

Coronal loop seismology

using damping of standing kink oscillations by mode coupling

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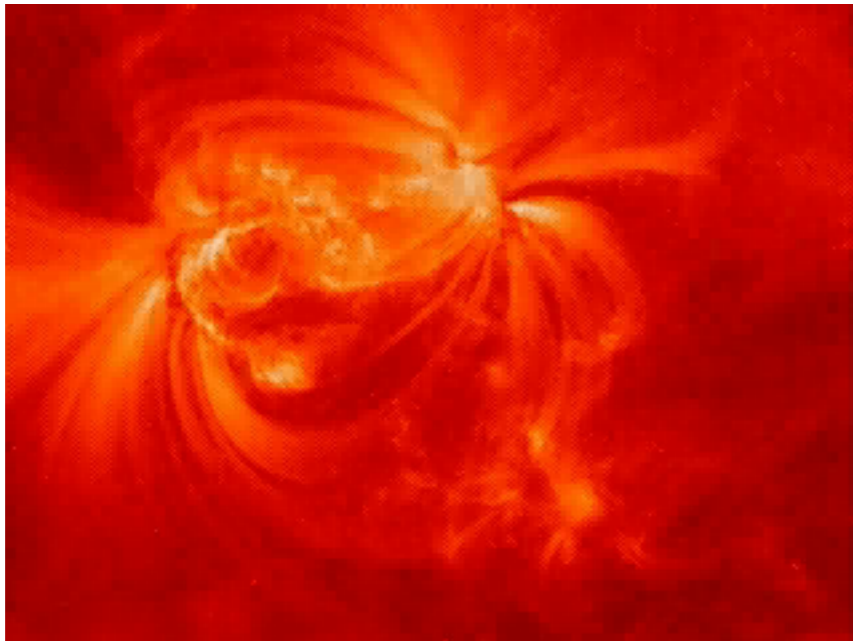
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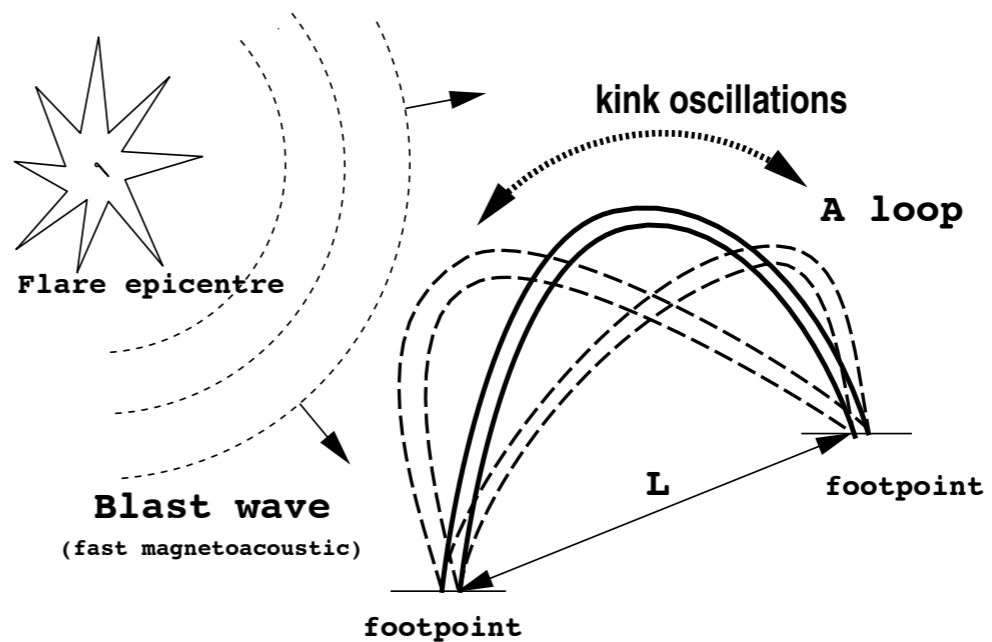
Overview

- Introduction (damped kink oscillations of coronal loops)
- Mode coupling by resonant absorption as damping mechanism
- Gaussian damping regime
 - numerical simulations
 - analytical theory
 - observational evidence
- Use of both exponential and Gaussian damping regimes for loop seismology
 - application to SDO/AIA observations
- Implications for heating

Standing kink modes of coronal loops



TRACE observation 14 July 1998



$$\frac{\omega}{k} = \frac{2L}{P} \approx \begin{cases} 1020 \pm 132 \text{ km s}^{-1} & \text{(14th July, 1998),} \\ 1030 \pm 410 \text{ km s}^{-1} & \text{(4th July, 1999).} \end{cases} \quad k = \pi/L$$

long wavelength limit: $ka \ll 1$

slow $\frac{\omega}{k} \approx C_{T0} \equiv \frac{C_{s0} C_{A0}}{(C_{s0}^2 + C_{A0}^2)^{1/2}}$

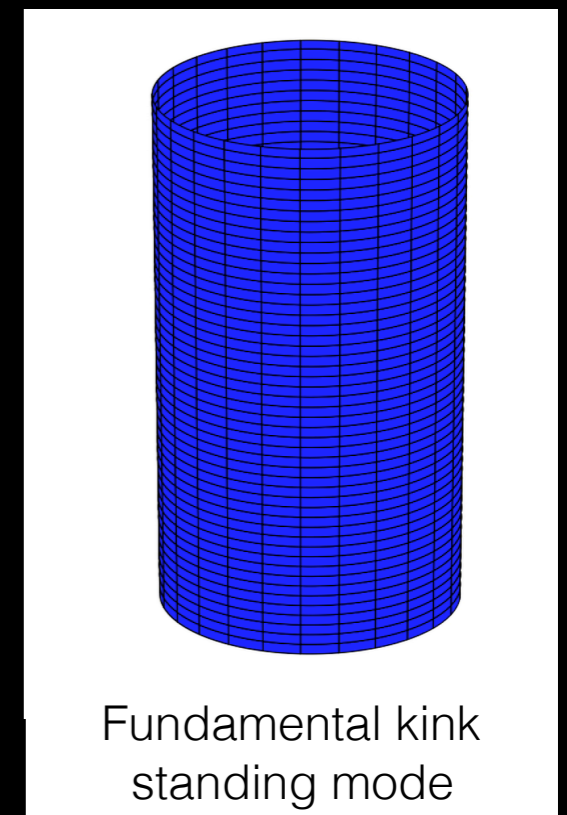
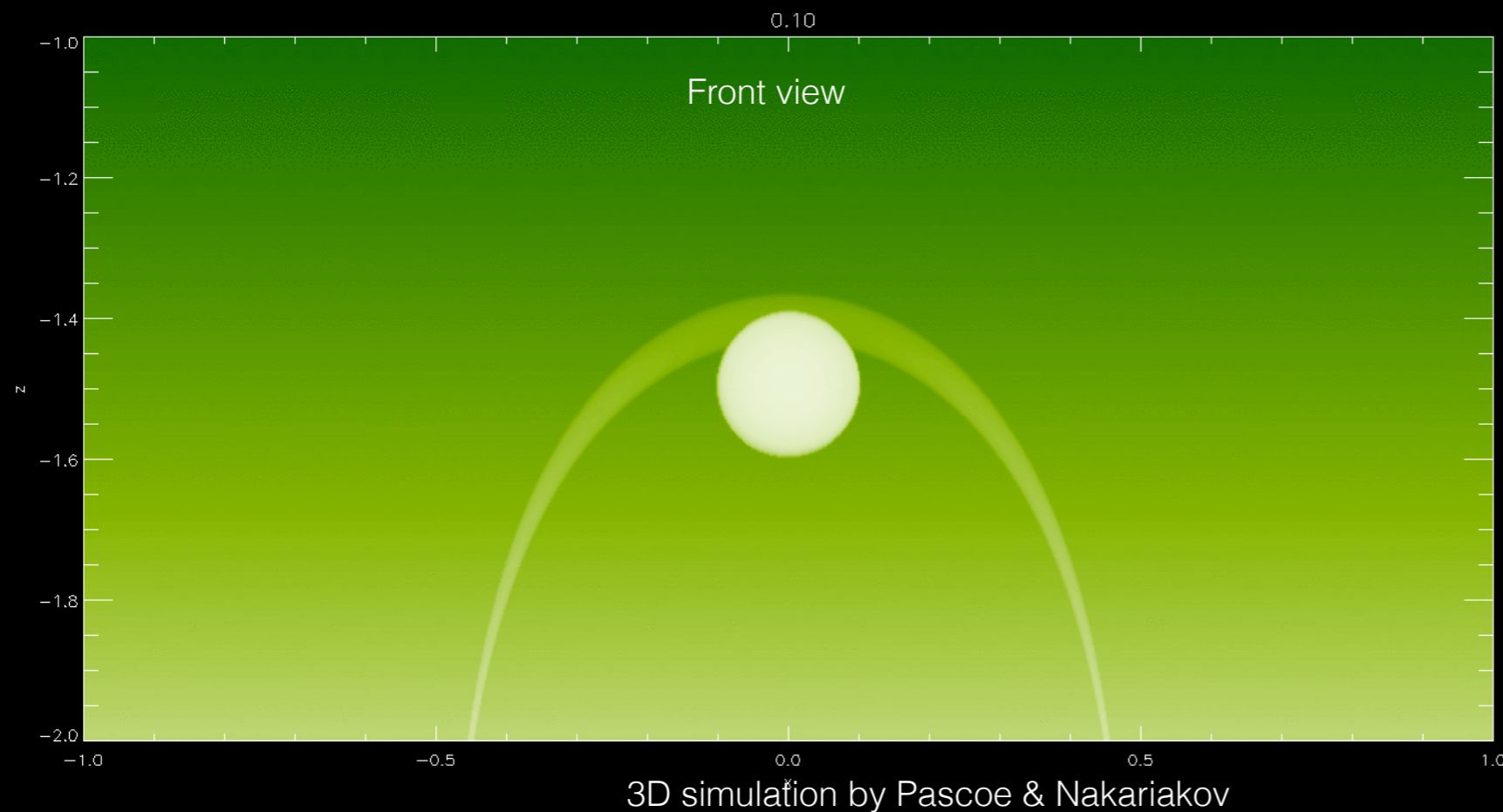
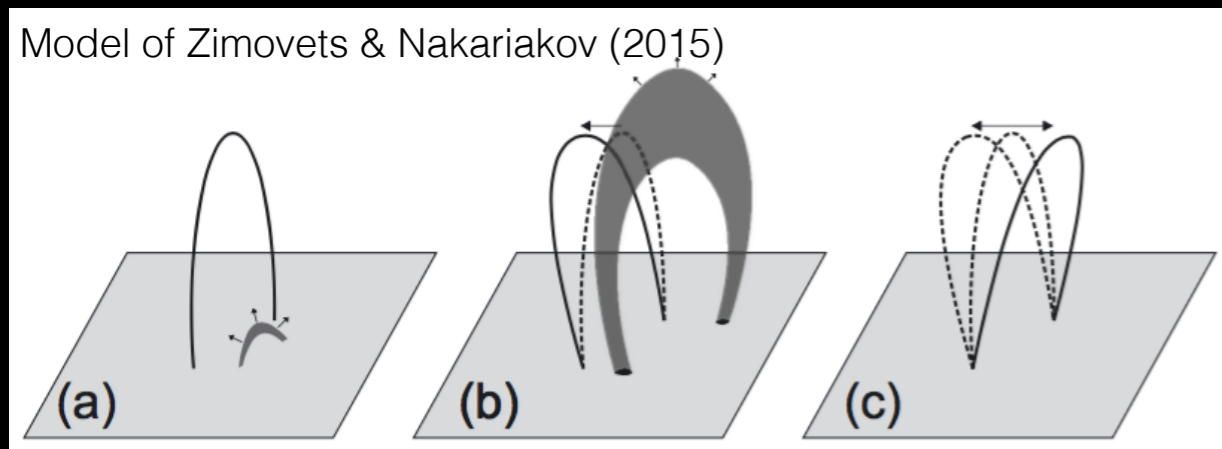
$\frac{\omega}{k} \approx C_k \equiv \left(\frac{2}{1 + \rho_e/\rho_0} \right)^{1/2} C_{A0}$ fast

$$B_0 = (4\pi\rho_0)^{1/2} C_{A0} = \frac{\sqrt{2} \pi^{3/2} L}{P} \sqrt{\rho_0(1 + \rho_e/\rho_0)}$$

$$B = 13 \pm 9 \text{ G}$$

Kink oscillations excited by coronal mass ejections (CMEs)

- Kink mode = transverse oscillation of loop
- 76% of oscillations associated with CMEs observed in white light emission
- 98% accompanied by lower coronal eruptions

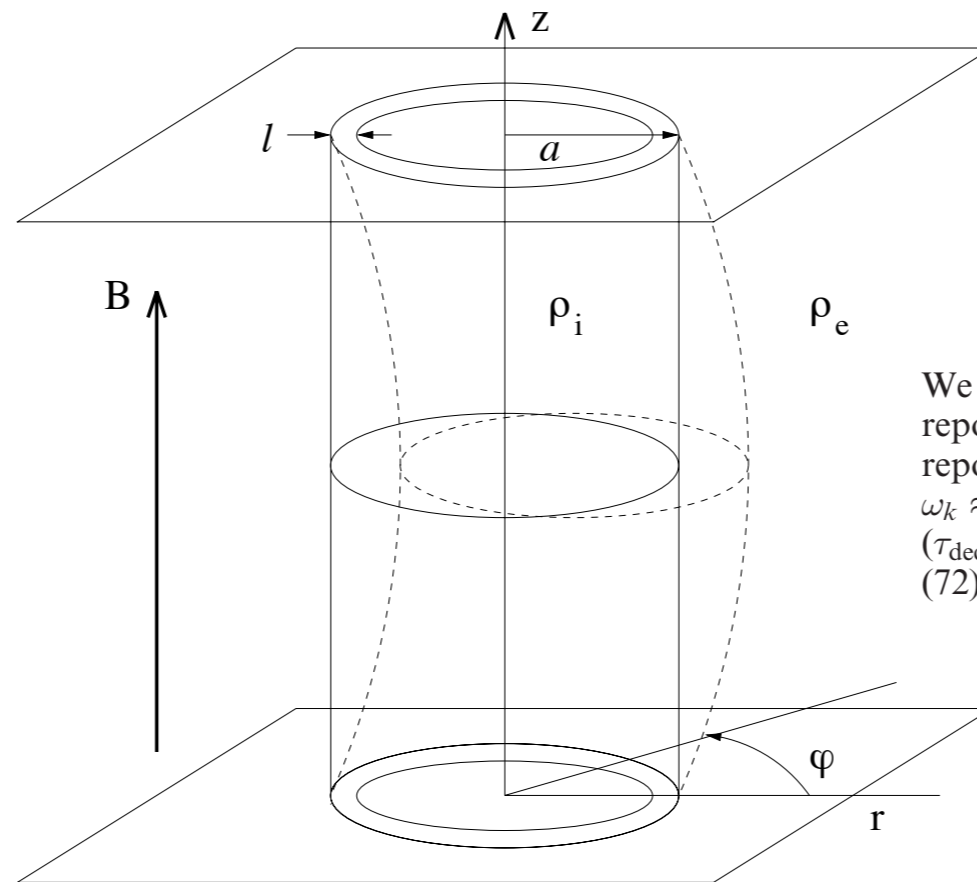


Resonant absorption as a damping mechanism

- Mode coupling studied by Hollweg & Yang (1988), Goossens et al. (1992) etc.
- Ruderman & Roberts (2002) considered in the context of the rapid damping of coronal loop oscillations observed by TRACE
- Modelled damping of kink modes is due to resonant absorption, acting in the inhomogeneous regions of the tube

Exponentially damped oscillation with

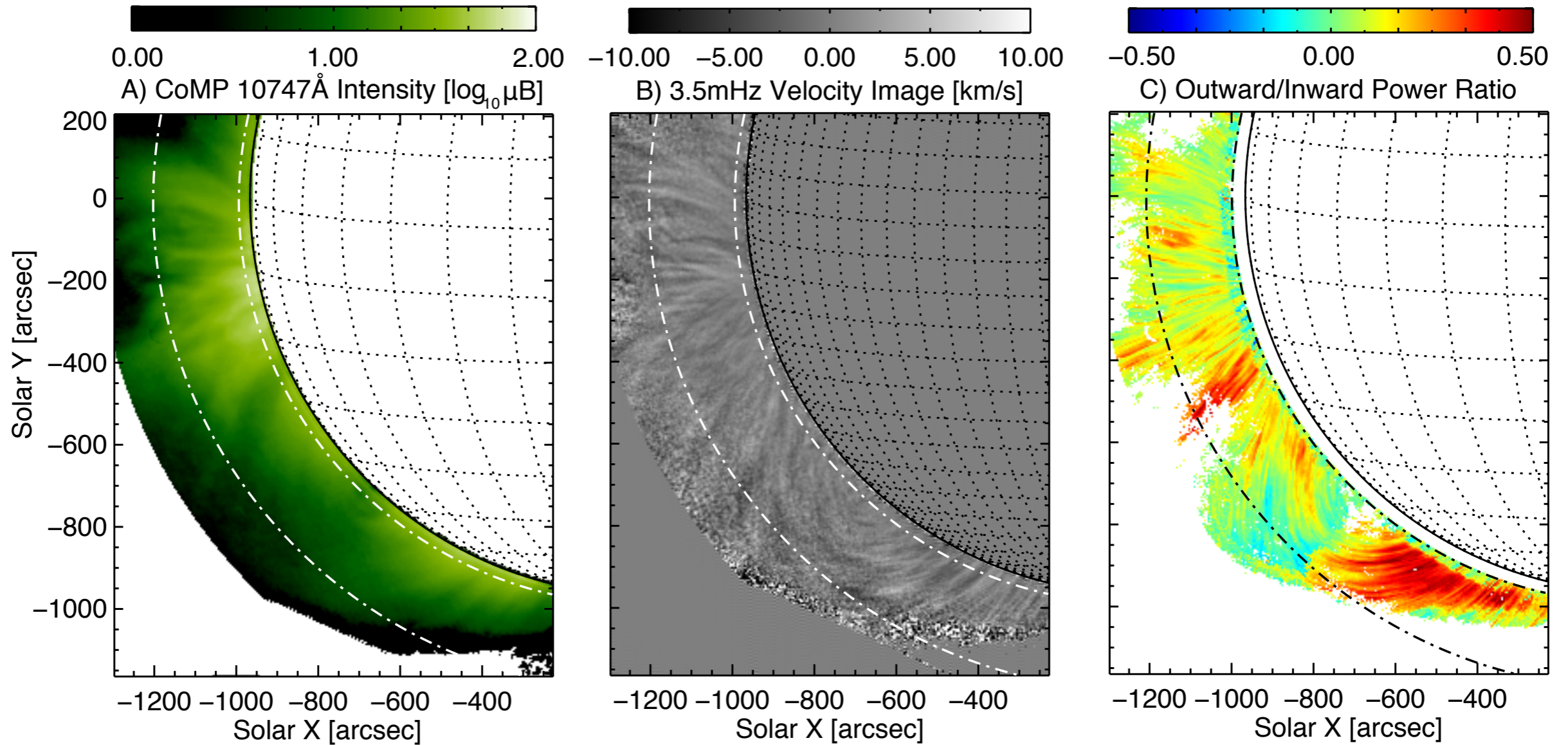
$$\tau_{\text{decay}} = \frac{2a\rho_i + \rho_e}{\pi\ell\rho_i - \rho_e} \tau$$



We consider this result in relation to the observational data reported by Nakariakov et al. (1999). These authors reported a coronal loop oscillation with frequency $\omega_k \approx 0.024 \text{ s}^{-1}$ ($\tau = 256 \text{ s}$) and decrement $\gamma \approx 0.0011 \text{ s}^{-1}$ ($\tau_{\text{decay}} = 870 \text{ s}$). Taking $\rho_i = 10\rho_e$, we obtain from equation (72) that $\ell/a \approx 0.23$.

- Transfer of energy from kink mode to Alfvén (azimuthal) oscillations within inhomogeneous layer
- Only loops with small inhomogeneous layers are able to support coherent oscillations

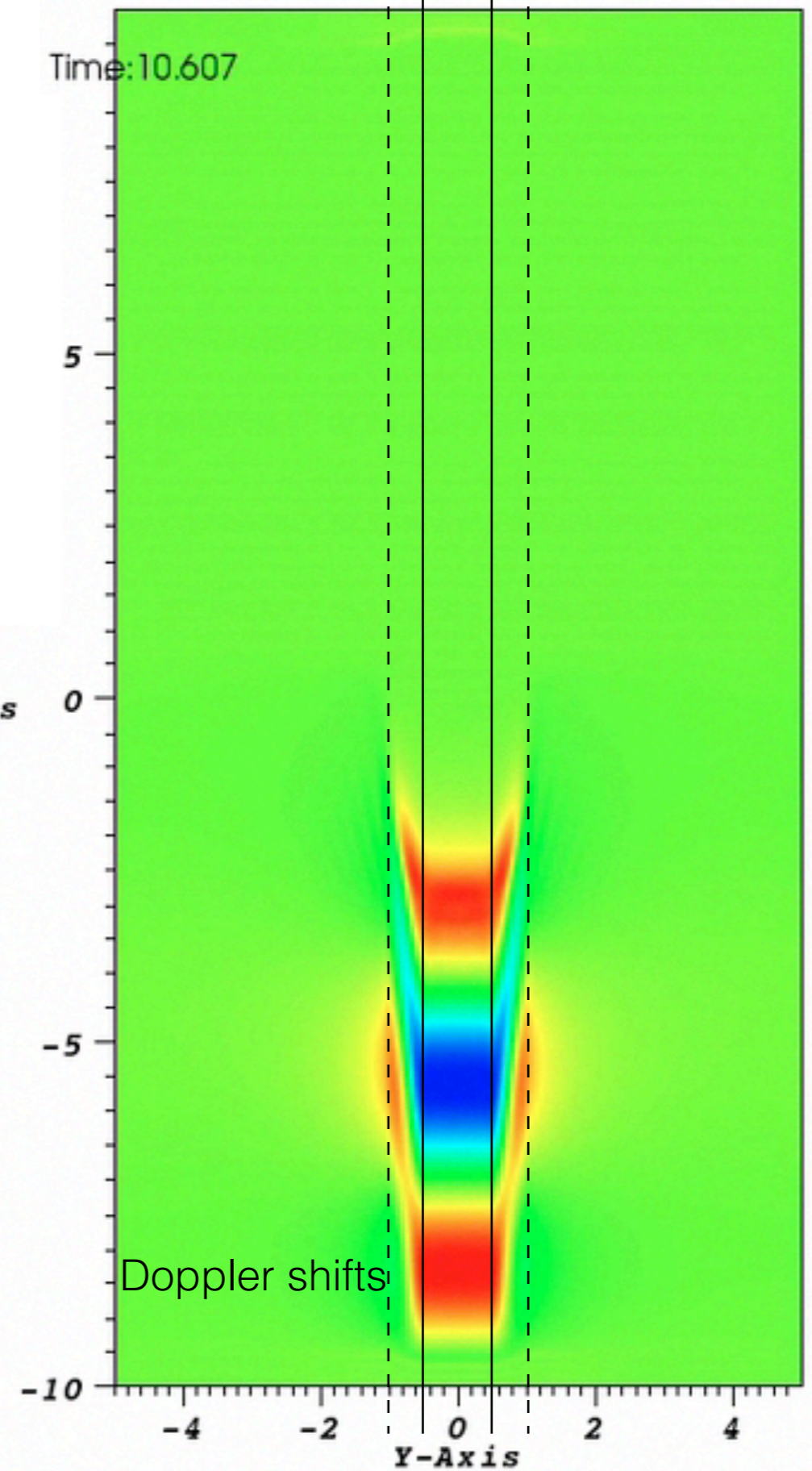
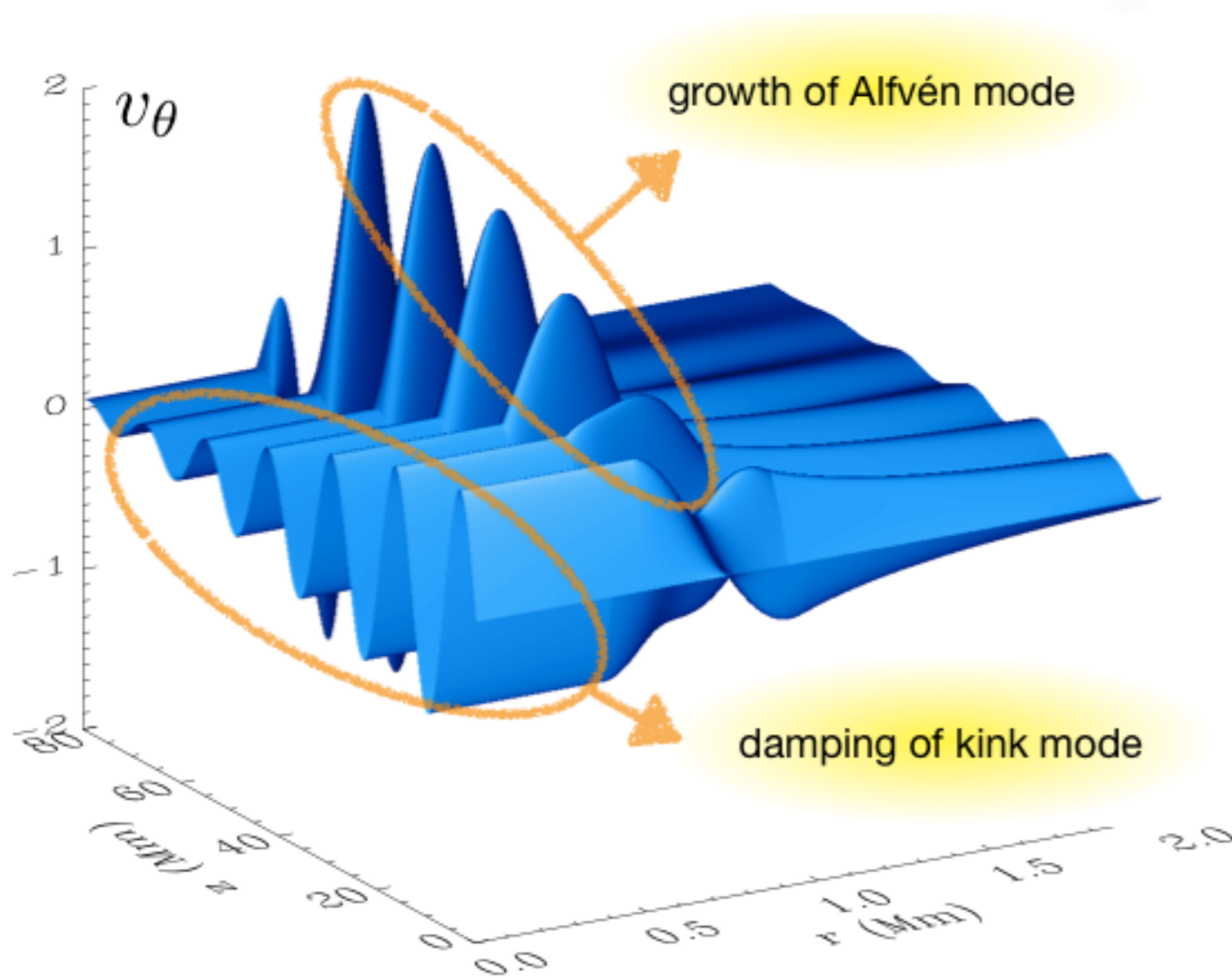
Ubiquitous propagating kink waves



- Ubiquitous transverse velocity perturbations propagating along field lines
- Broadband power spectrum centred on 5 minutes
- Strongly damped

Tomczyk et al. (2007)
Tomczyk & McIntosh (2009)

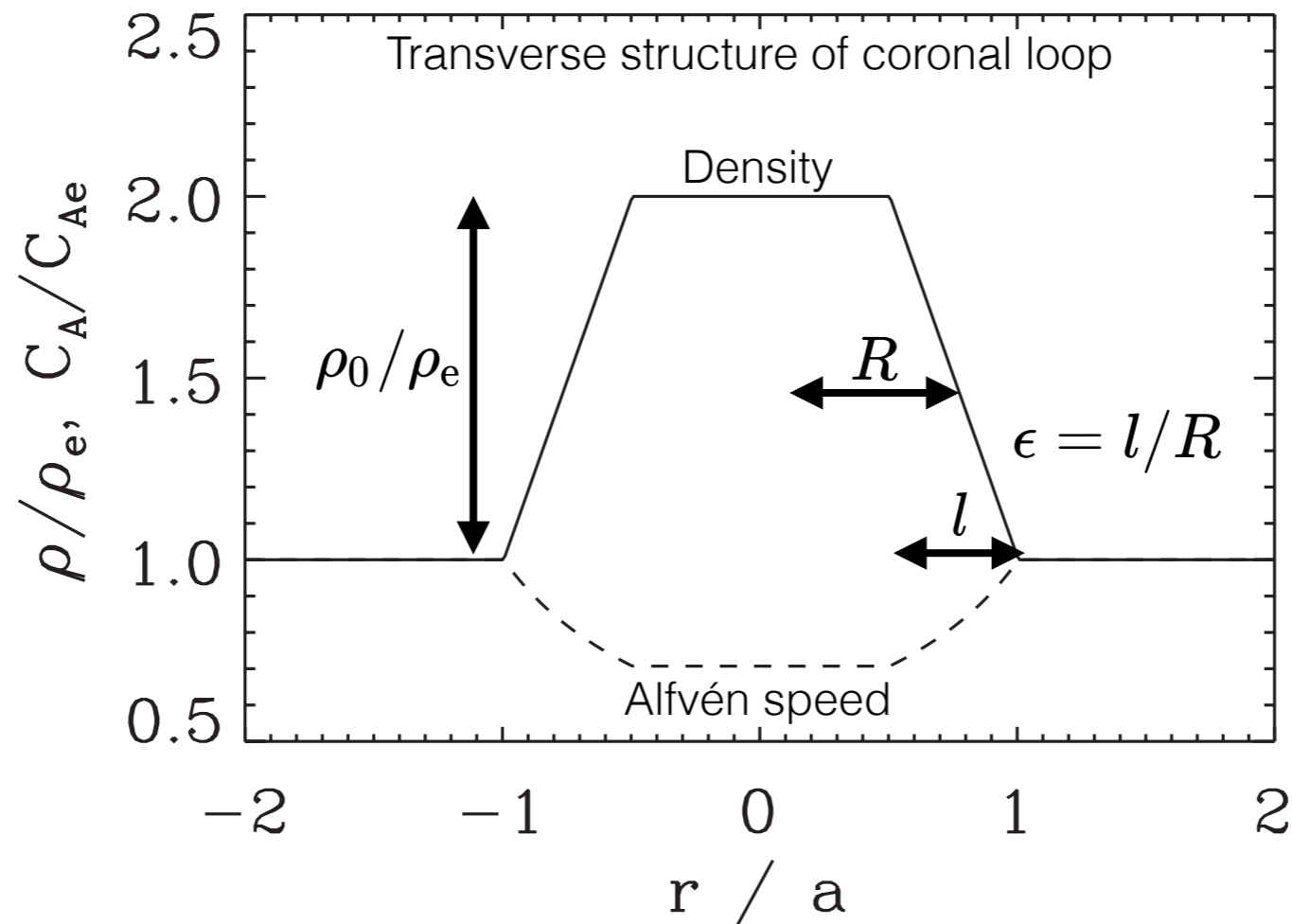
Wave energy being localised by resonance



- Finite wave energy in system ("kick" given by the CME) so kink wave damps as Alfvén wave grows
- Alfvén wave is difficult to observe (incompressible, weak Doppler signal) so only damped kink wave is seen

Damping of kink waves by resonant absorption

- The **inhomogeneous layer** provides a continuous range of Alfvén speeds, and resonance occurs where the Alfvén speed equals the kink speed



- Resonant absorption transfers energy from kink mode (collective motion) to Alfvén mode (localised, unresolved motion)
- Kink mode impulsively excited so transfer of wave energy to Alfvén mode causes damped kink oscillations

What is the kink mode damping envelope?

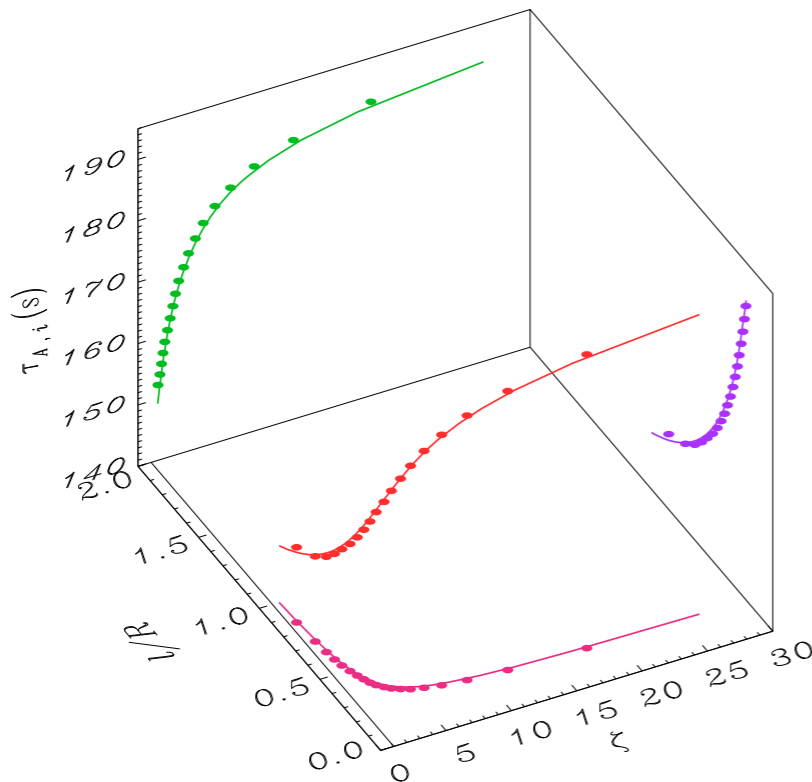
- Original application to coronal loops describes exponential damping envelope for kink oscillation:

$$A \propto \exp\left(-\frac{t}{\tau_d}\right) \quad \tau_d \propto \frac{1}{\epsilon} \frac{\rho_0/\rho_e + 1}{\rho_0/\rho_e - 1}$$

Ruderman & Roberts (2002)
Goossens et al. (2002)

- Seismological inversion problem is ill-posed since τ_d (one observable) depends on density contrast ratio and ϵ (two unknowns)

Arregui et al. (2007)
Goossens et al. (2008)



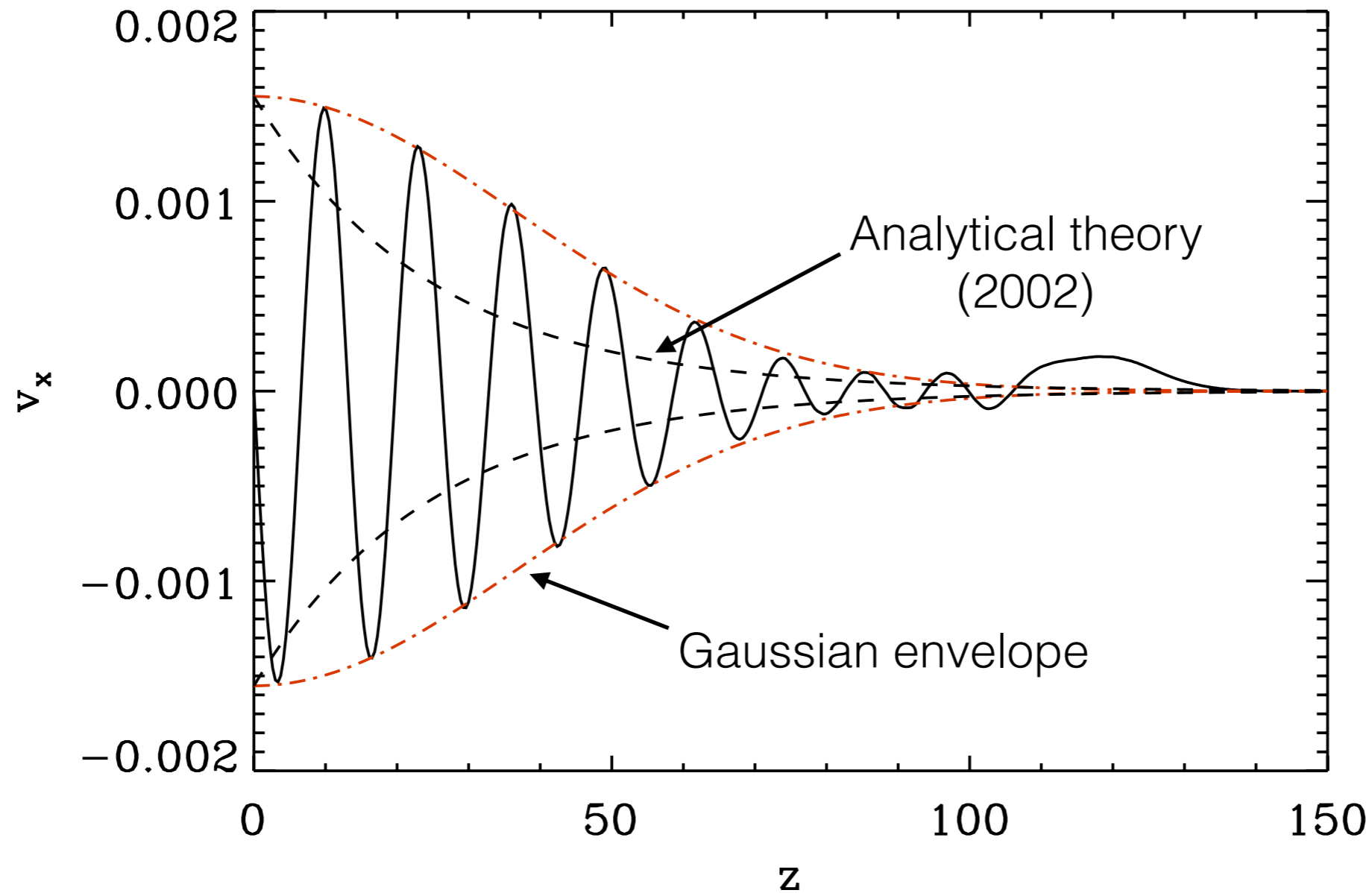
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Ruderman & Roberts (2002)

- Inversion problem has infinite solutions, though bounding values can be estimated (e.g. Arregui & Asensio Ramos 2014)

What is the kink mode damping envelope?

- Numerical simulations of kink oscillations in low density contrast loops discovered Gaussian damping regime:

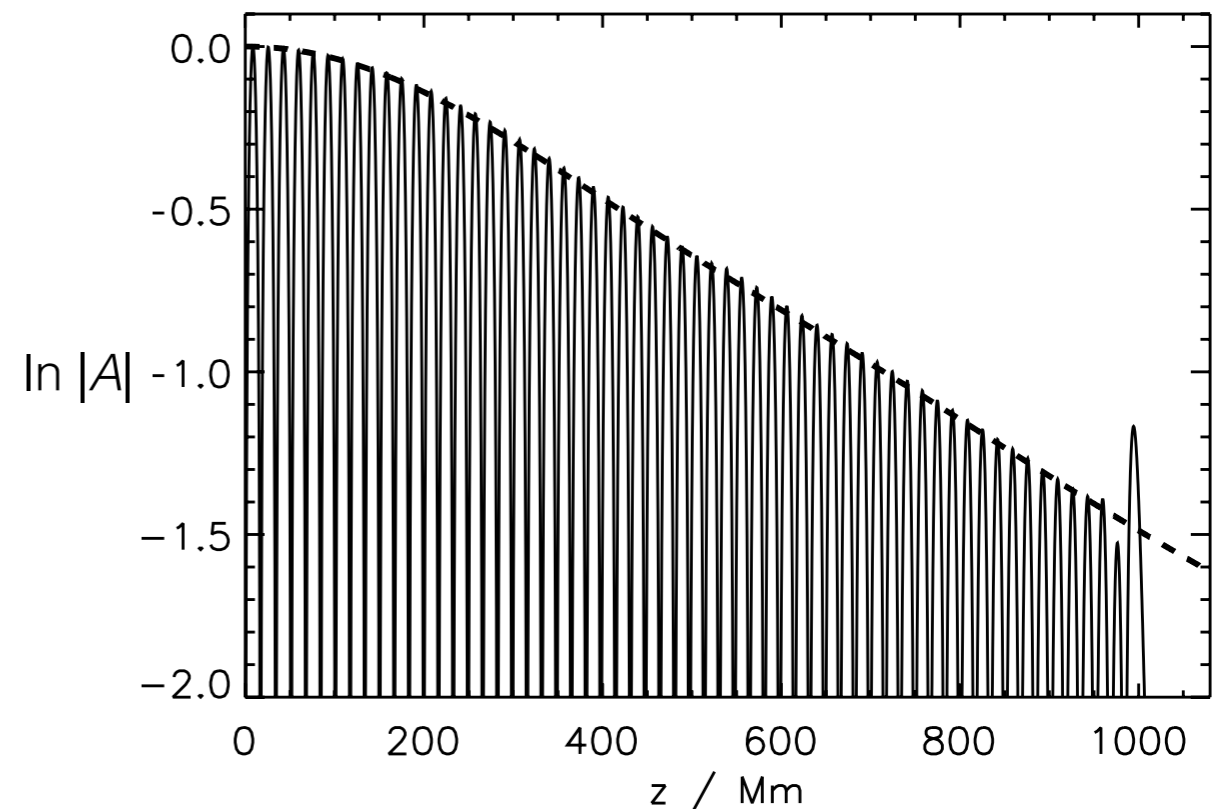
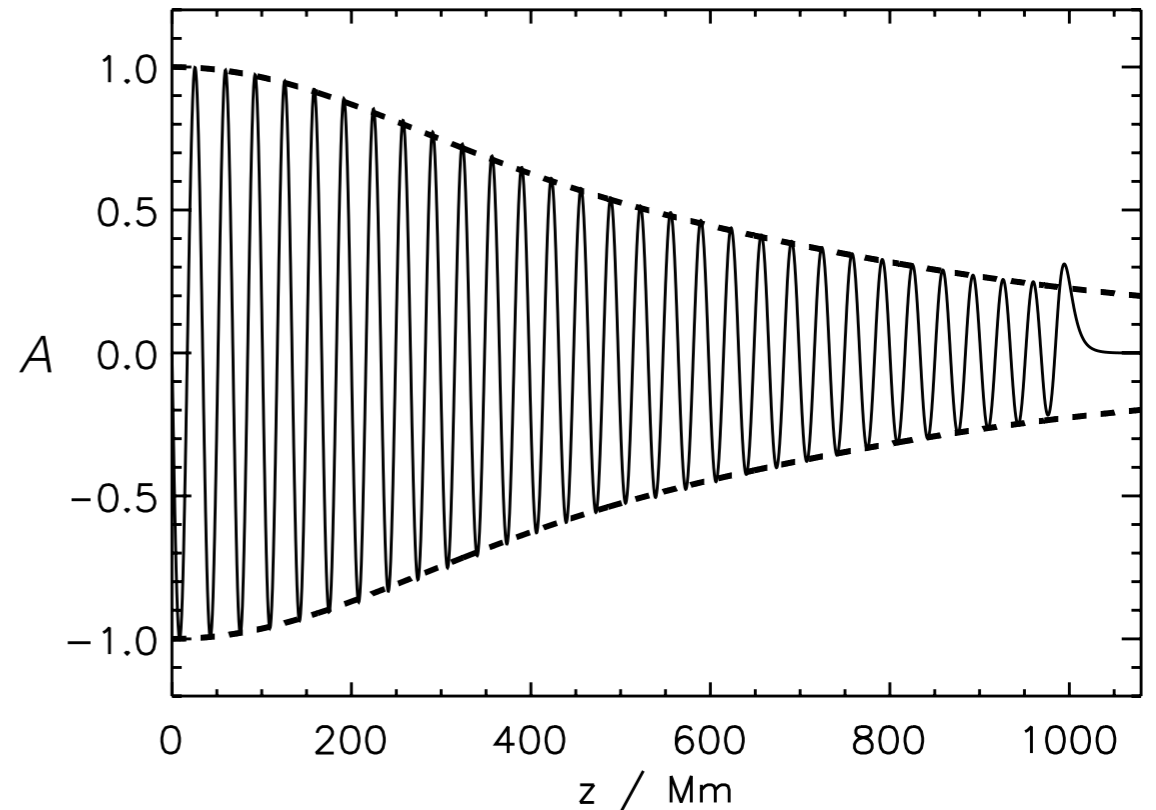


New theoretical damping envelope

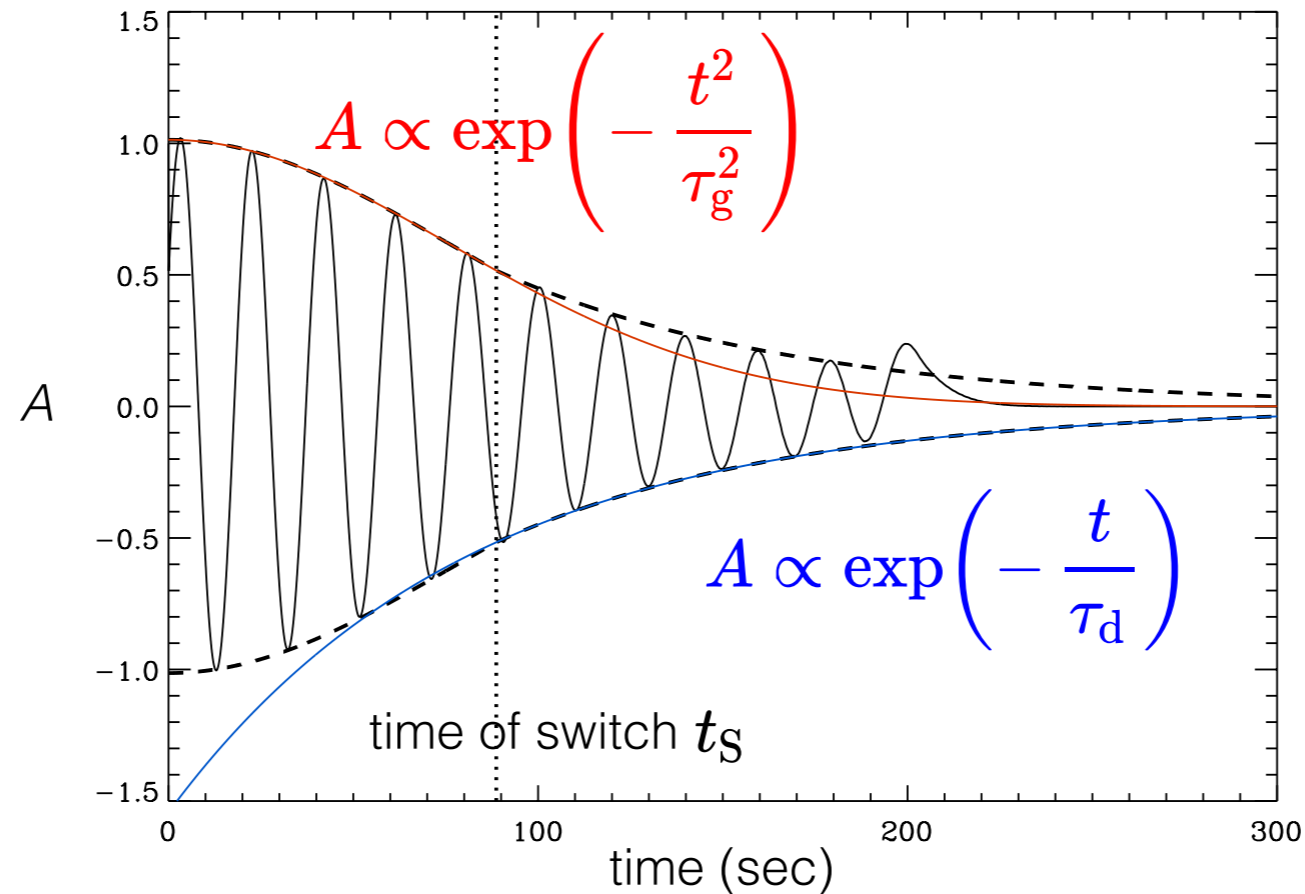
- New analytical solution (Hood et al. 2013) gives accurate envelope for all times:

$$\frac{d\tilde{\eta}}{dZ} = \frac{\epsilon}{2} \left\{ \frac{a(1 - \cos Z)}{Z} - \int_0^Z \tilde{\eta}(u) \frac{\sin(Z - u)}{Z - u} du \right\}$$

- Integro-differential equation allows precise testing of numerical simulations
- Gaussian profile describes initial state of the system
- Exponential damping profile describes asymptotic state of the system
- Inconvenient for seismological inversions for observations (solved numerically → slow)



Approximate damping envelope for seismology



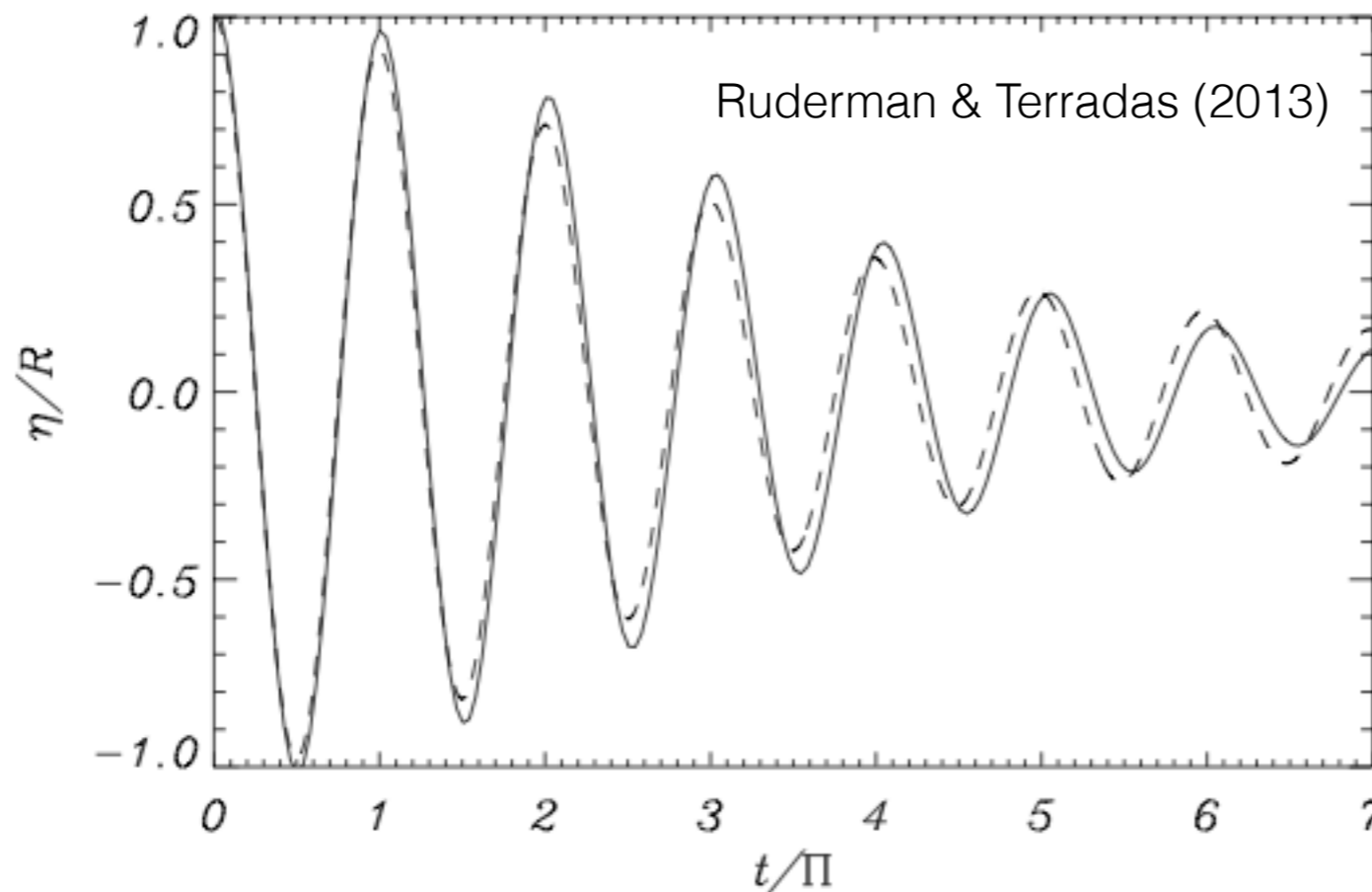
Numerical simulation of a kink mode showing **both** Gaussian and exponential damping regimes

$$t_s = \tau_g^2 / \tau_d$$

- Time of switch depends on density contrast ratio:
 - small density contrast (~ 2) \rightarrow slow switch \rightarrow mostly Gaussian envelope
 - large density contrast (~ 10) \rightarrow quick switch \rightarrow mostly exponential envelope
- If we detect both envelopes we have **two damping times** (well-posed inversion problem) \rightarrow unique solution for density contrast ratio and transition layer width

Gaussian damping regime for standing kink modes

- Same physical mechanism applies to standing modes



$$\frac{\tau_g}{P} = \frac{L_g}{\lambda} = \frac{2}{\pi\kappa\epsilon^{1/2}}$$

Pascoe et al. (2013)

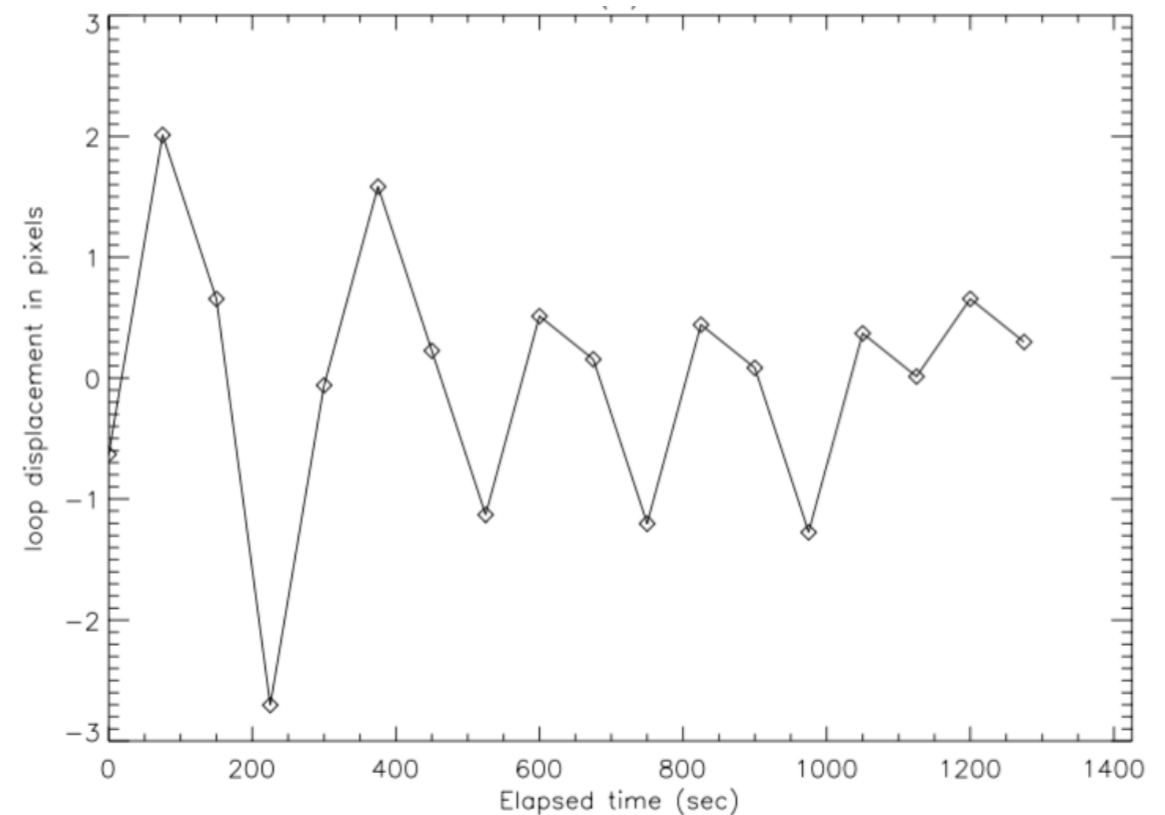
- For standing kink modes, a large density contrast is usually assumed (typically 10, e.g., Nakariakov et al. 1999; Ruderman & Roberts 2002; Goossens et al. 2002)
- This density contrast is consistent with an exponential damping profile (Gaussian profile for ~ 1.2 oscillations only)

Evidence of Gaussian damping regime from TRACE

- De Moortel et al. (2002) and Ireland and De Moortel (2002) analyse standing kink oscillations
- Shape of damping envelope taken to be a fitted parameter

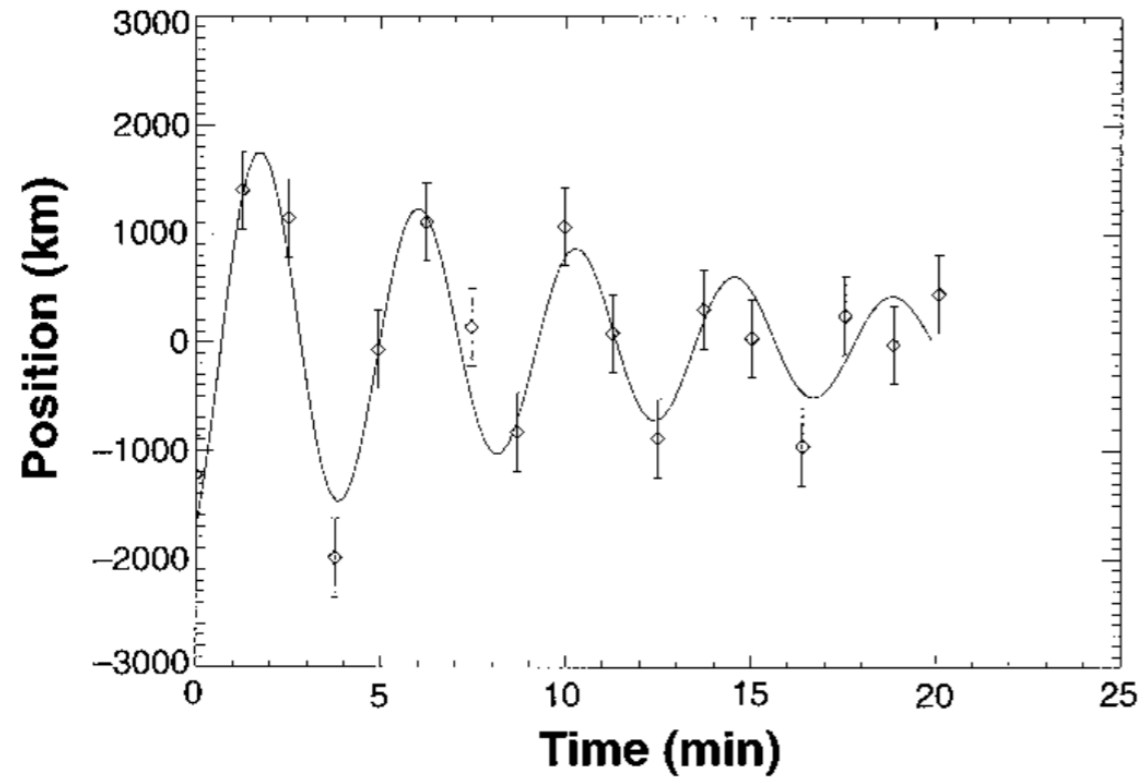
$$\propto \exp\left(-kt^N\right)$$

- 14 July 1998: $N \approx 2.0 \pm 1.2$
- 25 October 1999: $N \approx 0.43 \pm 1.17$
- 21 March 2001: $N \approx 2.0 \pm 1.3$



- Large errors due to low temporal resolution and noise for TRACE

Strongly damped kink oscillations

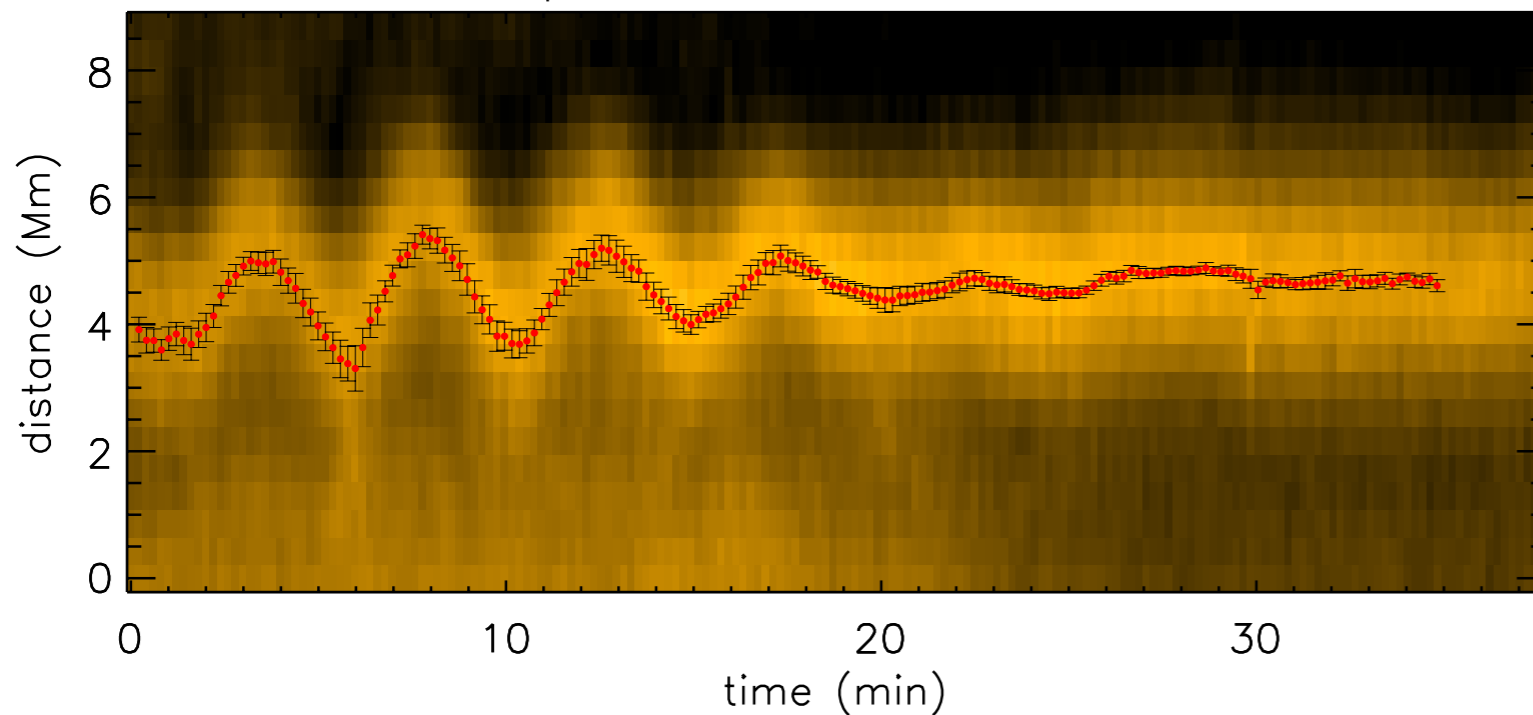


0.5 arcsec/pixel

TRACE

~75 second
cadence

Event 43 Loop 4 7-Jan-2013 06:38:11.34

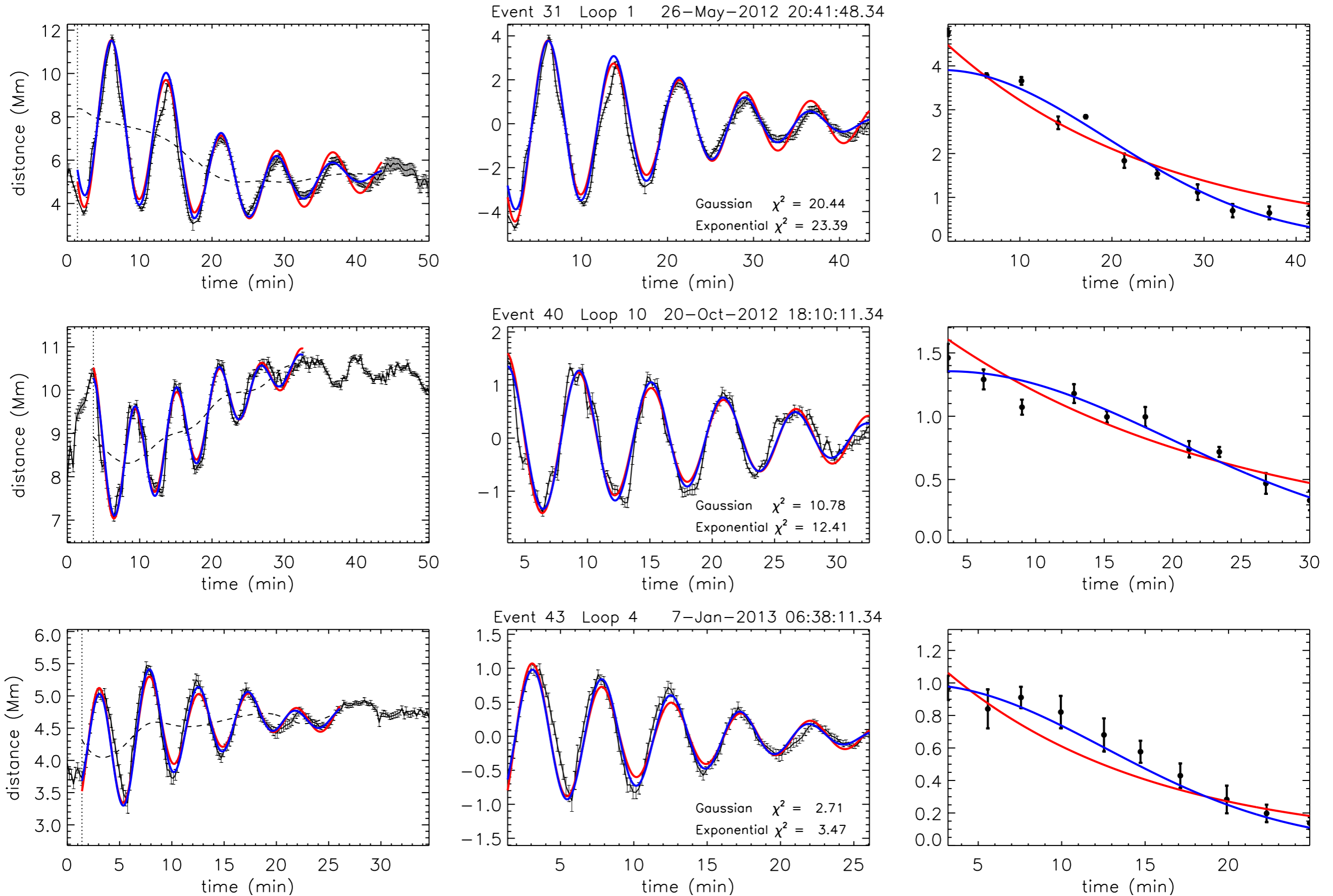


0.6 arcsec/pixel
(full disk)

SDO/AIA

12 second
cadence

Gaussian examples



Evidence of Gaussian damping regime from SDO

- Morton & Mooroogen (2016, A&A, in press) — statistical analysis of one of the previous loops supports Gaussian damping profile

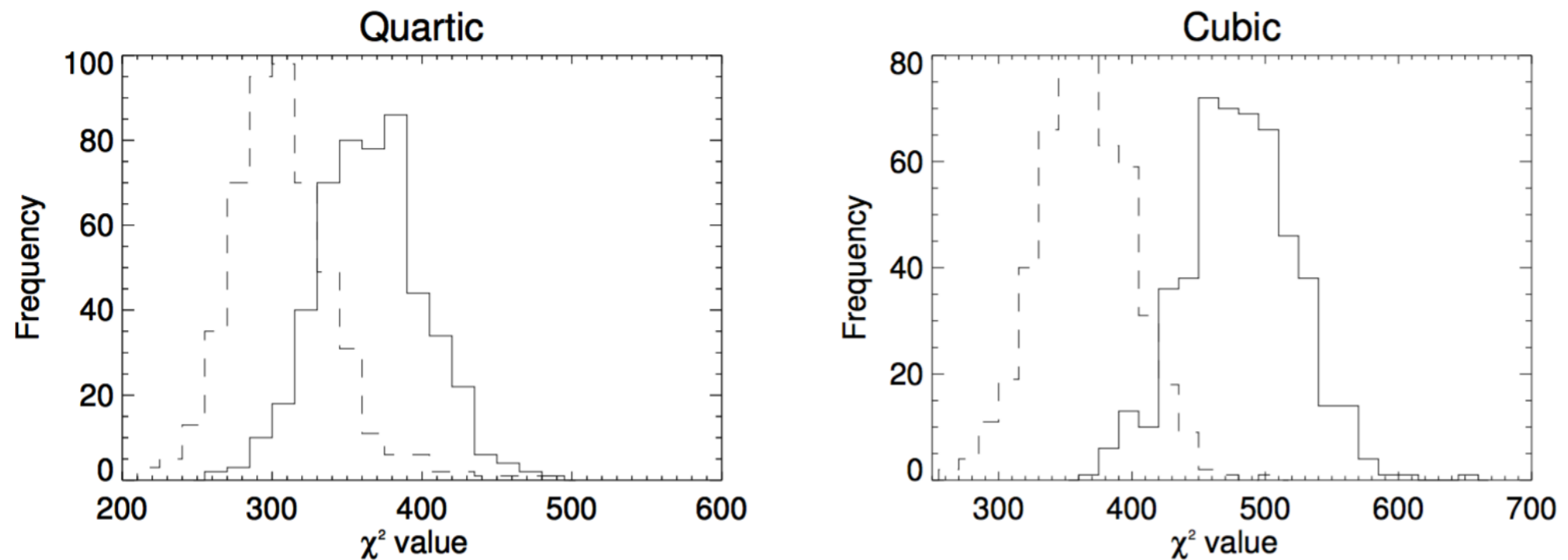
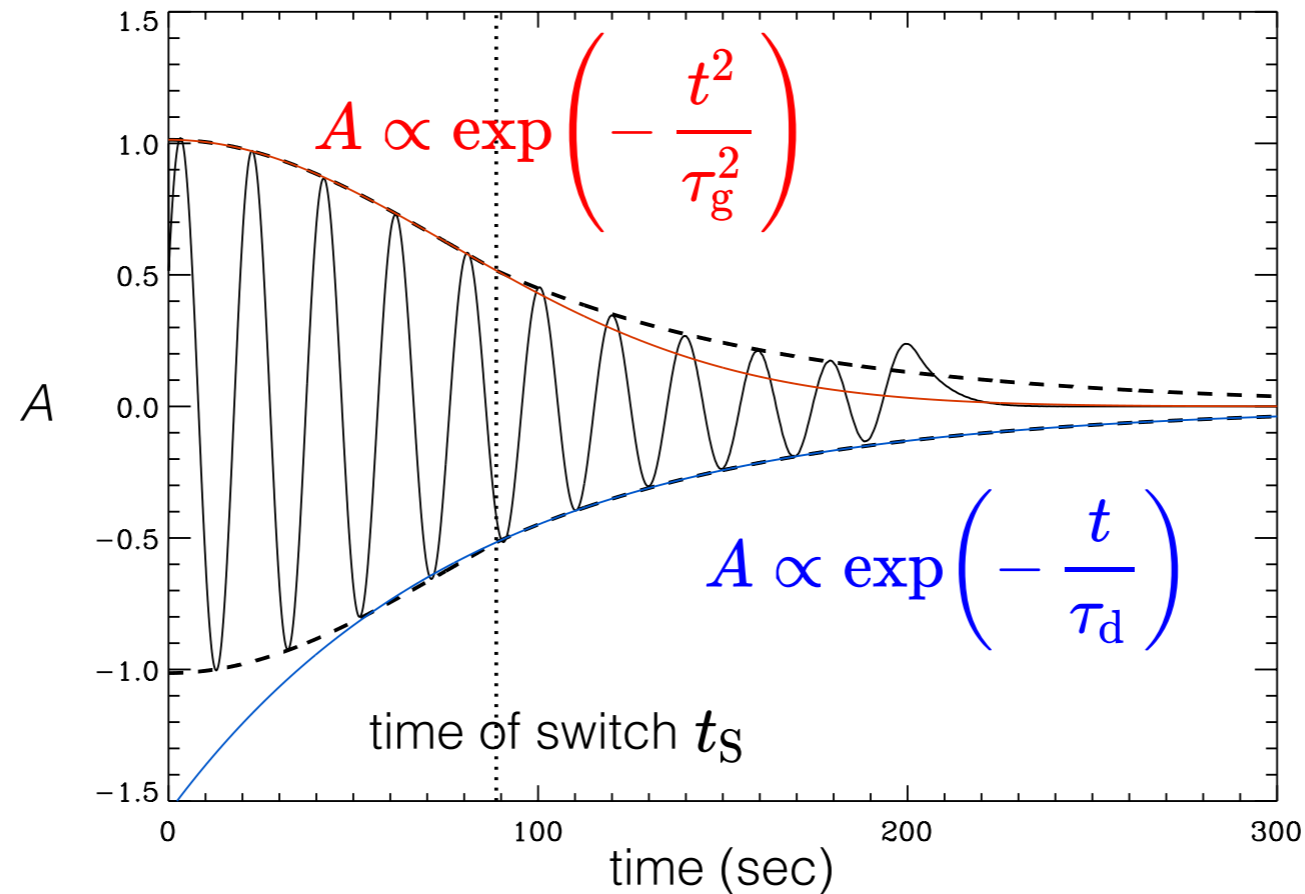


Fig. 3. Distribution of the χ^2 statistic for the quartic and cubic models. The solid lines show the distribution for the exponential damping, while the dashed line shows the distribution for the Gaussian damping.

- Model comparison:
 - Kolmogorov—Smirnov test (Morton & Mooroogen 2016)
 - Bayesian inference (e.g. Arregui et al. 2013, ApJ, 765, L23)

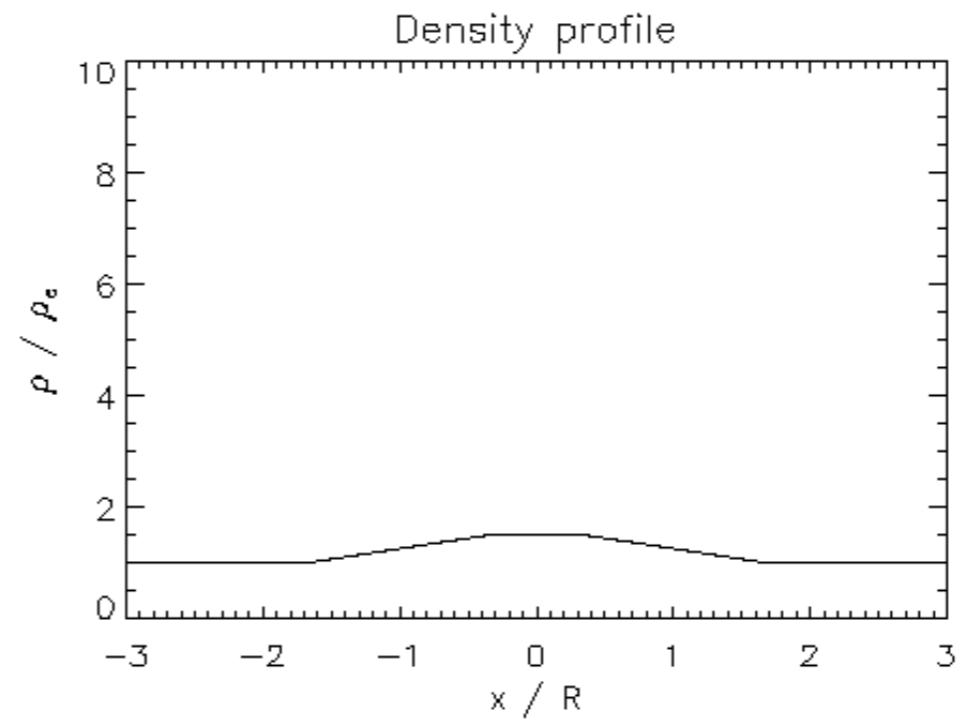
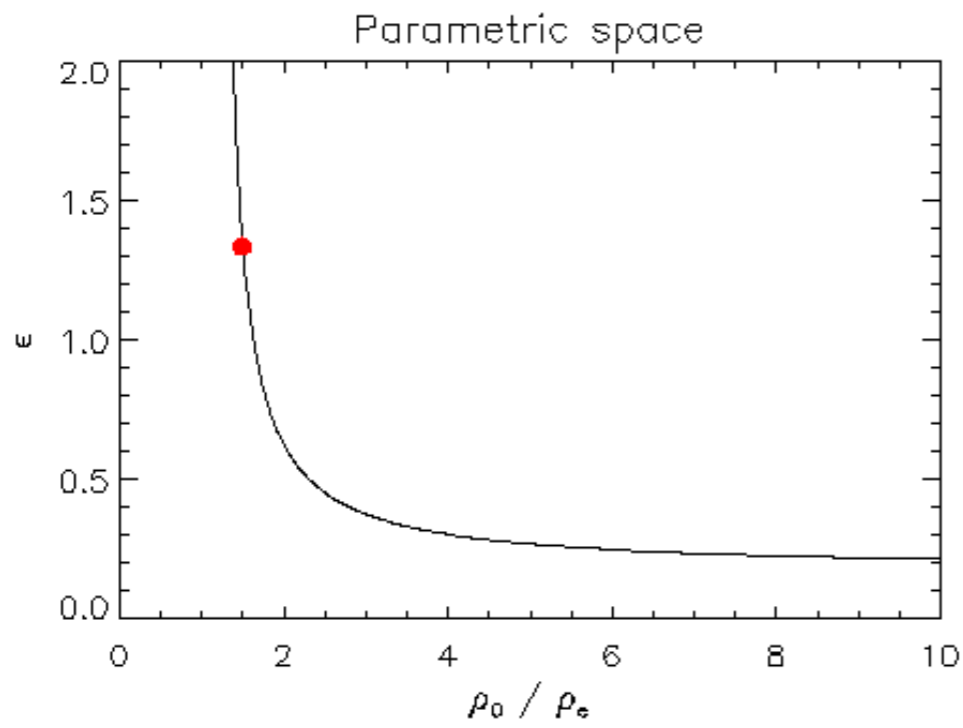
Approximate damping envelope for seismology



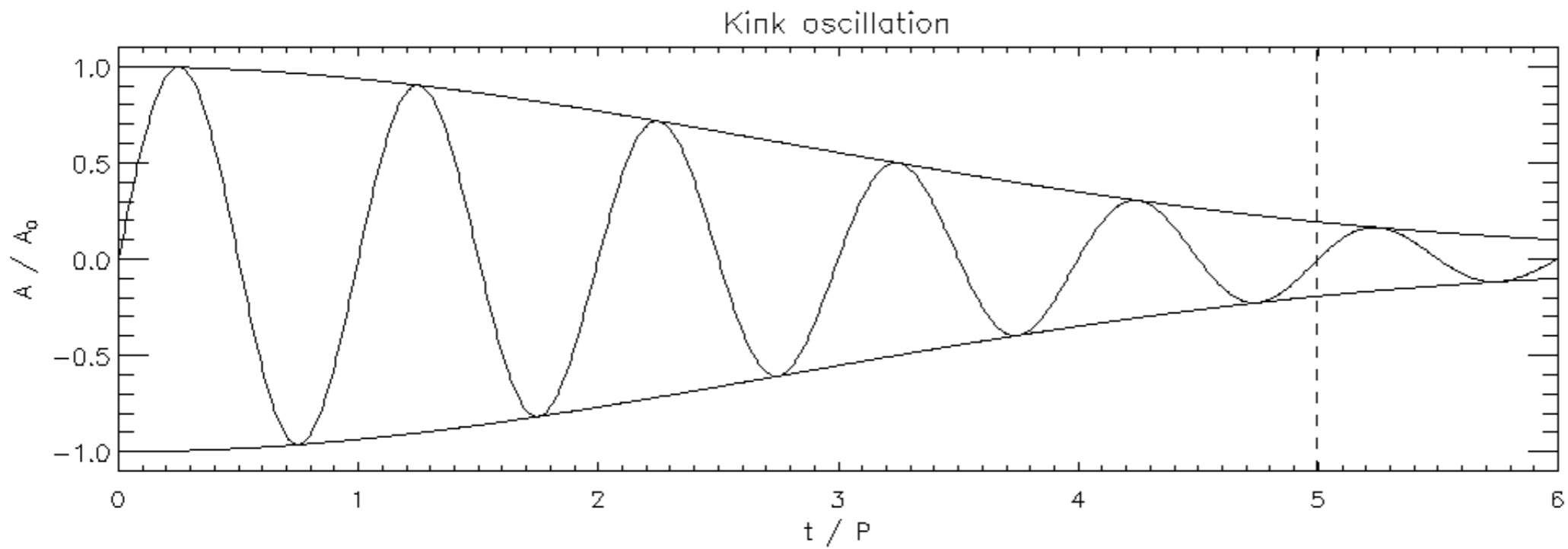
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Animation shows some of the transverse loop profiles which give the same damping rate (of 90% attenuation after $6P$)

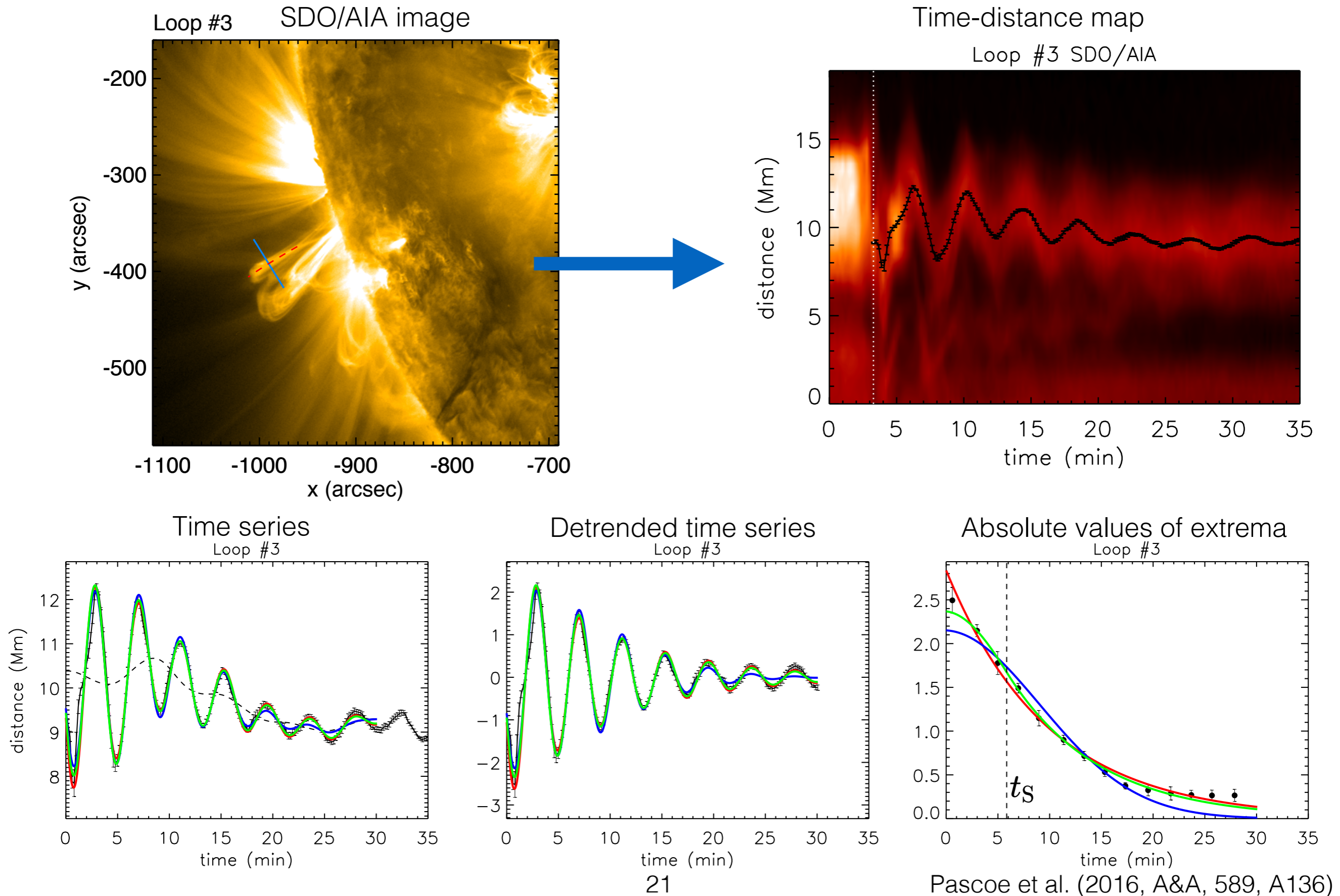


For each structure the **shape** of the kink mode damping envelope is different

- New method uses shape of damping envelope as well as damping rate to obtain structure information
- Envelope shape is characterised by the switch time between Gaussian and exponential damping regimes:

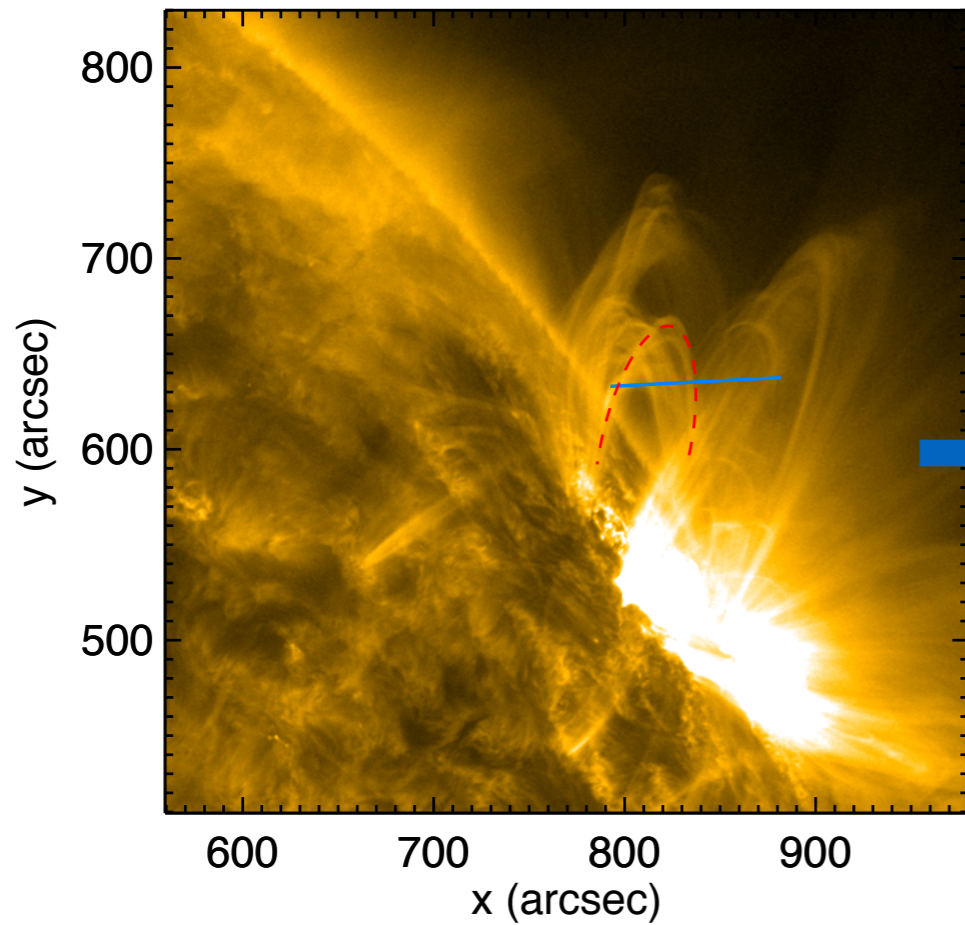
$$t_s = \tau_g^2 / \tau_d$$

General damping envelope fit to SDO data

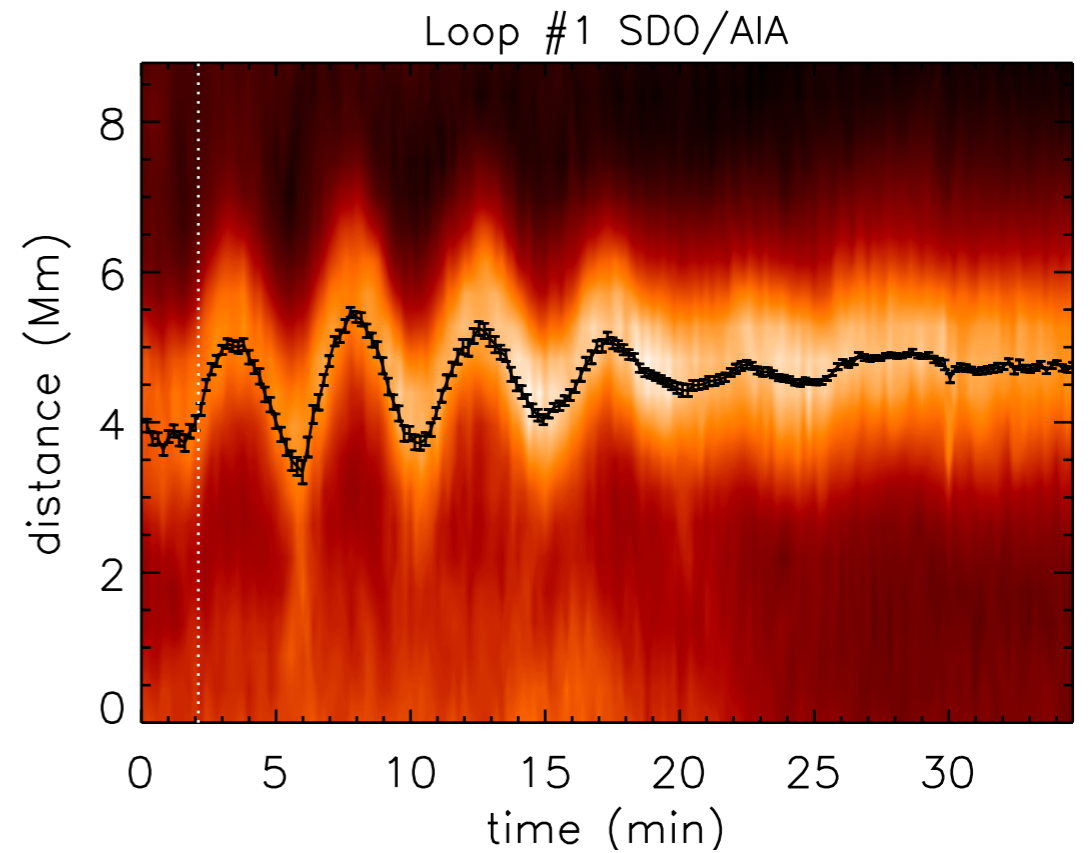


General damping envelope fit to SDO data

Loop #1 SDO/AIA image

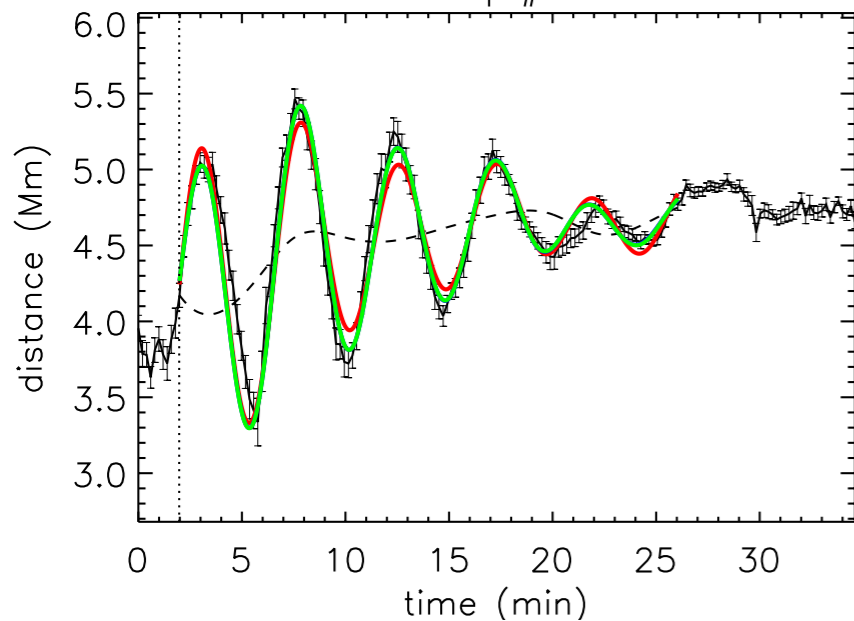


Time-distance map



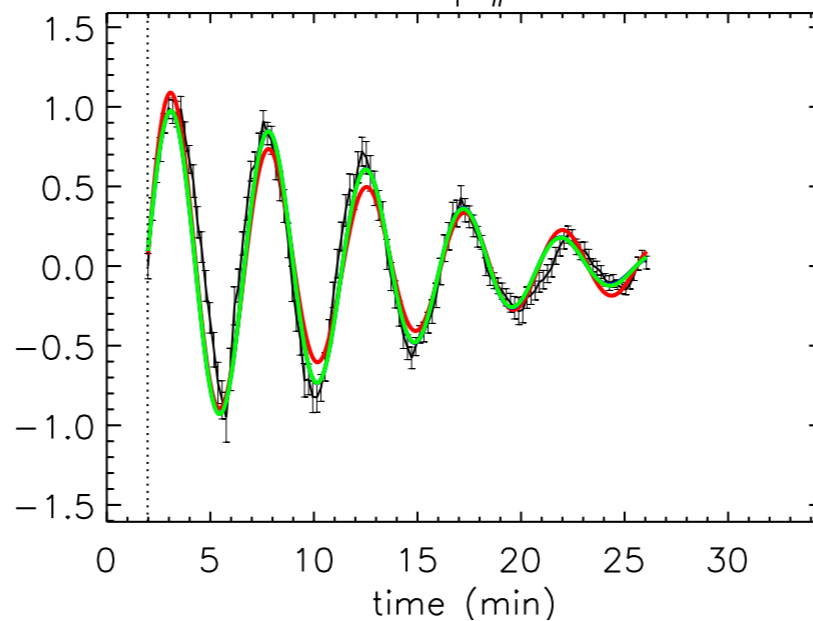
Time series

Loop #1



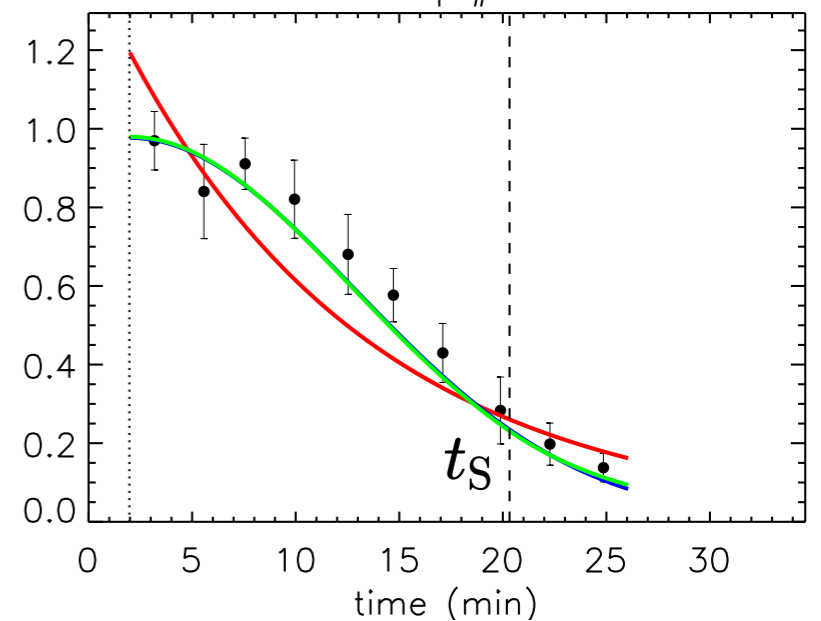
Detrended time series

Loop #1



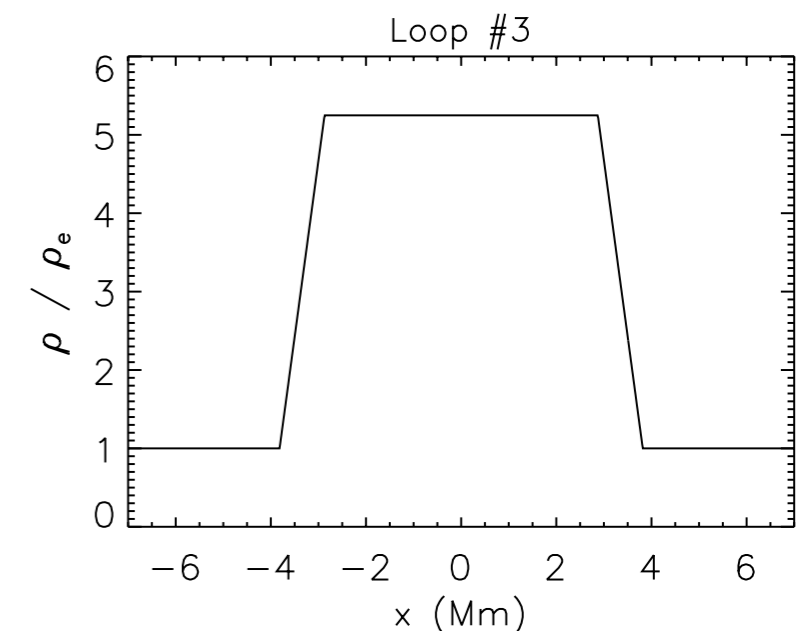
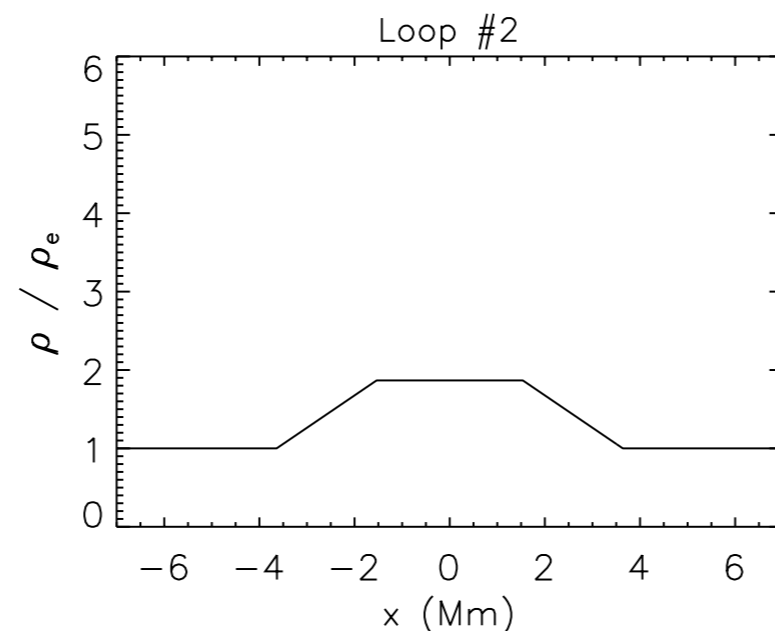
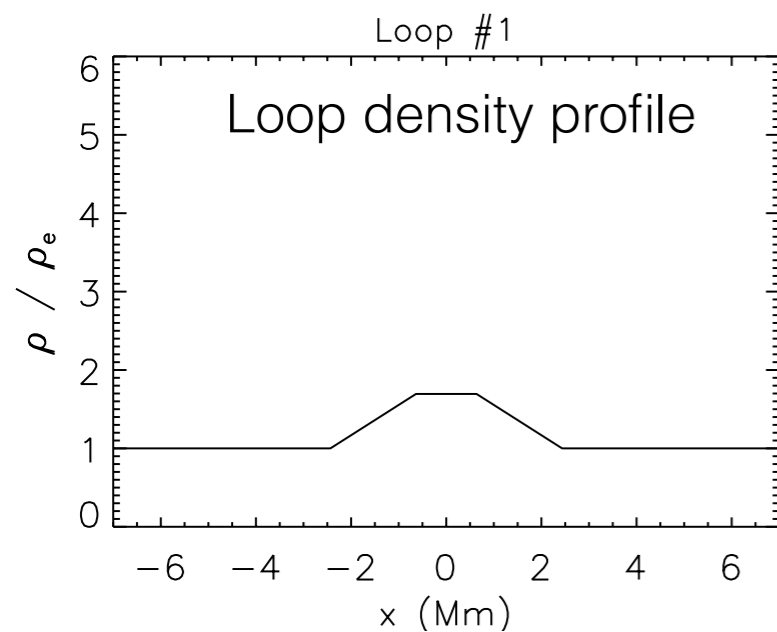
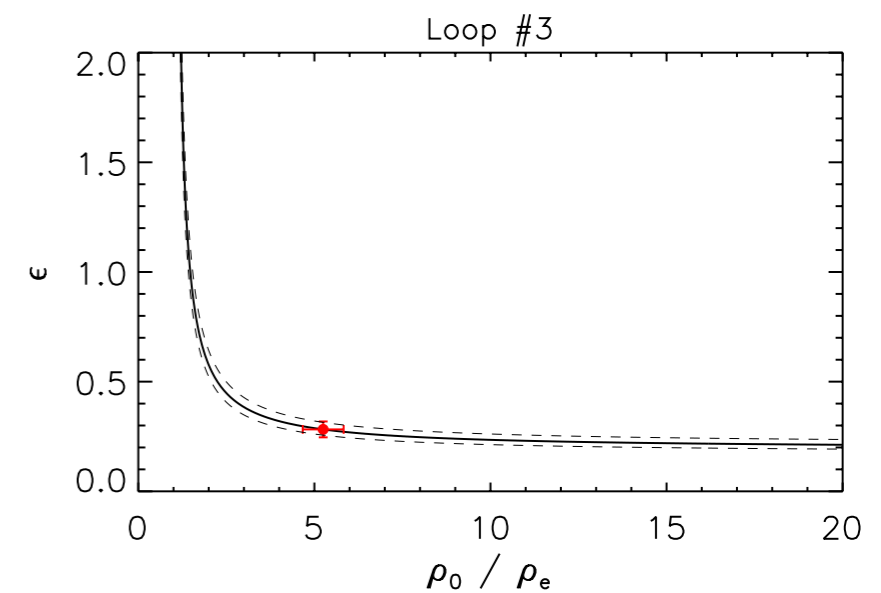
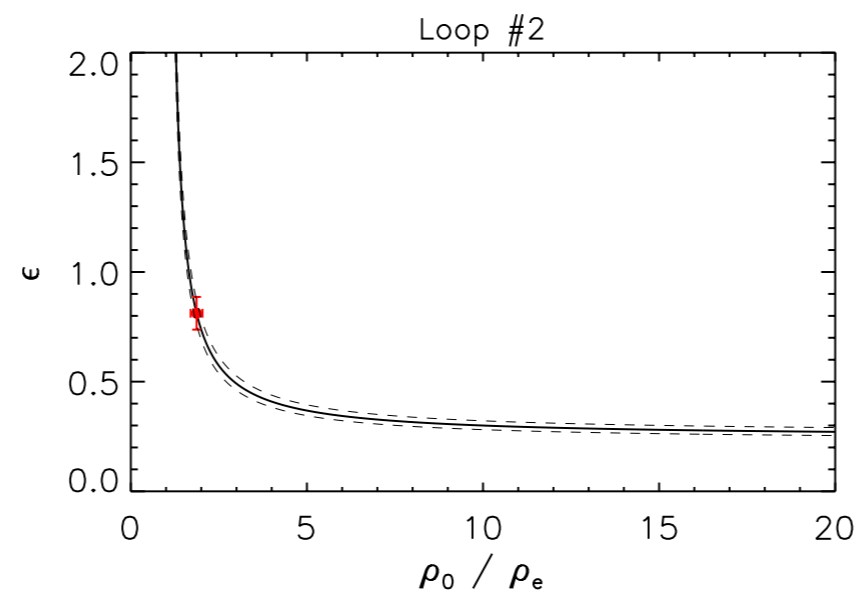
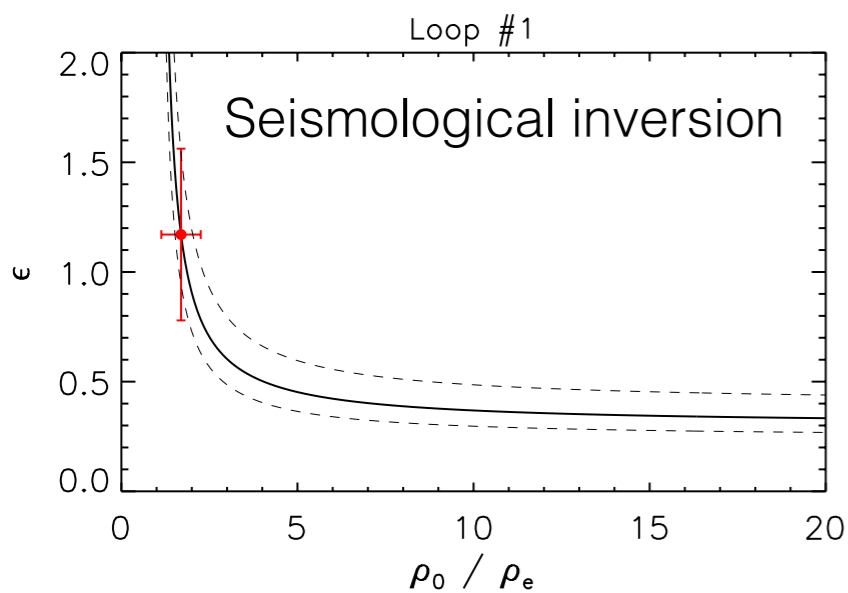
Absolute values of extrema

Loop #1

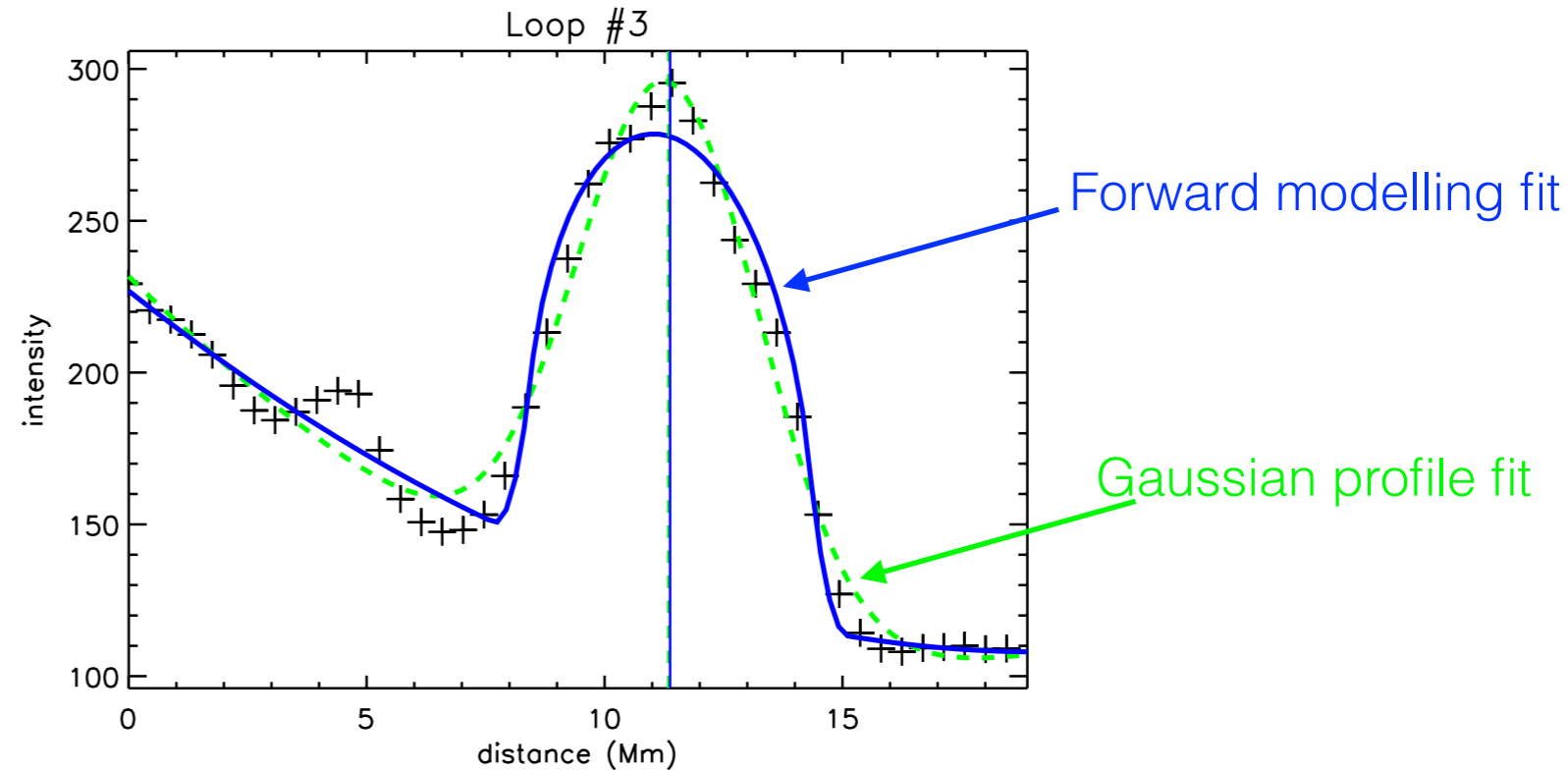
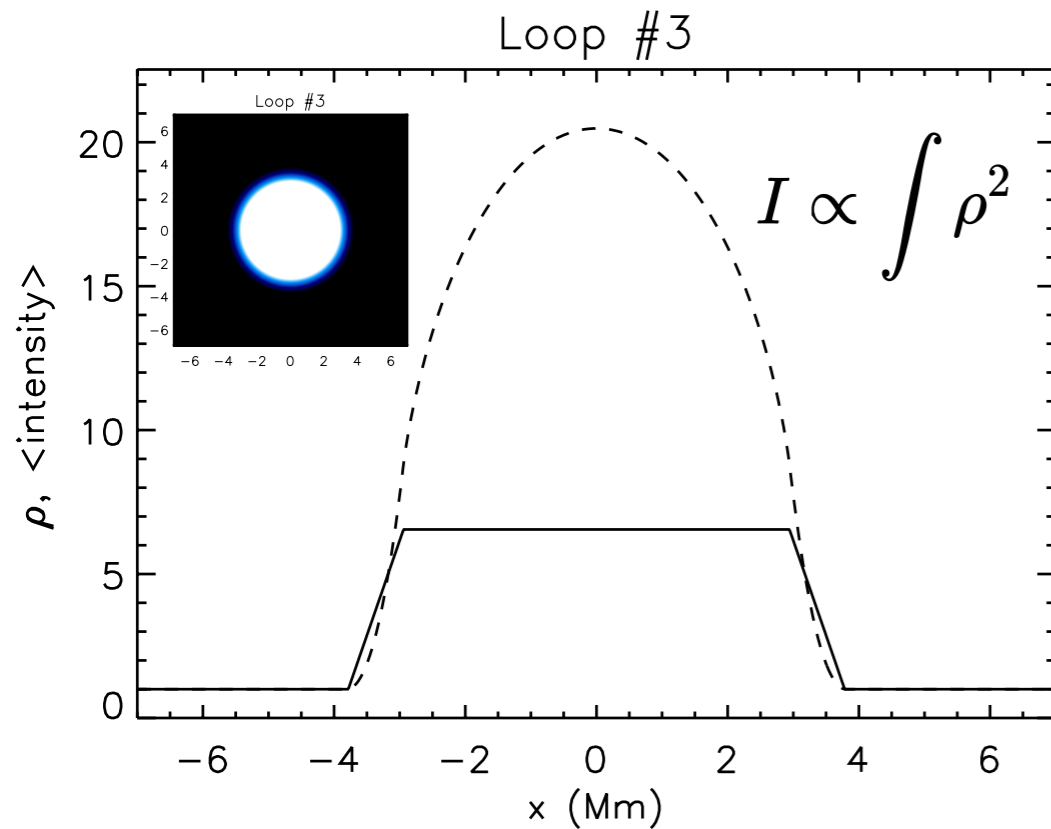


Seismology of coronal loop transverse structure

- Two damping regimes (Gaussian and exponential) used to calculate the transverse loop structure
- Three loops from catalogue of Goddard et al. (2016) analysed in detail:



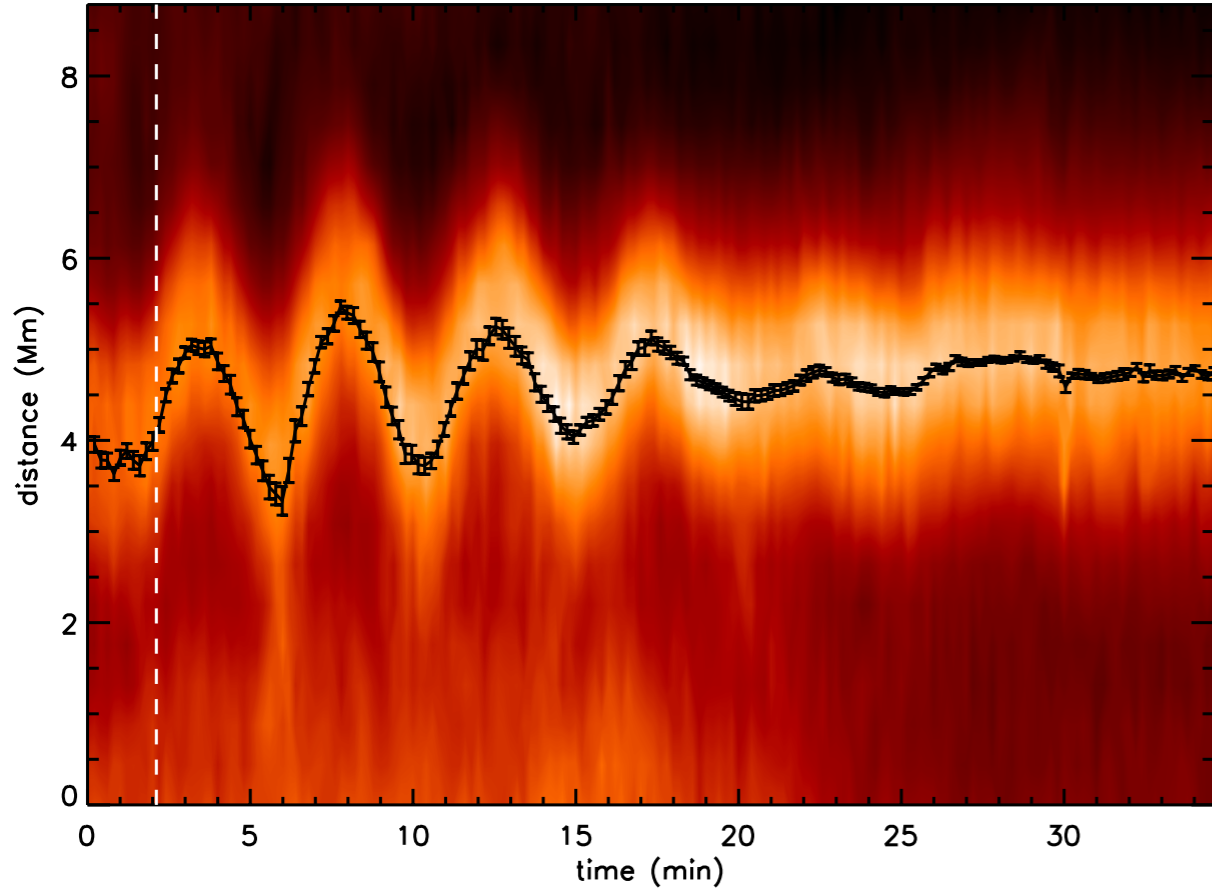
Forward modelling intensity profile



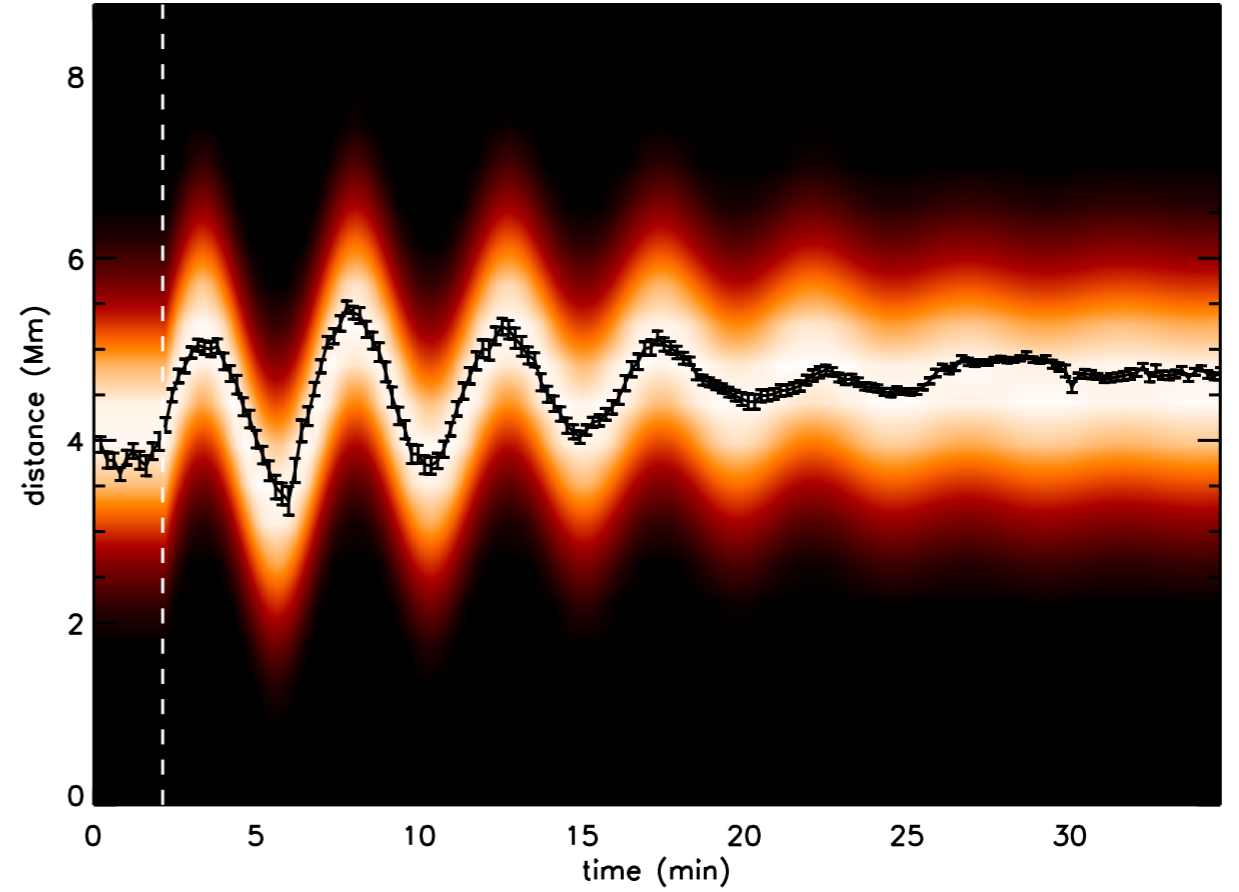
- Damping profile seismology gives $\epsilon = \frac{l}{R}$
- Forward modelling intensity profile and fitting to data gives l and R separately
- This method ensures same definition of l and R used in forward modelling as damping model

Forward modelling TD maps

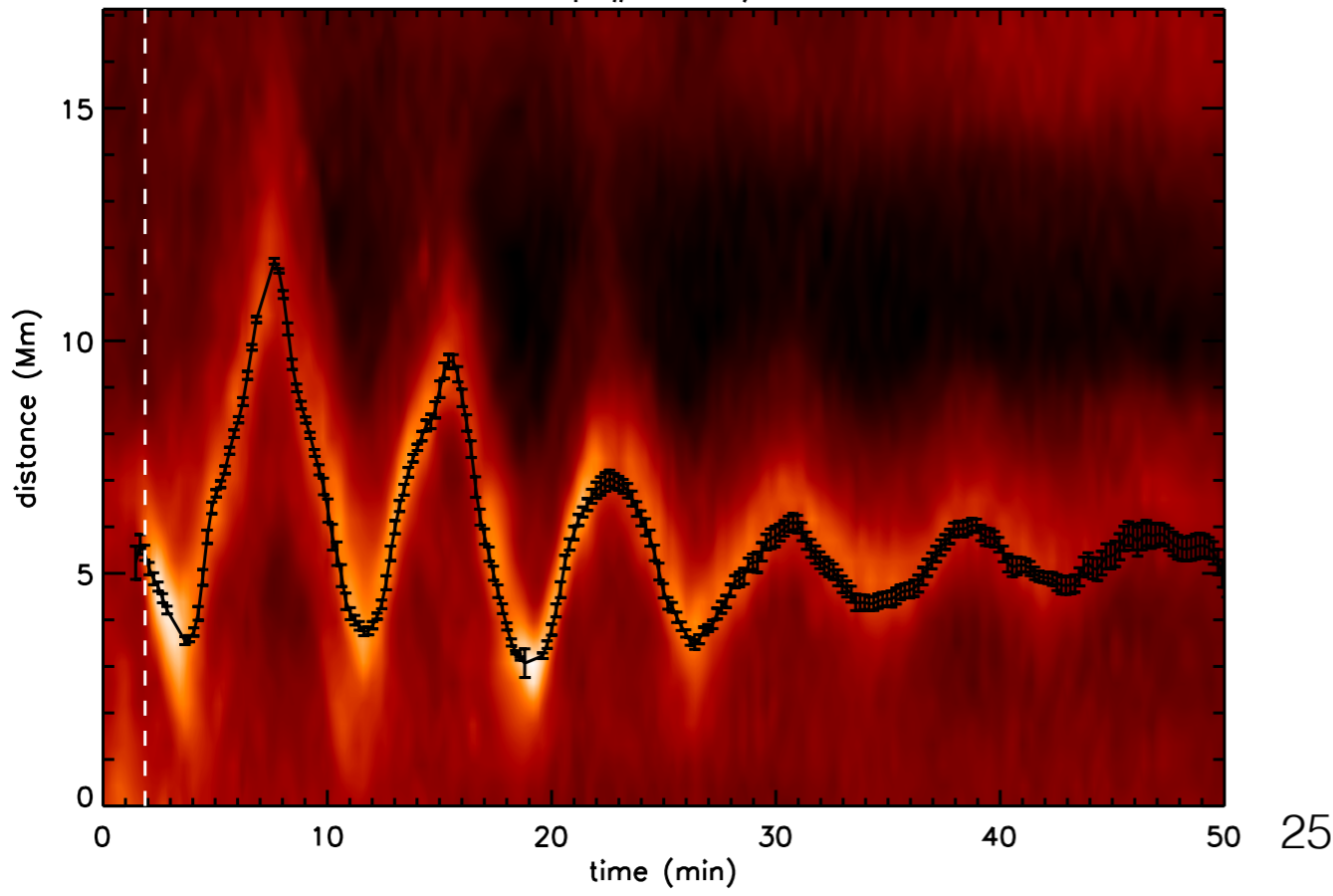
Loop #1 SDO/AIA



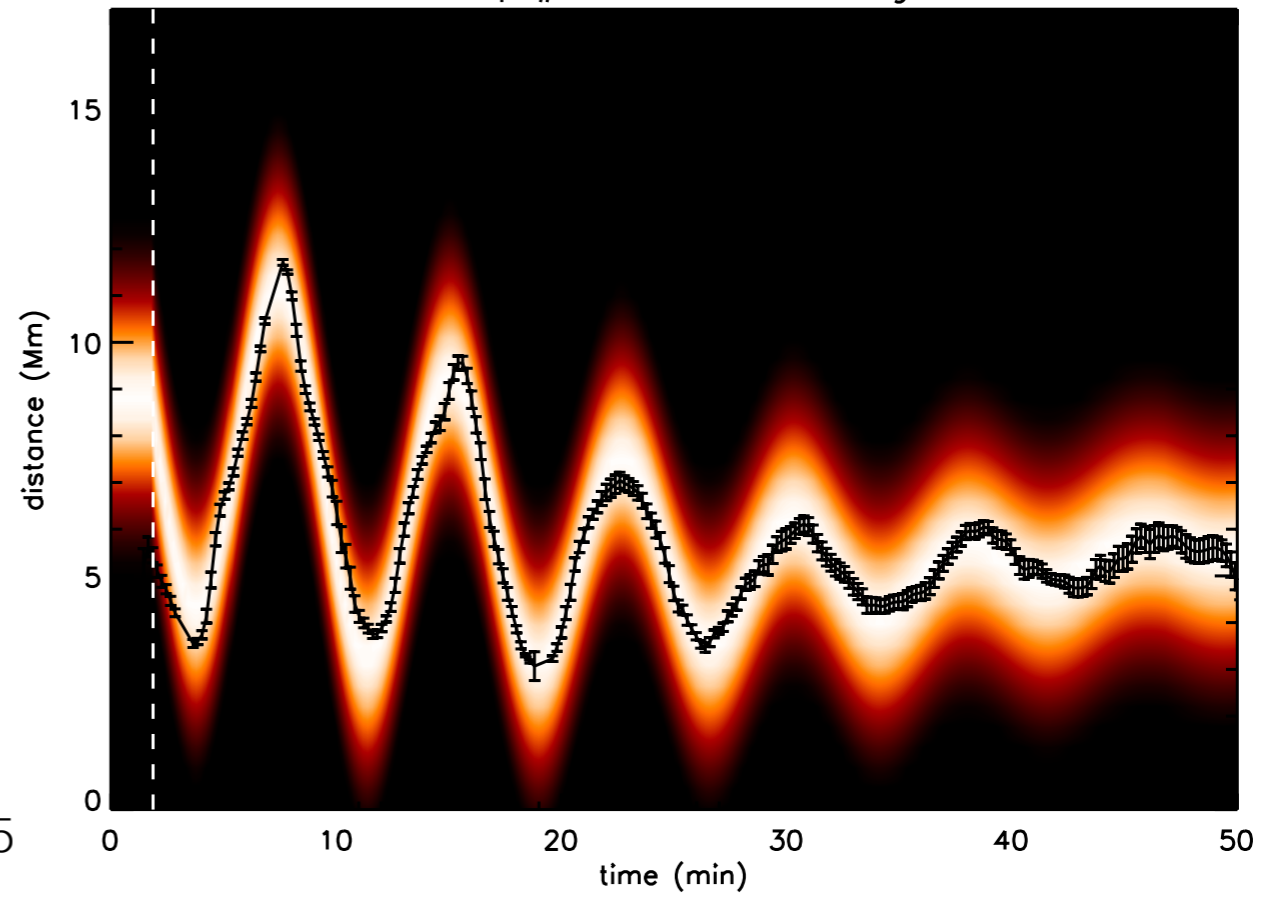
Loop #1 Forward Modelling



Loop #2 SDO/AIA



Loop #2 Forward Modelling



General spatial damping profile seismology

- Improved estimates of Alfvén speed by calculating actual density contrast ratio:

$$C_k = 2L/P$$

$$C_{A0} = C_k / \sqrt{2/(1 + \rho_e/\rho_0)}$$

$$C_{Ae} = C_{A0} \sqrt{\rho_0/\rho_e}$$

$$B_0 = C_{A0} \sqrt{\mu_0 \rho_0} = C_{A0} \sqrt{\mu_0 \bar{\mu} m_p n_0}$$

$$\bar{\mu} = 1.27$$

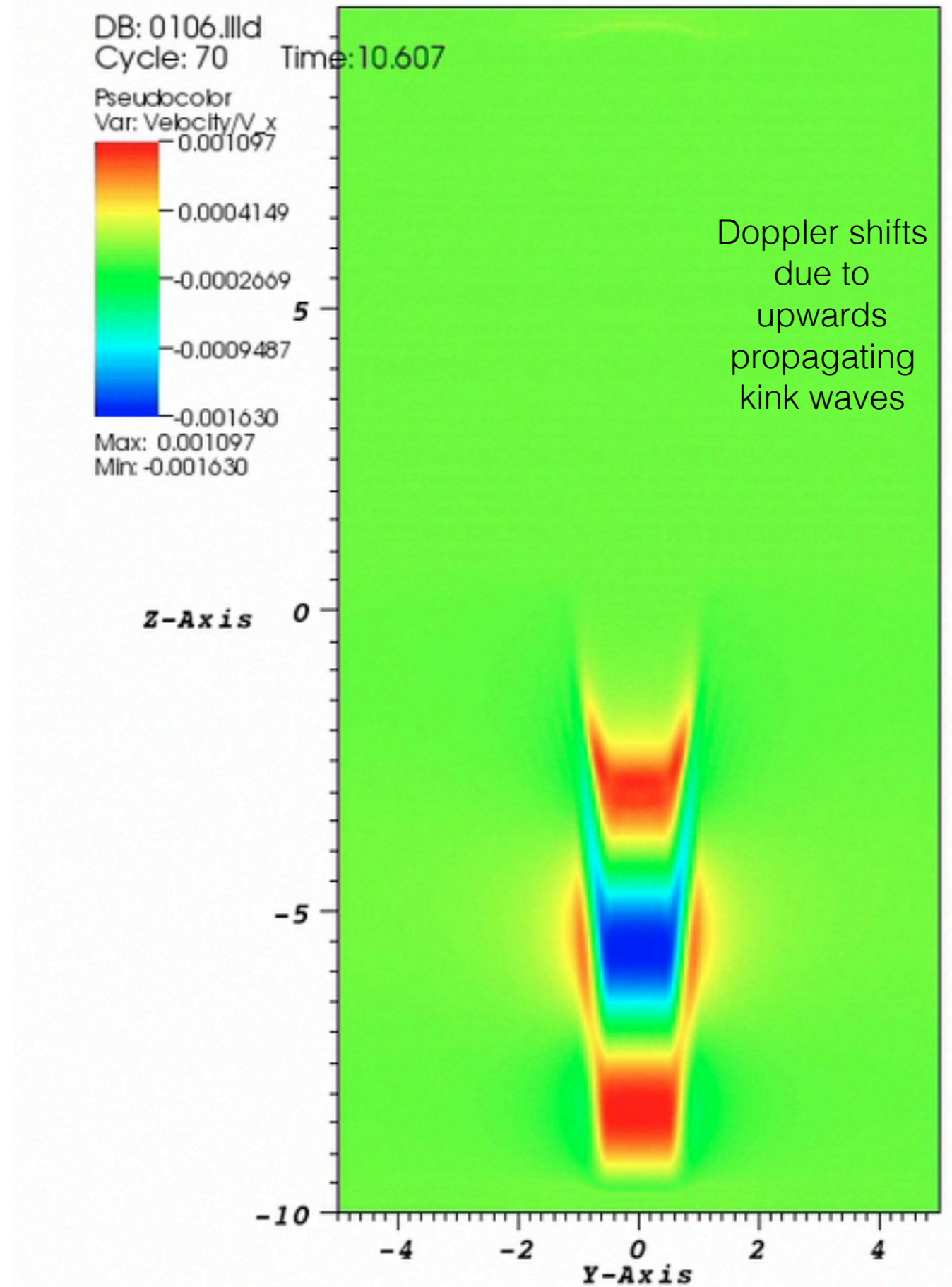
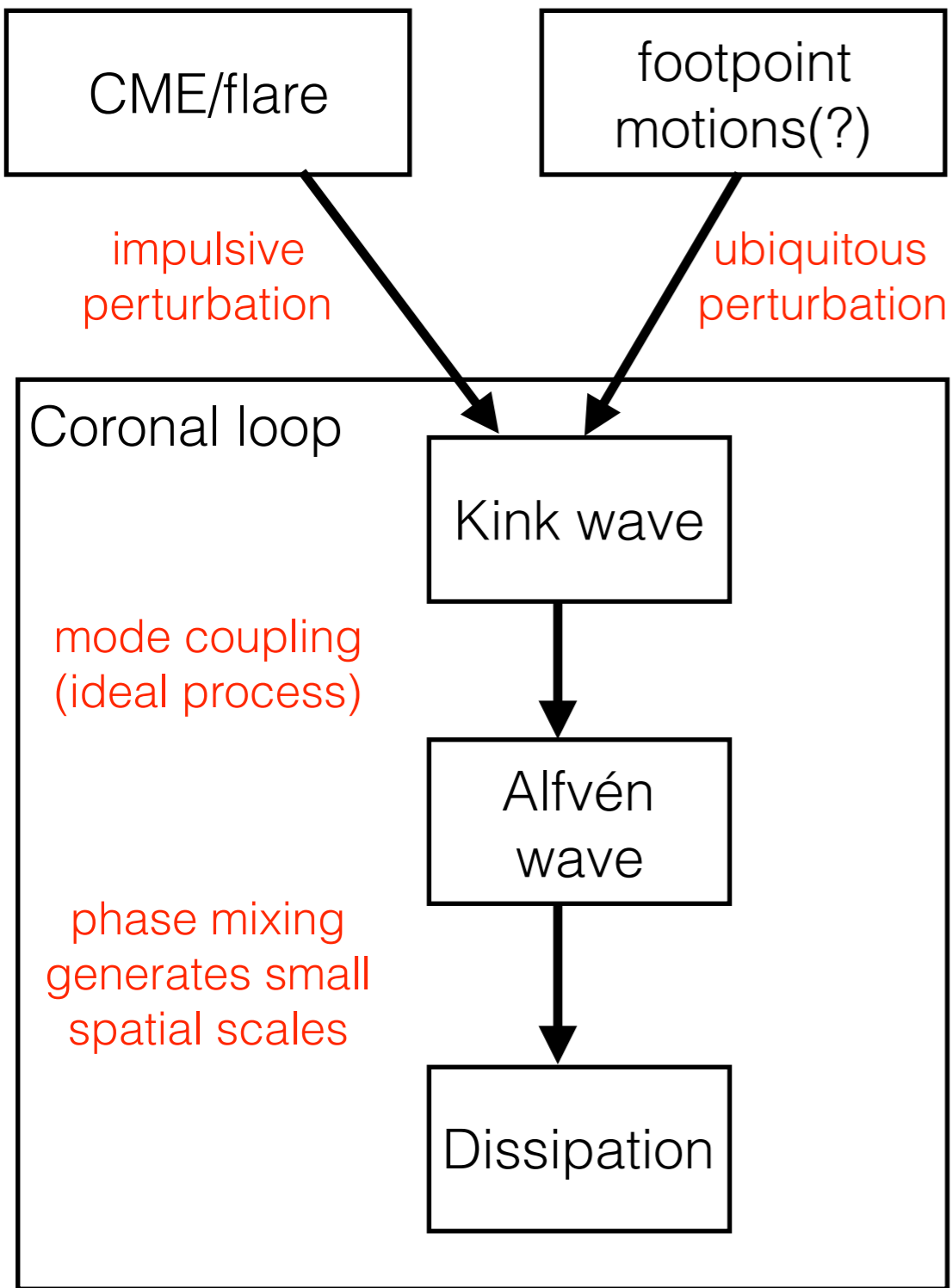
$$m_p = 1.6726 \times 10^{-27} \text{ kg}$$

$$n_e = 10^{15} \text{ m}^{-3}$$

Loop no.	A (Mm)	ϕ (rad)	P (min)	τ_g (min)	τ_d (min)	t_s (min)	c_0	c_1	c_2
Loop #1	0.98 ± 0.02	0.10 ± 0.03	4.73 ± 0.01	10.79 ± 0.27	6.35 ± 1.50	18.32 ± 4.27	4.05	0.075	-0.0022
Loop #2	-3.97 ± 0.02	0.00 ± 0.01	7.64 ± 0.01	18.49 ± 0.15	11.24 ± 1.07	30.40 ± 2.88	9.04	-0.237	0.0035
Loop #3	-2.36 ± 0.03	0.00 ± 0.01	4.18 ± 0.00	7.24 ± 0.27	9.21 ± 0.93	5.69 ± 0.49	10.69	-0.079	0.0009

Loop no.	ρ_0/ρ_e	ϵ	R (Mm)	l (Mm)	C_k (Mm/s)	C_{A0} (Mm/s)	C_{Ae} (Mm/s)	B_0 (G)
Loop #1	1.70 ± 0.56	1.17 ± 0.39	1.53 ± 0.13	1.79 ± 0.61	1.56 ± 0.22	1.40 ± 0.38	1.82 ± 0.65	9.39 ± 3.36
Loop #2	1.67 ± 0.22	1.10 ± 0.15	2.44 ± 0.45	2.67 ± 0.61	0.71 ± 0.14	0.63 ± 0.13	0.82 ± 0.19	4.22 ± 0.99
Loop #3	6.55 ± 0.80	0.25 ± 0.03	3.36 ± 0.28	0.84 ± 0.13	1.87 ± 0.25	1.42 ± 0.22	3.62 ± 0.65	18.72 ± 3.38

Phase mixing

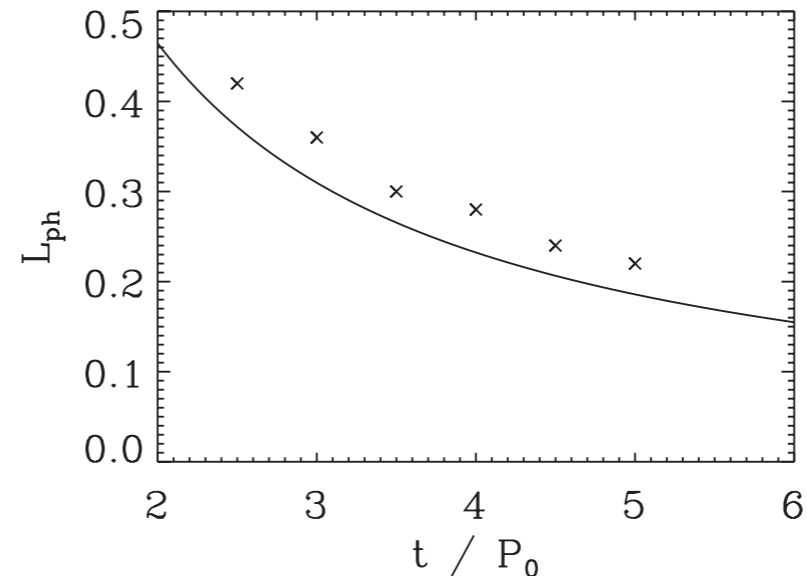


Phase mixing

- Alfvén wave generated inside inhomogeneous medium with continuous variation in local Alfvén speed — large gradients generated (e.g. Heyvaerts & Priest 1983; Cally 1991; Hood et al. 2005; Soler & Terradas 2015)
- Seismological method for determining density profile allows us to estimate phase mixing timescale

- Simulations: Mann et al. (1995):

$$L_{\text{ph}} = \frac{2\pi}{\omega'_A t}$$



Pascoe et al. (2010)

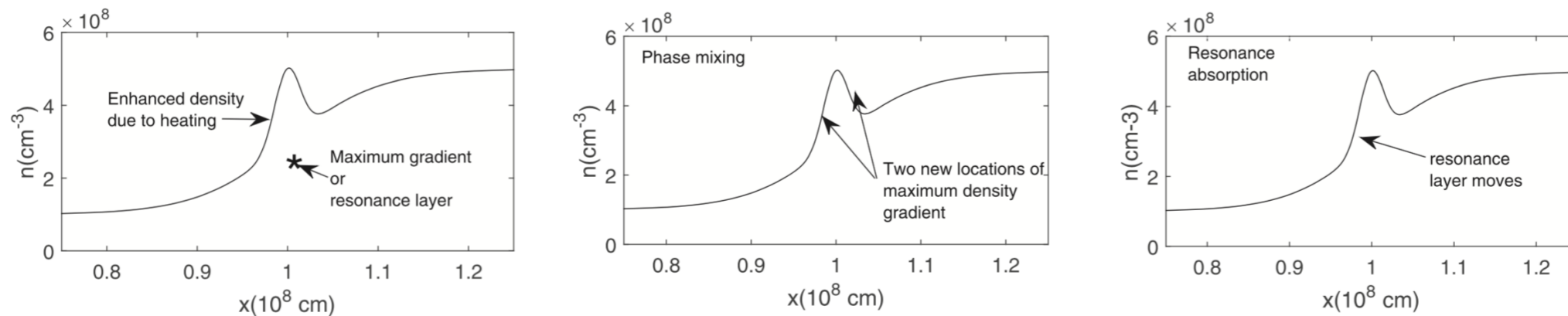
- Observations: Based on Mann & Wright (1995):

$$\tau_A = \frac{\epsilon L}{\pi(C_{Ae} - C_{A0})} \lesssim 200 \text{ seconds}$$

Pascoe et al. (2016)

Heating

- See next talk by Paolo Pagano for more details of heating
 - simulations including effects of resistivity, thermal conduction, radiative losses
- Cargill et al. (2016, ApJ, 823, 31) — wave damping creates fine structure



- Magyar & Van Doorselaere (2016, ApJ, 823, 82) — waves destroy fine structure

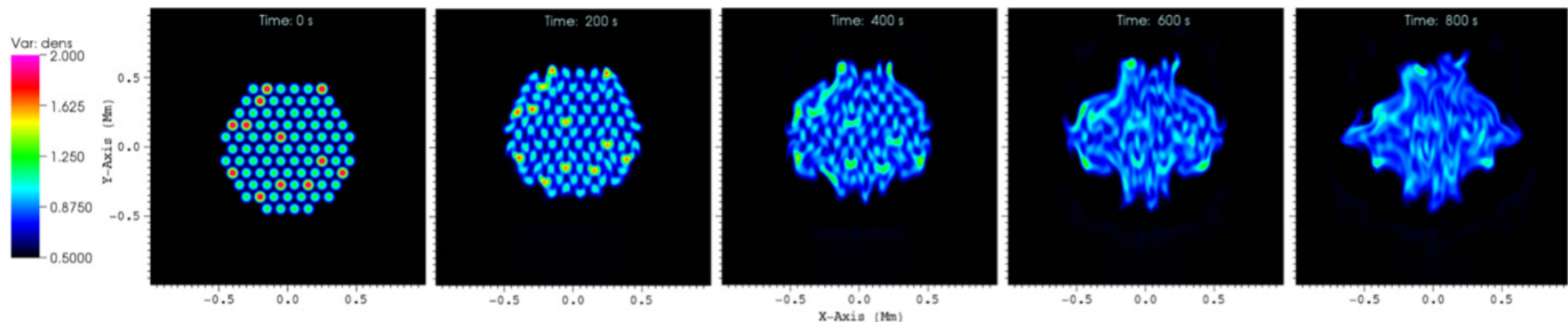
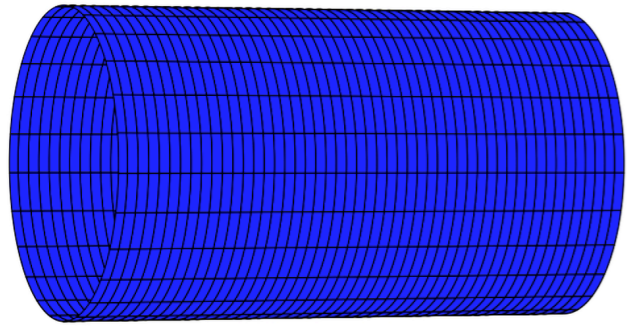


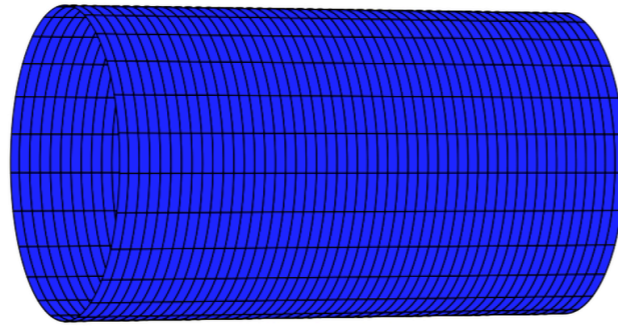
Figure 2. Plots of the loop cross-section density, at different times as indicated at the top of each panel. The slices are made at $z = 40$ Mm. The color bar is common for the plots, in units of $10^{-12} \text{ kg m}^{-3}$.

- (Mode coupling does not require symmetrical loops e.g. Terradas et al. 2008; Pascoe et al. 2011)

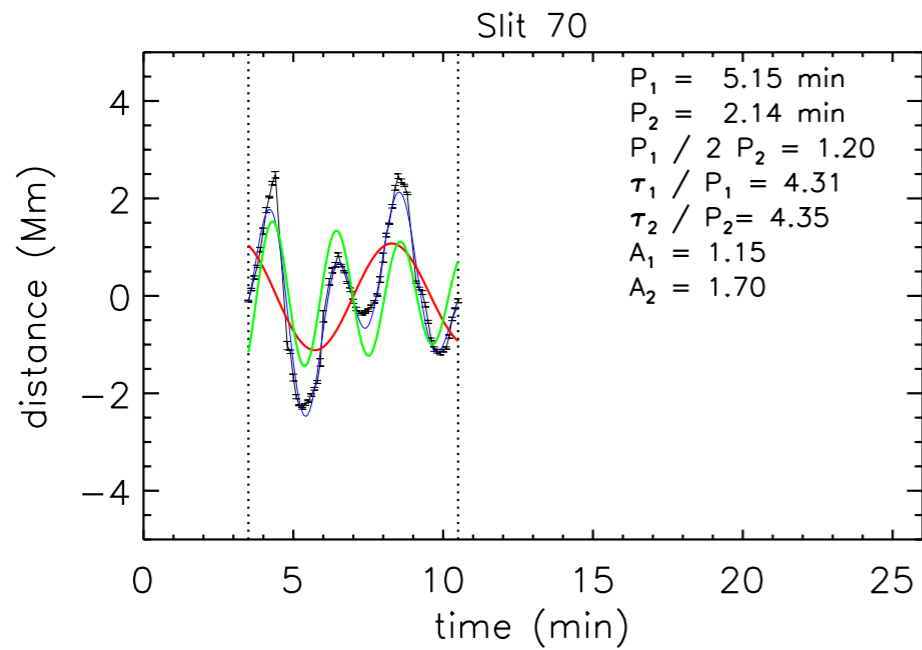
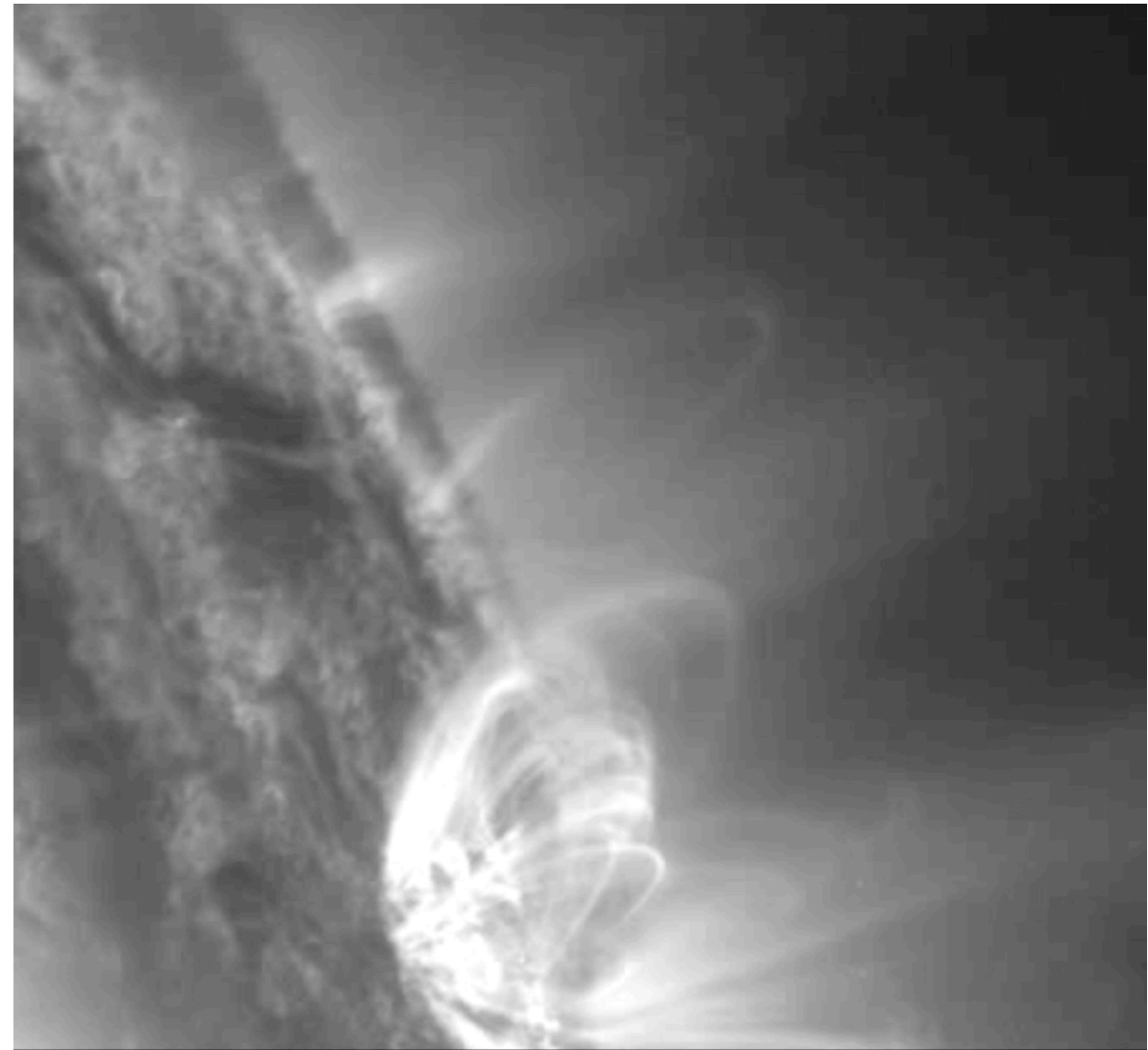
Second harmonic standing kink mode



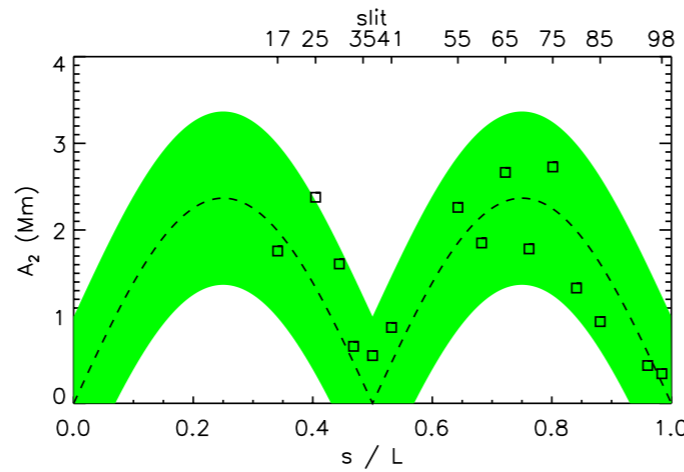
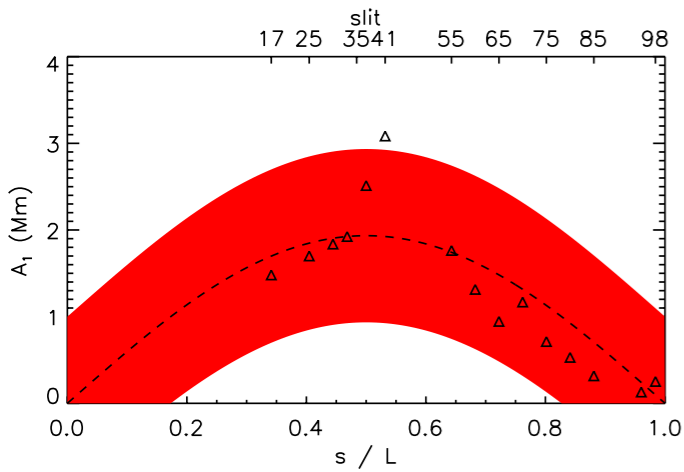
Fundamental kink mode



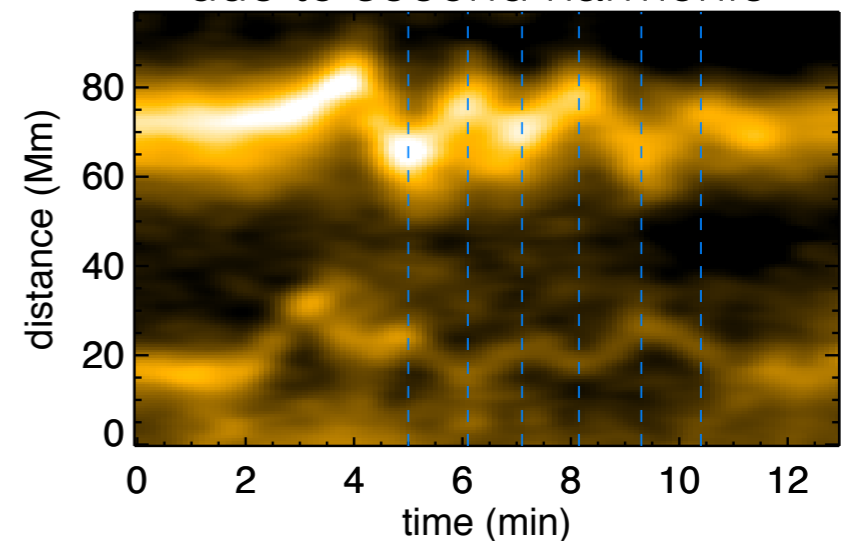
Second harmonic



Anti-phase motions of loop legs due to second harmonic



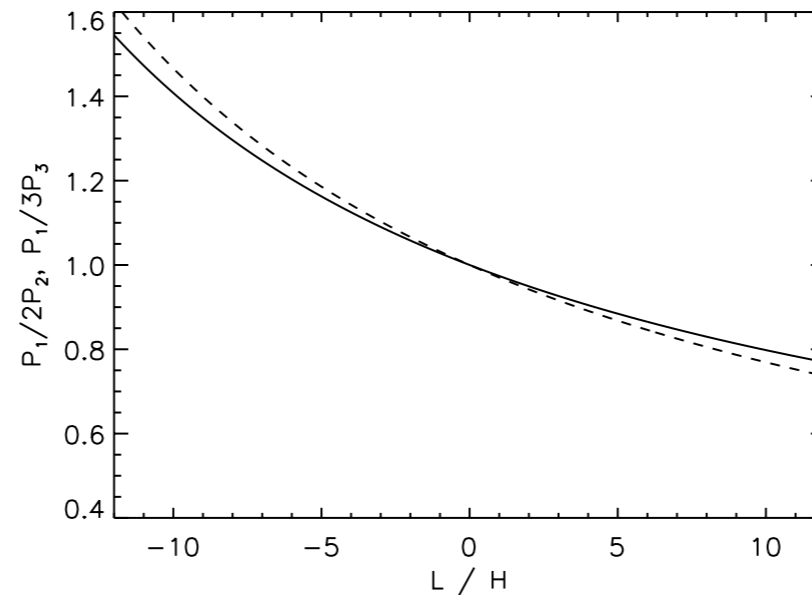
$$P_1 / 2P_2 = 1.15 \pm 0.22$$



Seismology of longitudinal structuring using period ratios

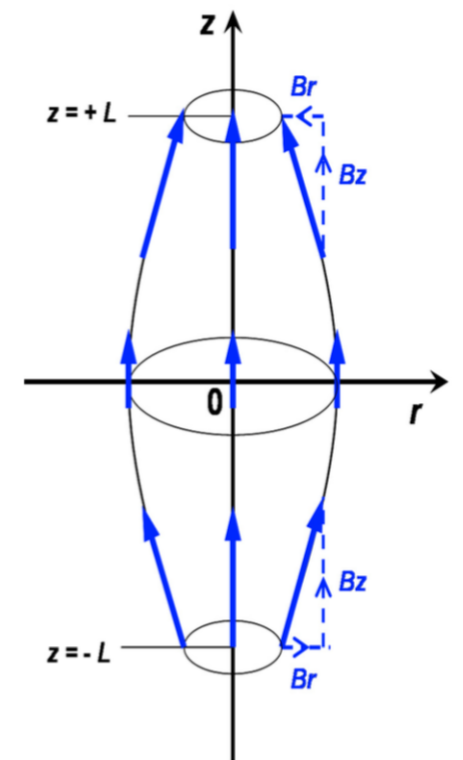
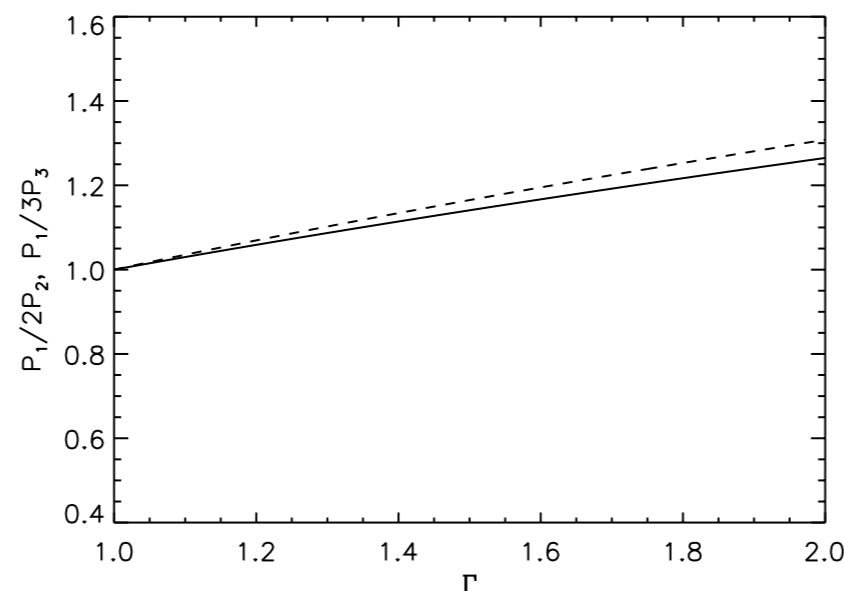
- Standing kink modes usually in long wavelength limit (dispersionless) but...
- Effect of density stratification (e.g. Andries et al. 2005; Safari et al. 2007; McEwan et al. 2008)

$$\rho \propto \exp(-L/\pi H \sin \pi z/L)$$



- Effect of Loop expansion (Verth & Erdélyi 2008, A&A, 486, 1015)

$$\Gamma = \frac{r_{\text{apex}}}{r_{\text{foot}}}$$



- Compare models using Bayes factors (e.g. Arregui et al. 2013, ApJ, 765, L23)

Summary

- Kink oscillations of coronal loops are damped by coupling of energy to Alfvén waves
- Numerical simulations and analytical theory predict:
 - Gaussian damping envelope for low density contrast loops
 - Exponential damping envelope for high density contrast loops
- Evidence of Gaussian damping regime recently discovered in SDO/AIA data
- Observation of both damping envelopes has allowed the transverse loop structure to be seismologically calculated for the first time
- Transverse structure essential for understanding corona e.g. improved estimates of magnetic field strength, heating rate based on phase mixing