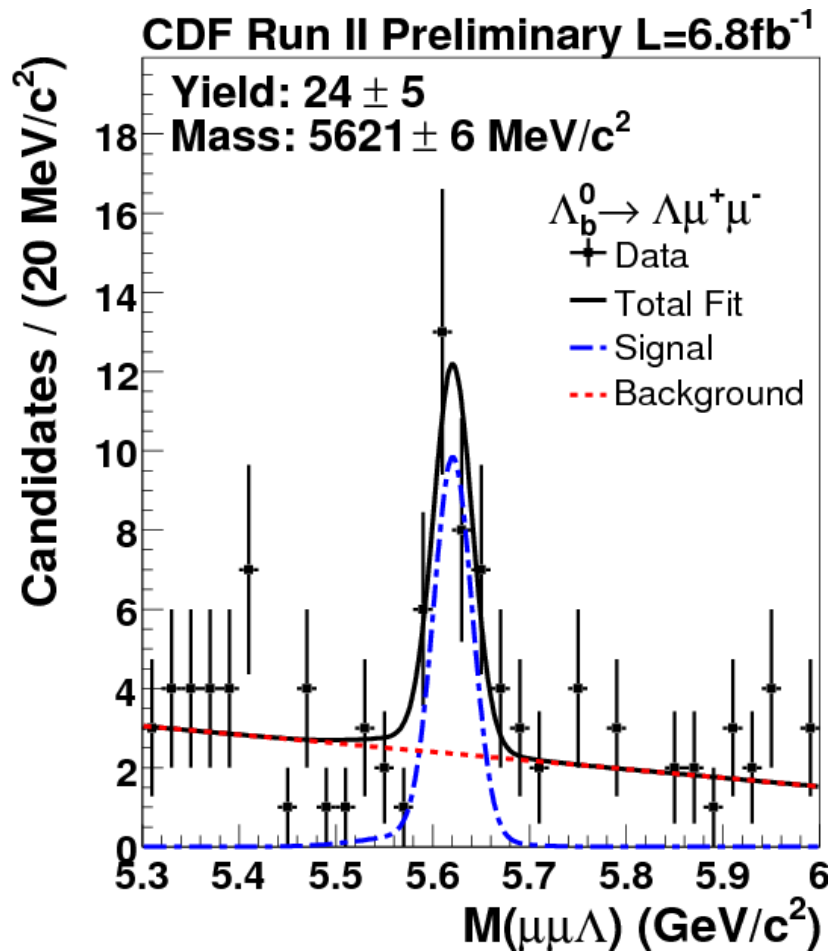
$b \rightarrow s |^+ |^-$



$b \rightarrow s |^+ |^-$



$b \rightarrow s |^+ |^-$



$$B(\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-) = [1.73 \pm 0.42(\text{stat}) \pm 0.55(\text{syst})] \times 10^{-6}$$

$$B(B_s^0 \rightarrow \phi \mu^+ \mu^-) = [1.47 \pm 0.24(\text{stat}) \pm 0.46(\text{syst})] \times 10^{-6}$$

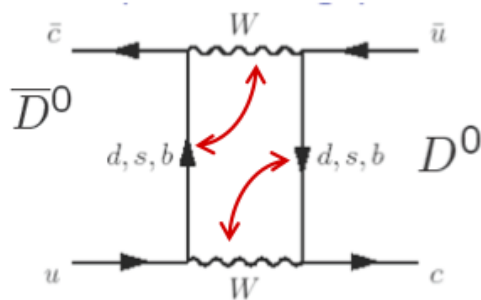
Charm mixing and CPV

- Basic physics is same as in B mixing
- Important because it tests down type quarks in loops
- Complementary to B mixing
- CP violation in SM is very small
 - Mixing in first two generations is real
 - Bottom quark contribution strongly CKM suppressed
- Theory predictions difficult as long distance contributions play important role
- Experiments resolving mixing difficult because of very slow mixing

Mixing in Standard Model is Very Small

- Off-diagonal mass matrix element – two leading terms:

$\Delta C=2$ (short-range)
(contributes mostly to x)



Down-type quarks in loop:

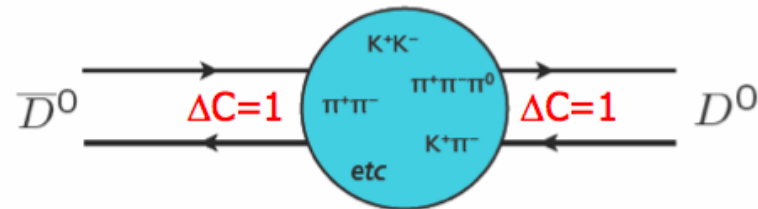
b : CKM-suppressed ($|V_{ub}V_{cb}|^2$)

d, s : GIM-suppressed

$$x \propto (m_s^2 - m_d^2) / m_c^2 \sim 10^{-5}$$

(almost 2 orders of magnitude less than current sensitivity)

Hadronic intermediate states (long-range)



Difficult to compute (need to know all the magnitudes and phases, ...)

Most computations predict x and y in the range 10^{-3} – 10^{-2} and $|x| < |y|$

Recent predictions: $|x| \leq 1\%$, $|y| \leq 1\%$
(consistent with current observation)

x_D, y_D at 1% consistent with SM, BUT
 CPV at 10^{-3} levels would be signal for NP

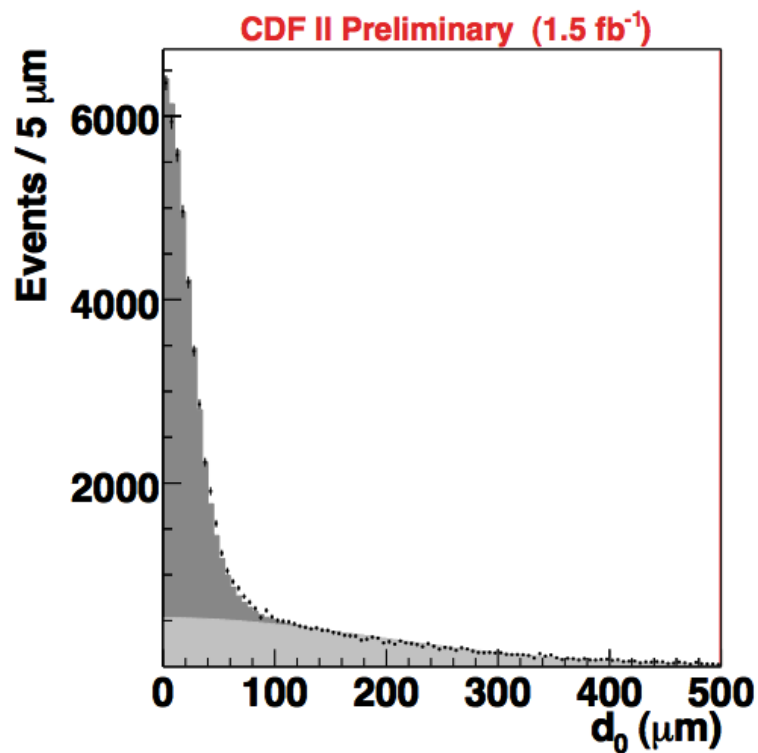
$$x = \Delta m / \Gamma$$

$$y = \Delta \Gamma / 2\Gamma$$

Brian Meadows



Charm mixing - technique

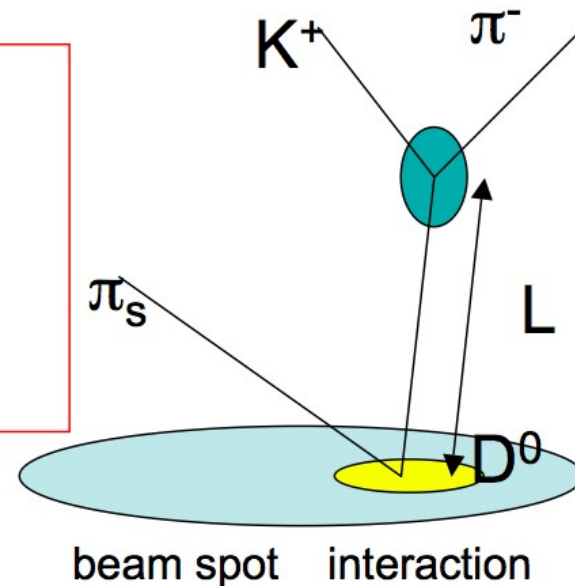


BaBar

3-D flight path

$L \sim 200 \mu\text{m}$

$\sigma_L \sim 100 \mu\text{m}$



Wrong Sign (WS) Decays $D^0 \rightarrow K^+ \pi^-$

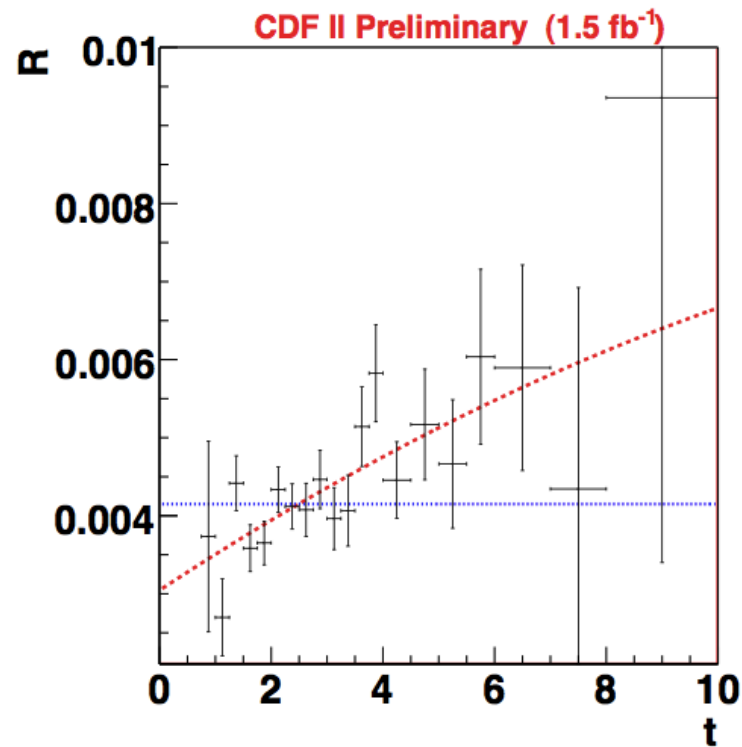
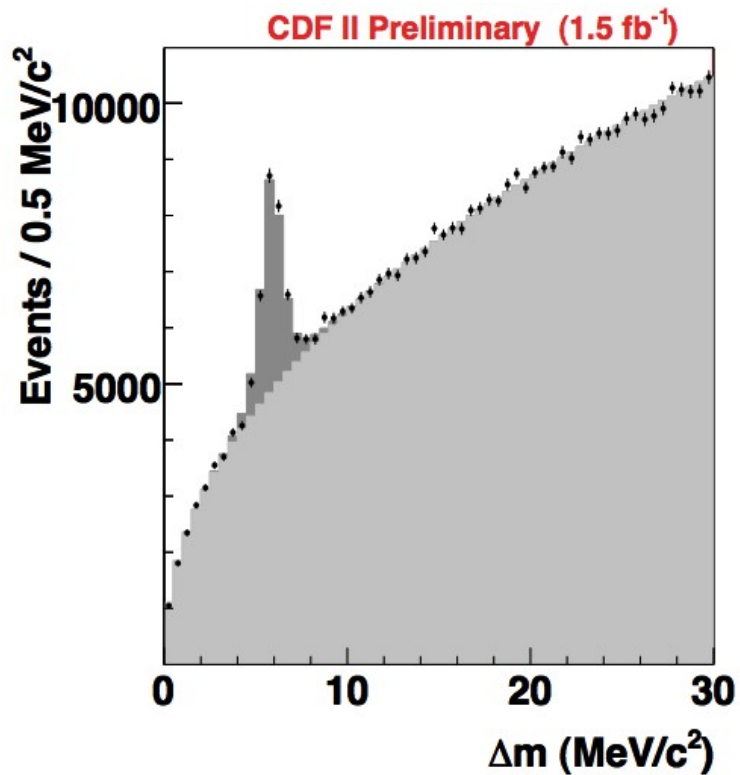
- The **WS** decay rate R_{WS} is:

$$R_{WS} = e^{-\Gamma t} |\mathcal{A}_{\bar{f}}|^2 \left[\underbrace{1}_{\text{Direct decay}} + \underbrace{\lambda_{\bar{f}} y'_D(\Gamma t)}_{\text{Interference}} + \underbrace{\frac{|\lambda_{\bar{f}}|^2}{4} (x_D'^2 + y_D'^2) (\Gamma t)^2}_{\text{Decay through Mixing}} \right]$$

- Since $|\lambda_{\bar{f}}| \gg 1$, all three terms are comparable
- For “right-sign” (**RS**) decays $D^0 \rightarrow K^- \pi^+$ though, $|\lambda_f| \ll 1$, so 2nd two terms are negligible and R_{RS} is approximately exponential. $R_{RS} \approx e^{-\Gamma t} |\mathcal{A}_f|^2$

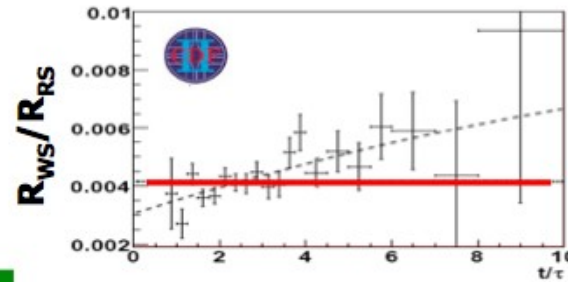
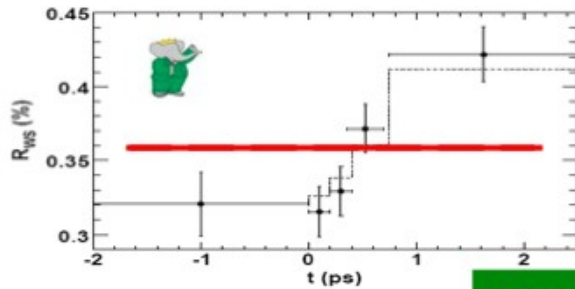


Charm mixing in DCS decays



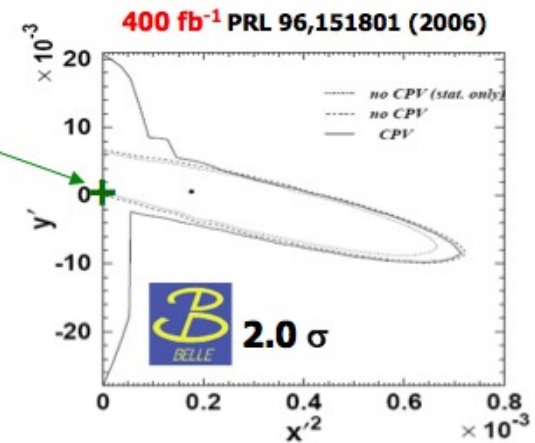
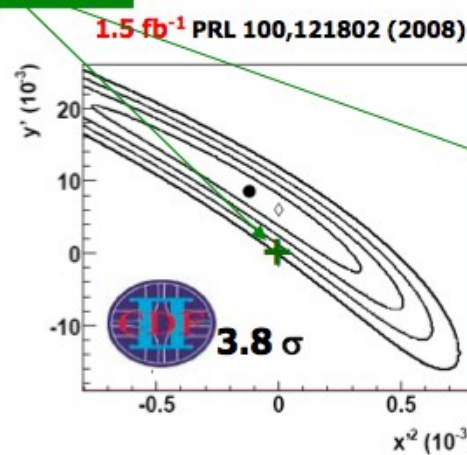
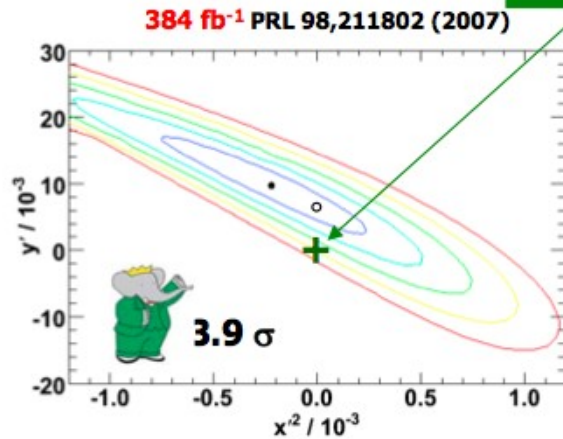
Evidence for Mixing in $D^0 \rightarrow K^+\pi^-$

Mixing seen by Babar and CDF in time-dependence of the R_{WS}/R_{RS} ratio



— No Mixing

No Mixing

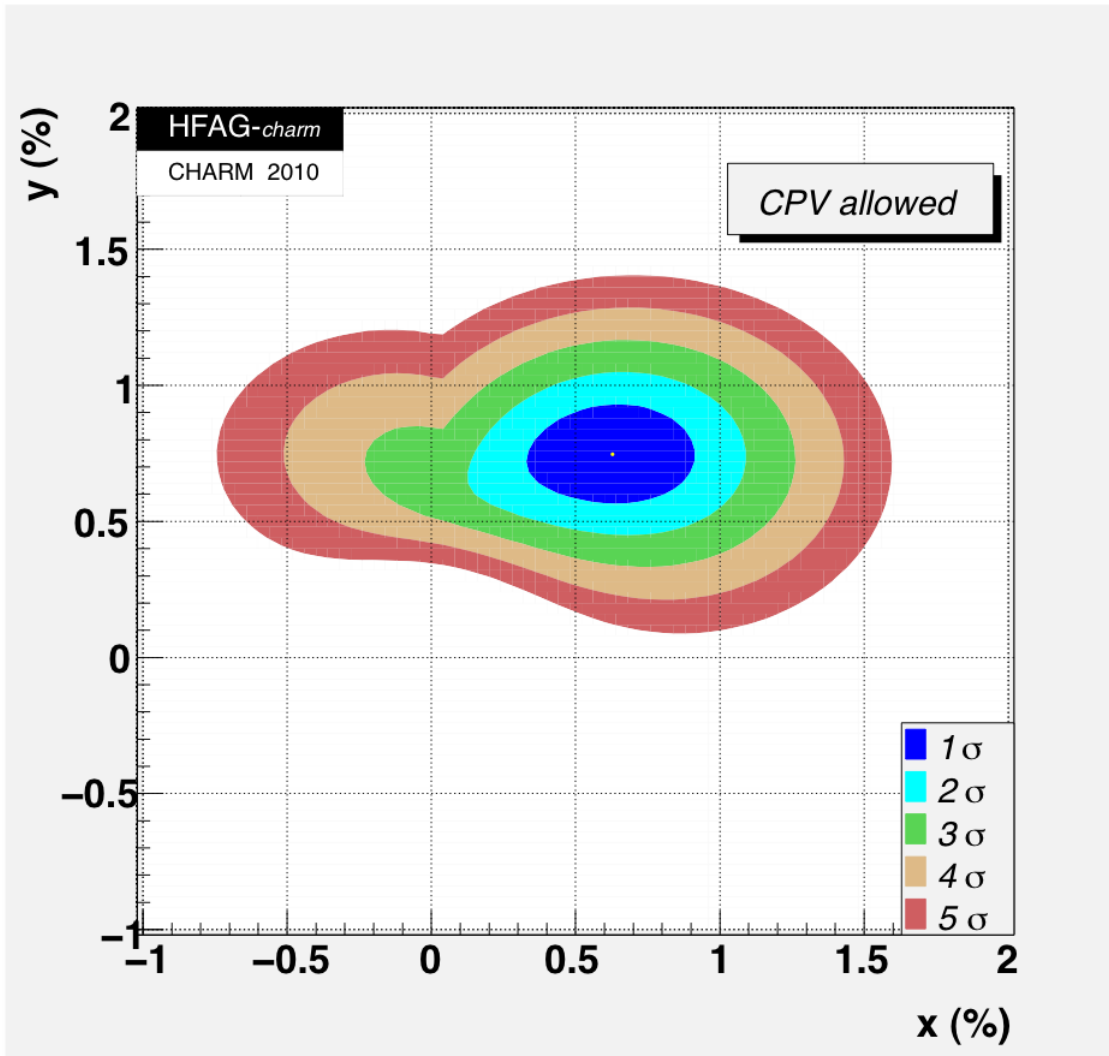


Belle result was the most sensitive, BUT evidence for mixing not significant !



Charm mixing - global status

$$x = \Delta m / \Gamma$$
$$y = \Delta \Gamma / 2\Gamma$$



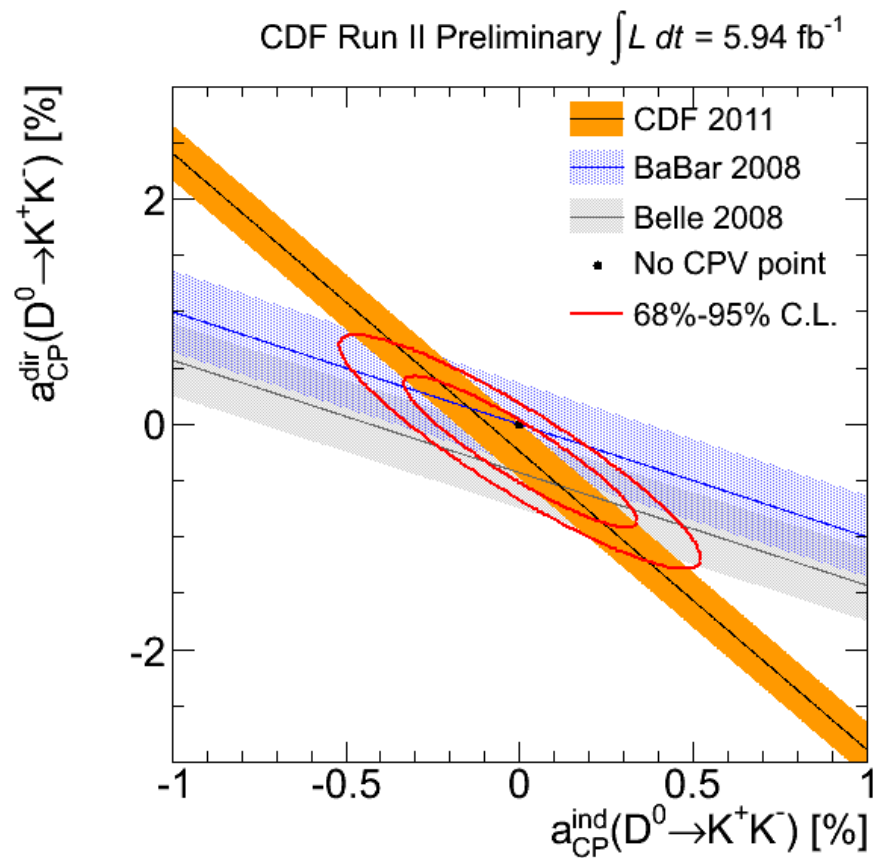
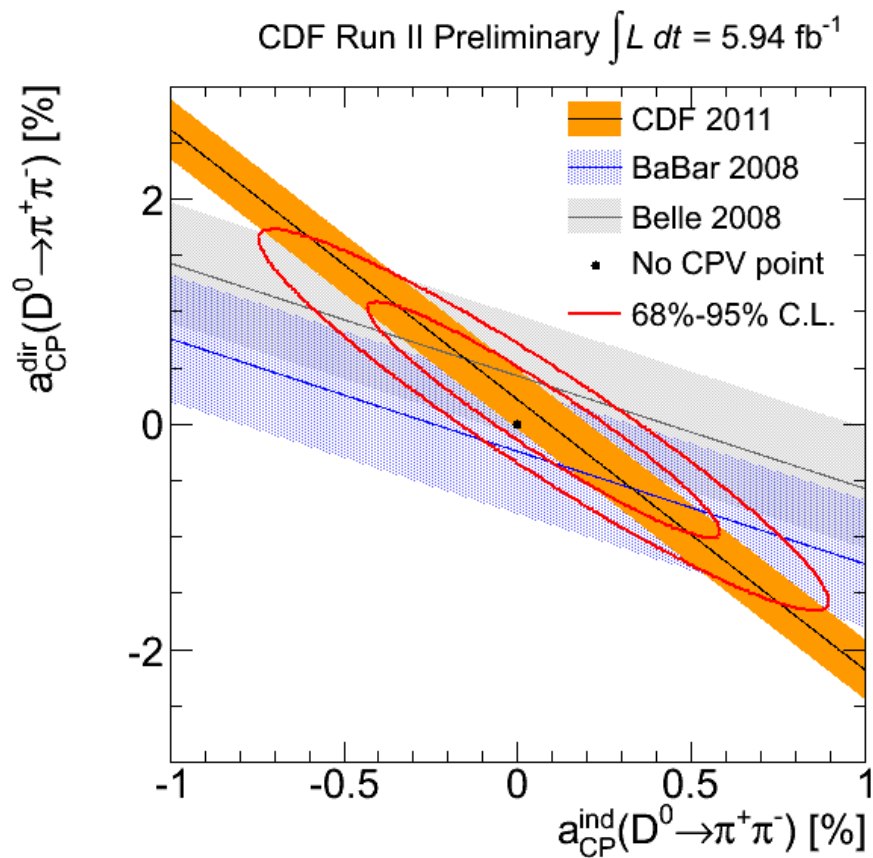
- Concentrate on D^0 here
- What is usually done is time integrated measurement in common final state
 - Tagged by D^*
- Time integrated measurements sensitive also to time-dependent CPV as oscillation is very slow

$$A_{\text{CP}}(h^+h^-) = \frac{\Gamma(D^0 \rightarrow h^+h^-) - \Gamma(\bar{D}^0 \rightarrow h^+h^-)}{\Gamma(D^0 \rightarrow h^+h^-) + \Gamma(\bar{D}^0 \rightarrow h^+h^-)}$$

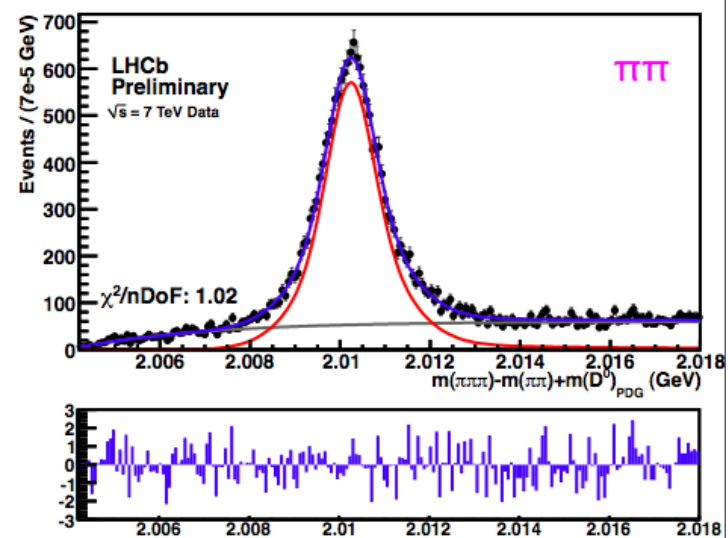
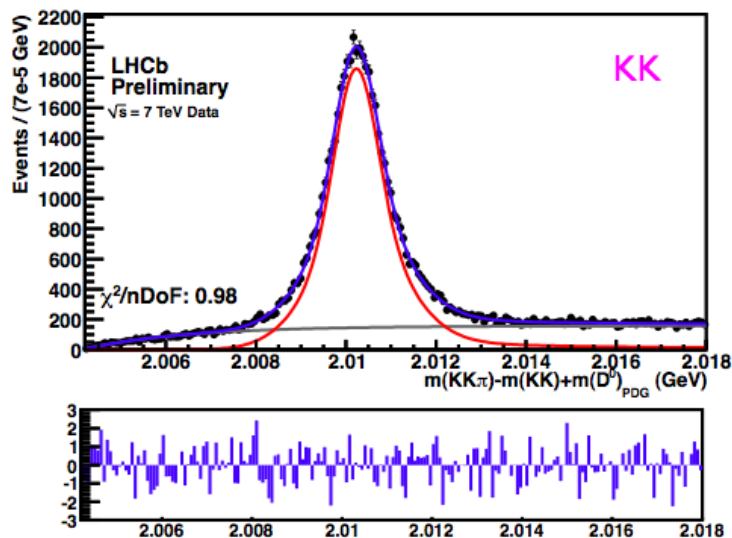
$$A_{\text{CP}}(h^+h^-) = a_{\text{CP}}^{\text{dir}}(h^+h^-) + \int_0^{\infty} A_{\text{CP}}(t)D(t)dt \approx a_{\text{CP}}^{\text{dir}}(h^+h^-) + \frac{\langle t \rangle}{\tau} a_{\text{CP}}^{\text{ind}}(h^+h^-).$$

- Method uses also $K\pi$ decays to measure detector asymmetry
- LHCb has additional complication of production asymmetry

CDF result



- Measure difference $A(KK)-A(\pi\pi)$
- Detector and production asymmetry cancels in first order
- Mainly sensitive to direct CP violation as indirect is independent of final state



Preliminary: 2010 data, 38 pb^{-1}

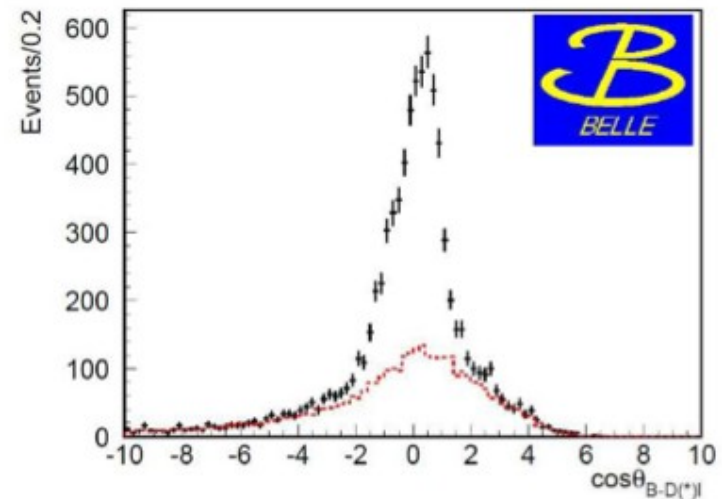
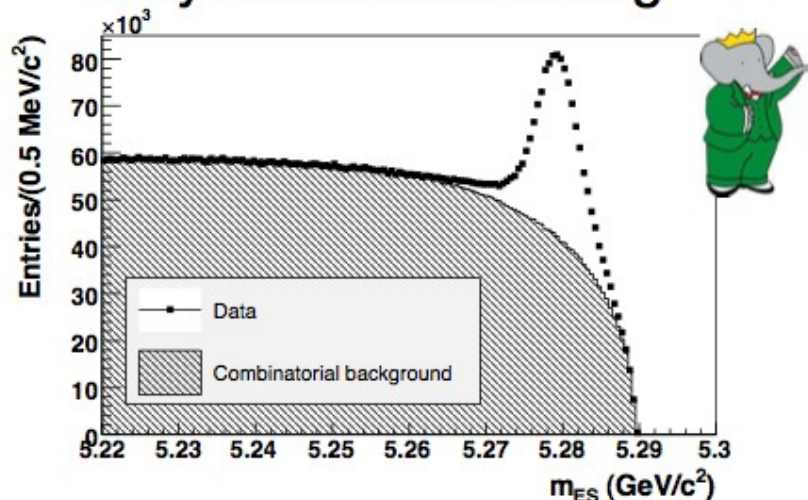
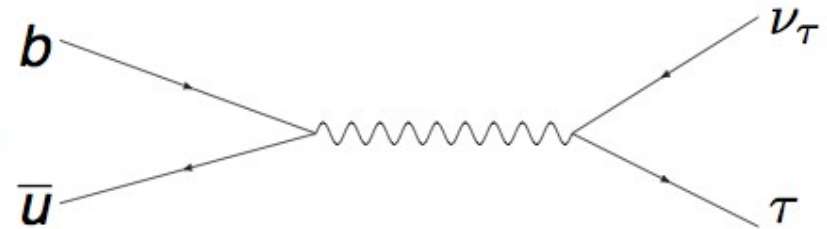
$$A_{CP}(KK) - A_{CP}(\pi\pi) = (-0.275 \pm 0.701 \pm 0.25)\%$$

$B \rightarrow \tau \nu$

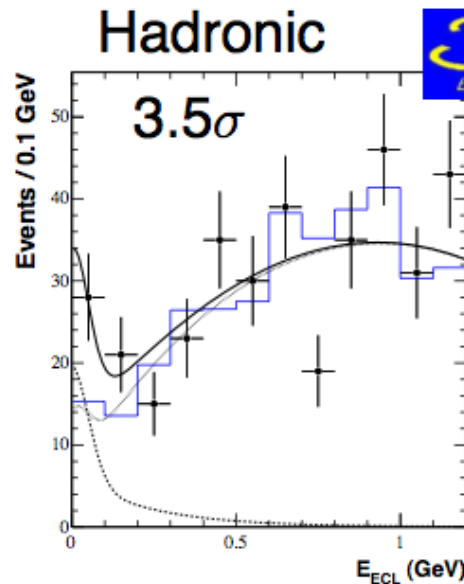
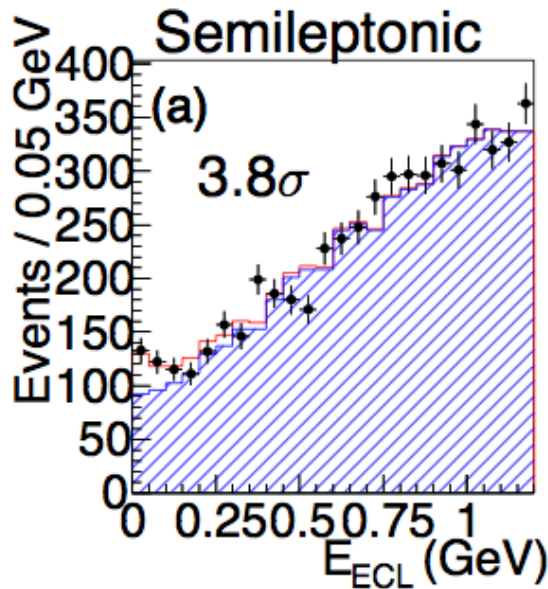
- SM branching fraction given by

$$BF = \frac{G_F^2 m_B}{8\pi} m_l^2 \left(1 - \frac{m_l^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$

- One can extract f_B or $|V_{ub}|$
- SM $BF = 1.20 \pm 0.25 \times 10^{-4}$
- Practically only B-factories thanks to clean environment
- After reconstructing tag B and all charged particles from signal B , only neutrinos missing



$B \rightarrow \tau \nu$



$$S: 1.65^{+0.38}_{-0.37} \times 10^{-4}$$

$$H: 1.79^{+0.56}_{-0.49} \times 10^{-4}$$

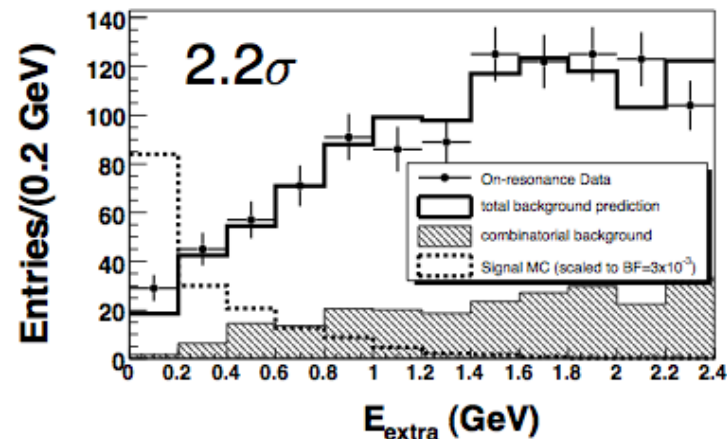
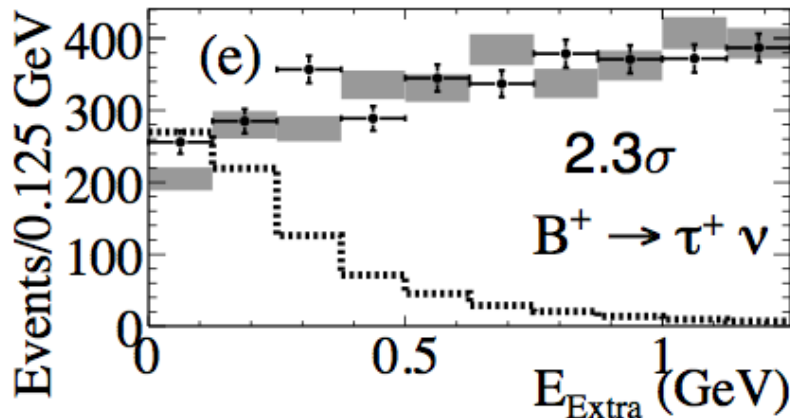
hep-ex/0809.3834,
PRL 97, 251802 (2006)

$$S: 1.7 \pm 0.8 \pm 0.2 \times 10^{-4}$$

$$H: 1.8^{+0.9}_{-0.8} \pm 0.4 \pm 0.2 \times 10^{-4}$$

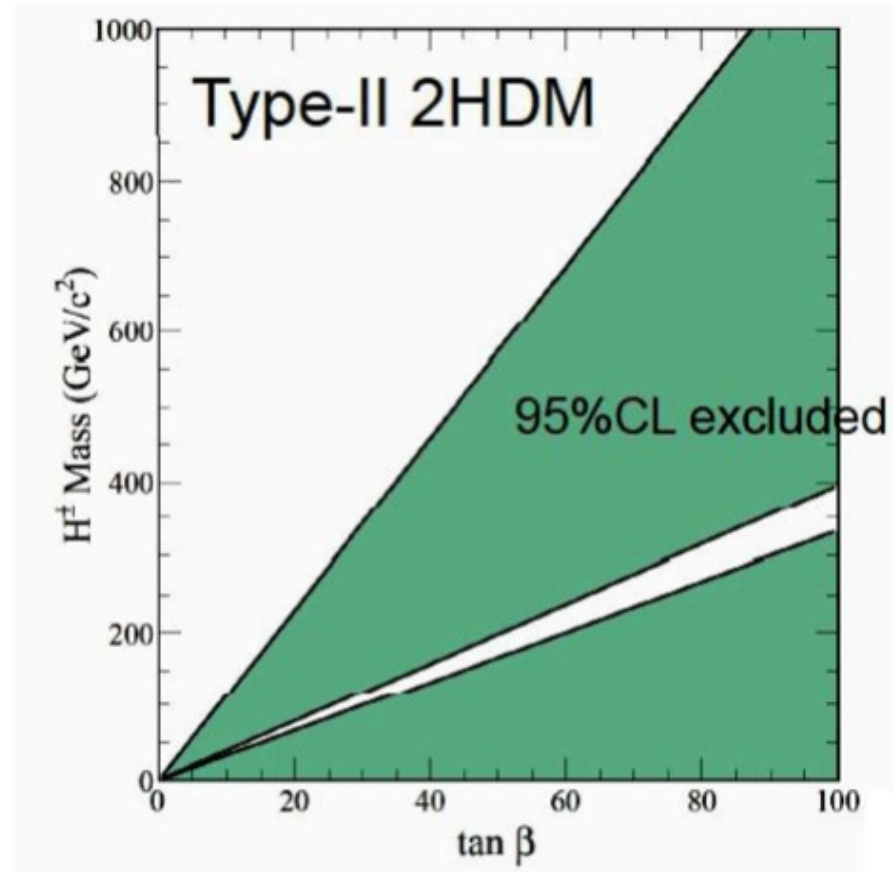
PRD 81, 051101 (2010)

PRD 77, 011107 (2008)

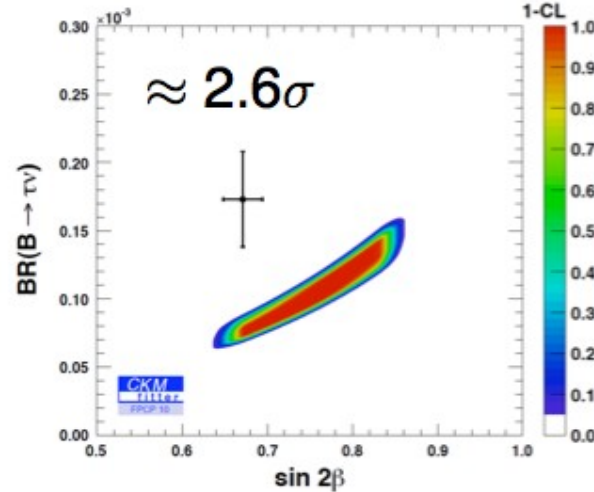
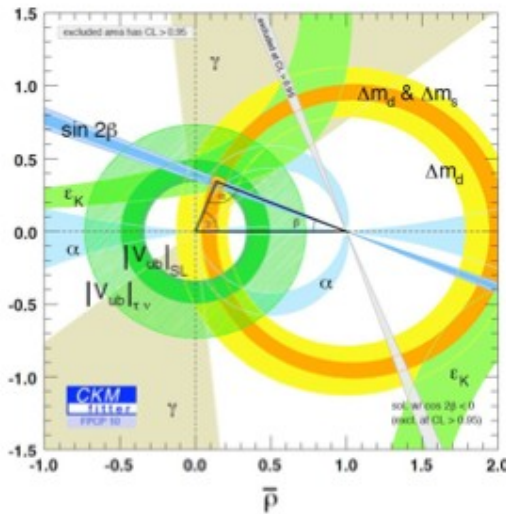


$B \rightarrow \tau \nu$

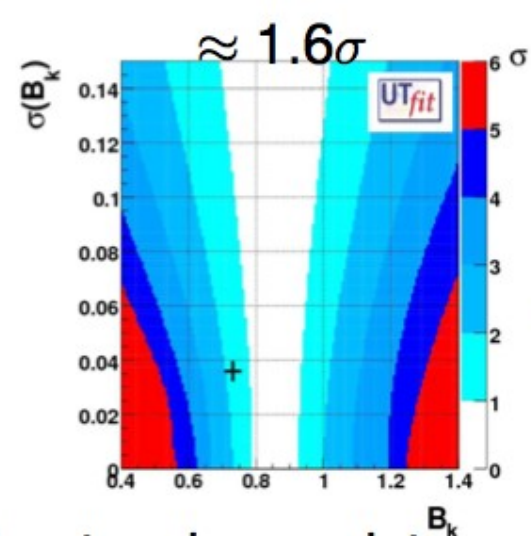
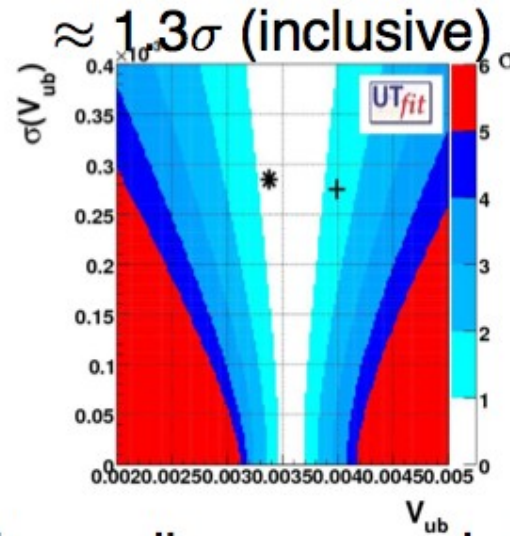
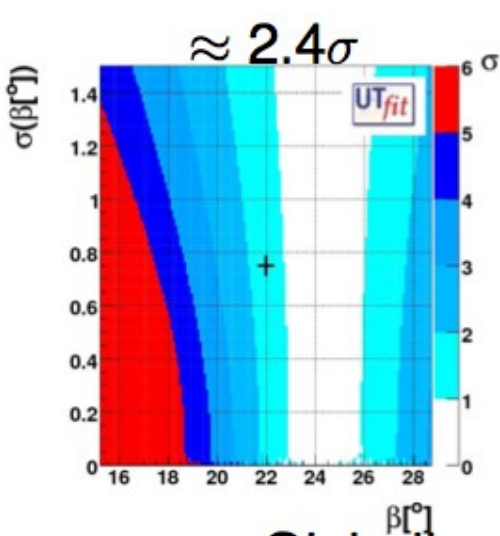
- Naive average of exp. results
 $BF_{exp} = 1.73 \pm 0.35 \times 10^{-4}$
- SM prediction
 $BF_{SM} = 1.20 \pm 0.25 \times 10^{-4}$
- Effect of charged Higgs
 $BF_{exp} = BF_{SM} \times r_H$
$$r_H = \left(1 - \frac{m_B^2 \tan^2 \beta}{m_H^2} \frac{1}{1 + \epsilon_0 \tan \beta} \right)^2$$
- For Type-II 2HDM $\epsilon_0 = 0$



Global status of the triangle



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Globally no large discrepancy, but few tensions exist

Final remarks

- Lot of work was done in last few years to discover new physics
- Many places where some tension with SM was observed with small statistics
- Latest results from this summer now agree quite well with SM
- Despite the agreement with SM, there is still quite some room for new physics